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THEORY OF AND COMPUTER AIDED NEUTRON ACTIVATION ANALYSIS

by

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Submitted to the Institute for Graduate Studies in Science and Engineering
in Partial Fulfillment of
the Requirements for the Degree of
Master of Science
in
Nuclear Engineering

Bogazici University Library



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1983

KEYWORDS

NEUTRON ACTIVATION ANALYSIS

GAMMA SPECTRA

GAMMA SPECTROMETERS

LI DRIFTED GE DETECTORS

PROGRAMMING

QUANTITATIVE ANALYSIS

A B S T R A C T

Neutron activation analysis, NAA, has become a widely used and important analytical technique, not only in trace analysis but also in determining alloy elements and minor constituents. The micro qualitative and quantitative analysis can usually be performed by NAA in various fields, such as biology, criminology, geochemistry, cosmochemistry, industry etc. Thus, the field of application of the method is very wide.

In this study, first the basic concepts and principles of NAA is introduced. The detection methods and various types of detectors are discussed. The use of computers in NAA problems are also given. Then, a predeveloped computer code, Gretel, which is set up for routine batchwise processing of spectrometric data obtained by GE(Li) detectors is introduced. This code, programmed originally for the IBM 370/165 computer, is adopted to the UNIVAC 1106 computer in BU for further studies. The UNIVAC version of the program, the input and the output data are finally recorded onto a magnetic tape.

Also, another computer code prepared for NAA, Corgam, which was studied previously is introduced briefly. The comparison of these two codes are given at the end of this study.

Ö Z E T

Günümüzde Nötron Aktivasyon Analizi, NAA, yöntemi ile çeşitli malzemelerde, bu malzemeleri teşkil eden elementlerin cins ve miktarlarının tayini büyük önem kazanmıştır. Diğer analitik yöntemlere kıyasla daha hassas sonuçlar veren bu yöntem, bugün biyoloji, kriminoloji, jeokimya ve kosmokimya gibi birçok alanlarda kullanılmaktadır.

Bu çalışmada, NAA'nın temel kavram ve prensipleri, analizde kullanılan spektrometre sistemleri incelenmiş ve değişik detektör tipleri tanıtılmıştır. Nötron aktivasyon analizi sonuçlarının değerlendirilmesinde karşılaşılan güçlüklerden biri, uzun zaman alan hesaplamalardır. Bilgisayarların hızlı hesaplama yetenekleriyle bunların giderilmesinden ayrıca söz edilmiştir. Bu amaçla, daha önce hazırlanmış ve Ge(Li) tipi spektrumları inceleyen bir bilgisayar programı, Gretel, tanıtılmış ve programda gerekli bir takım düzeltmeler, çizim altprogramlarının program mantığını bozmadan çıkartılması gibi, yapılarak yalnız tasarlandığı üzere IBM 370/165 tipi bilgisayarlarla değil, UNIVAC-1106 tipi bilgisayarlarla da kullanılabilmesi sağlanmıştır. Program bu yeni şekliyle bir manyetik teybe kaydedilmiş ve gelecek kullanımlar için saklanmıştır.

Ayrıca, NAA problemleri için hazırlanmış bir diğer bilgisayar programı olan Corgam kısaca tanıtılmış ve bu iki programın kıyaslaması da çalışmanın sonunda sunulmuştur.

ACKNOWLEDGEMENT

The author would like to express her gratitude to Doç. Dr.Fahir I.Borak, her thesis advisor, who offered invaluable and encouraging guidance and assistance throughout this work. Special thanks are extended to Prof.Dr.Turan Enginol, Head of NE Department, and to Dr.Vural Altın for their constructive criticisms and suggestions. A special extension of thanks is also given to all staff of the BU.computer center for their help. A very special expression of love and gratitude is extended to her parents for their encouragement and unconditional help during not only this work but also previous studies.

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I. INTRODUCTION

1.1. GENERAL

The nuclei of atoms are stable only when they contain a certain number of neutrons and protons. In other words, the number of excess neutrons in the nucleus of an atom determines whether that atom is radioactive or not. However, all elements found in nature with atomic number greater than 83 (bismuth) are radioactive. They belong to chains of successive decays, and all the species in one such chain constitute a radioactive family or series. Besides the naturally occurring radioactivity which was first discovered by H. Becquerel, [Fr-62, Ha-69], there are also some radioactive nuclei artificially produced.

The name of radioactivity implies that radiation is emitted. As will be seen in the next chapter, radiation may be in the form of α , β and γ rays. Since these radiations penetrate and cause ionisation and damage in the body, they may be used for various types of testing, such as tracer techniques. Activation analysis is the most sensitive tool for analysis of trace elements in many materials.

The activation analysis may be defined as a method of measuring concentrations of constituents in a given sample by measuring the characteristic radiations emitted by the radioactive nuclei resulting from nuclear transmutations. First, originally Hevesy and Levi used thermal neutron activation to determine the concentration of dysprosium in impure

yttrium, in 1936 [He-36]. Then, increasing application of radionuclides to analytical problems led to a rapid growth of activation analysis. The activation product provides a specific method for its identification and measurement. For example, for a given problem, a specific method can be selected based on the properties of the sensitivity required in the analysis, the presence of interfering reactions, the location of a suitable irradiation facility and cost.

The activation method is now used frequently in combination with gamma-ray spectrometry in order to simplify or even to eliminate chemical separations. The digital computers are also used for the analysis of gamma-ray spectra and many computer programs are prepared not only to resolve the gamma-ray spectra, but also to calculate the weight of the activated elements from their induced radioactivities. These programs are prepared especially for three types of systems [Le-65, Be-73]. These systems may be:

1. A large centralised computer facility which can analyse batchwise punched tape on which gamma spectrometric data have been collected.
2. A time-shared teleprocessing system linking various spectrometric installations to a centralised computer.
3. A minicomputer used at the laboratory site for both data acquisition and analysis.

The programs prepared for these systems, deal with especially NaI(Tl) type or Ge(Li) type spectra.

Through its nuclear character, the neutron activation technique is a powerful analytical tool not only in trace

analysis but also in determining alloy elements and minor constituents. Therefore, the field of application of the method is very wide. It can be used to determine traces in ultrapure elements, in biological materials and minerals, in industrial product identification by trace characterization, in geochemistry and cosmochemistry, etc [Ko-69, Ch-60, Fa-73].

In this thesis, all these subjects mentioned above will be discussed briefly in the separate subchapters. In addition, a computer code for Ge(Li) type spectra, Gretel, will be examined in a detailed form and compared with another one prepared for NaI(Tl) type spectra.

1.2. OUTLINE OF THE THESIS

The main purpose of the thesis is to investigate the theory and application of the neutron activation analysis, NAA, and make a detailed study of Gretel which is a computer code used to analyze the gamma spectra obtained by Ge(Li) detectors.

The thesis is composed of five chapters and a number of appendices. These can be summarized as follows:

The first chapter is the general introduction of the thesis, covering its purpose and results, and the method of neutron activation analysis.

The second chapter is mainly based on the theory of NAA; the basic principles of activation analysis and gamma-ray spectrometry are examined in detail. First, in the subchapter 2.1 natural radioactivity, radioactive decay and radioactive equilibrium are discussed briefly. In the subchapter 2.2 interaction of neutrons is studied. In the 2.3 interaction of gamma rays with matter and the main effects of this interaction, such as photoelectric effect, Compton scattering and pair-production, are introduced briefly. Then, the detection of gamma rays by means of different types of counters, especially scintillation and semiconductor counters are explained in detail.

The subchapters 2.5 and 2.6 introduce the gamma-ray spectrometry and the basic principles of NAA. Since the use of computers in NAA is an important tool, the last subchapter examines this subject in a summary form.

In chapter 3, Gretel is introduced and investigated in detail; theory of the code, analysis of gamma spectra obtained by Ge(Li) detectors with computer, the smoothing procedure to minimize the influence of the statistical counting fluctuations and the principles for searching of the peaks and the computation of the peak areas found by two different methods are discussed in the following subchapters. In the last subchapter 3.3, the modification of the code which was recorded originally by the use of IBM 370/165 to UNIVAC-1106 is given in detail.

In the following chapter 4, introduction of another code, CORGAM, that has been studied in a previous thesis is given and comparison of these two codes in a tabular form is made. Then, in the chapter 5, discussion and conclusion of the thesis and some suggestions for the use of Gretel are given in a short form.

Finally, some mathematical methods used in Gretel and the listing of Univac version of the code, the listing of input and output data are given as appendices A to K.

1.3. NEUTRON ACTIVATION ANALYSIS

Activation analysis is based on the principle that when a material is irradiated by the nuclear particles produced in a nuclear reactor, particle accelerator or other suitable source, some of the atoms present in the material will interact with the bombarding particles and be converted into different isotopes of the same element or isotopes of different elements, depending on the nature of the bombarding particles. In many cases, the isotopes produced are radioactive. The subsequent emission of beta or gamma radiations from these nuclei is characteristic of the particular isotope. If each different induced radioactivity can be distinguished or separated from all other radioactivities produced, then the amount of each radioactivity is a measure of the quantity of the parent isotope present in the material [Wa-63].

The sensitivity of the neutron activation analysis method can be considered to be about parts per billion, ppb, for reactor irradiations and parts per million, ppm, level for irradiations with neutron generators. In fact, the sensitivity and the accuracy of the analysis are frequently based on the ability to distinguish the radiation of the radioisotopes of interest from the other activated constituents of a sample. Besides NAA, analysis with photons and charged particles can be considered [Ho-71].

Besides the high sensitivity of thermal activation analysis, this technique offers a number of other advantages such as the simplicity of the nuclear reaction, transparency of most materials to thermal neutrons and non-destruction of the sample. In principle, in most cases only neutron capture reactions have to be considered and there is no change in the chemical nature of the irradiated samples.

Unfortunately, thermal neutron activation can not solve a number of important trace analysis problems, for instance those of the light elements, such as carbon, nitrogen and oxygen. These problems can sometimes be solved by applying photons or charged particles such as protons, deuterons, etc. If neutrons and protons essentially analyze the whole mass of the sample, this is not true for charged particle analysis, as these have a very short range. While all analytical techniques require some chemical treatment before the analysis; the NAA may require the chemical treatment only after irradiation and often after removal of surface contamination. One further advantage of neutron activation analysis, NAA, is in many cases, the selectivity or even specificity of its nature. First, the activation is a reaction of the nucleus and thus independent of the chemical state of the considered element. Further, the relatively simple nature of the nuclear reaction allows one to establish with great certainty the isotope or isotopes which gave rise to the measured species. Besides chemical criteria, nuclear ones can increase the certainty of identification of elements, as, for instance, half-life measurements, type of Disintegration and energy measurements of the emitted radiation.

Gamma-ray spectrometry appears to be a powerful tool in this respect. By applying gamma scintillation spectrometry to activation analysis, it becomes possible to analyze simultaneously a number of different elements. With the use of germanium lithium drifted detectors, gamma spectrometry has developed in a spectacular way. These detectors have a very high resolving power [So-72, Ho-71].

The procedure is often summarized as follows: Irradiate a sample and a standard under the same conditions; count

both under the same conditions; the ratio of these rates will be equal to the ratio of the masses of the elements of interest. If one takes activation in a thermal reactor as an example, it should be remembered that among many other parameters the neutron energy varies from a thermal energy of 0.025 eV up to about 20 MeV. Consequently activation not only occurs with thermal neutrons, but also with epithermal and fast neutrons. Besides, the activation can vary from place to place in the reactor. The flux gradient may not be the same in standard or sample or even over the entire sample itself. Many other complexities may be met during activation, such as unsuspected reactions, production of the isotope of interest by second order reactions or radioactive growth of a daughter isotope, absorption of radiations, instability of the instrumentation, non-linearity in detector response, etc.

From this short list of possible sources of error it is concluded that activation analysis is a complex analytical technique and that experience is required to obtain satisfactory results. But all these complexities have not prevented it from developing rapidly. At the beginning, it was mainly used in the field of trace determination in high purity materials; applications are now increasingly used in different fields such as biological sciences, geochemistry, criminology, environmental sciences, etc. Since 1969, cosmochemistry is also added to these fields. For instance, the analysis of Apollo Lunar material is widely performed by activation analysis [Ho-71]. More and widespread applications are to be expected in the future.

2. THEORY

2.1. RADIOACTIVITY

Some nuclei undergo spontaneous decay or disintegration. This process is known as radioactivity. The discovery of radioactivity goes back two historical events: 1) The discovery of x rays by W.C. Roentgen in 1895, 2) In 1896, H. Becquerel reported his experimental results; after exposure to bright sunlight crystals of the uranyl double sulfate emitted a radiation which blackened a photographic plate after penetrating black paper, glass and other substances. Then, Pierre and Marie Curie concluded that the uranium rays were an atomic phenomenon characteristic of the element and not related to its chemical or physical state. They, then introduced the name "radioactivity" for the phenomenon [Fr-62, Ha-65].

There are three kinds of radiations emitted by the radioactive substances. These rays can be distinguished from each other in two ways: First, by the difference in the ease with which the rays can pass through matter. Second, by the direction in which their path is bent by the application of a magnetic field. The most easily absorbed rays are α particles which are deflected by both electric and magnetic fields. This deflection shows that its charge is positive. It is also concluded that it is composed of a doubly ionized He atom, He^{++} . The second type of radiation requires roughly 1000 times as much matter to bring it to rest. These particles, named β^- , are reflected much more readily by magnetic fields than α

particles but in the opposite direction. They, therefore, carry a negative electrical charge. The properties of these particles are similar to those of fast moving electrons. The third type of radiation, called γ ray is even more penetrating than β^- particles. The gamma-rays can only be stopped by several centimeters of lead. These rays are not deflected by magnetic or electric fields. They are electromagnetic radiations of the same kind as X rays, light and radio waves, but of very short wavelength.

The emission of α and β particles cause a change in atomic number or mass number of the nucleus. If a nucleus emits an α particle, whose charge is 2 and whose mass is 4 on the atomic weight scale, then obviously the nucleus that remains after the event will carry two units less of positive charge and four units less of mass. Since a β^- particle is negatively charged, its emission will cause an increase of one unit in the nuclear positive charge. On the other hand, the emission of a γ -ray produces no change in atomic number of the nucleus; but only a decrease in its energy content. In most cases, the emission of gamma-rays is a secondary consequence of either an α or a β decay. Following the initial decay, the residual nucleus is frequently left with excess energy which may emit one or more γ -rays.

Besides the naturally occurring radioactivity, some radioactive nuclei are artificially produced. First in 1934 I. Curie and F. Joliot discovered that boron, aluminum and magnesium could be made radioactive by bombardment with the α rays [Fr-62]. At that time, many laboratories put into operation devices for the acceleration of hydrogen ions and helium ions to energies at which nuclear transmutations were produced. Then, many artificially produced radioelements have found important applications in some fields, such as che-

mistry, physics, biology, medicine and engineering.

2.1.1. Radioactive Decay

Radioactive decay is a statistical process. The probability of one nucleus in a group of nuclei decaying in a finite time interval is independent of the time. In other words, the probability of decay at two different times t_1 and t_2 will be the same.

The probability P of decay of one nucleus in a time interval Δt can be written as,

$$P = \lambda \Delta t \quad (2.1)$$

where λ is decay constant of the isotope. The number of nuclei decaying from a group of N nuclei will be,

$$-\Delta N = NP = N\lambda \Delta t \quad (2.2)$$

If the time interval is very small, eg. (2.2) can be written as,

$$-dN/dt = \lambda N \quad (2.3)$$

integrating this equation and taking N_0 as the number of nuclei present at time $t=0$, one obtains,

$$N = N_0 e^{-\lambda t} \quad (2.4)$$

This equation shows that the number of radioactive nuclei decreases exponentially with time.

According to eg. (2.3), the number of nuclei decaying per unit time is equal to (λN) . This is called the activity of the sample [Ou-75, Ha-65]. Multiplying both sides of eg. (2.4) by λ ;

$$\lambda N = \lambda N_0 e^{-\lambda t}$$

or,

$$A = A_0 e^{-\lambda t} \quad (2.5)$$

where A is the activity at time t and A_0 is the initial activity at time $t = 0$. Equation (2.5) shows that the activity also changes with time exponentially.

2.1.2. Radioactive Equilibrium

Among the radioactive nuclei there are many examples of chains of decays in which a radioactive parent decays to a daughter which is also radioactive, and so on. In such a case, it is especially important to know how much of that is produced.

In the case of a radioactive parent decaying to form a radioactive daughter, the daughter product is produced at the rate of the parent's decay. However, the daughter itself decays, and this decay is proportional to the amount of daughter that is present.

Considering N_1 , the number of parent atoms and N_2 , the number of daughter atoms as functions of decay constants and time, it is possible to write the decay rate of the parent, according to eg. (2.3), as,

$$dN_1/dt = -\lambda_1 N_1$$

or taking $N_1 = N_{1,0} e^{-\lambda_1 t}$

$$dN_1/dt = -\lambda_1 N_{1,0} e^{-\lambda_1 t} \quad (2.6)$$

The analogous equation for the daughter is,

$$dN_2/dt = R_F - \lambda_2 N_2 = \lambda_1 N_{1,0} e^{-\lambda_1 t} - \lambda_2 N_2 \quad (2.7)$$

where R_F represents the rate of formation of the daughter and it is equal to $(\lambda_1 N_{1,0} e^{-\lambda_1 t})$; the second term $(-\lambda_2 N_2)$ represents the decay term.

Rearranging the equation (2.7) which is a 1st order linear differential equation, for the solution of which(*), one obtains,

$$N_2 = \frac{\lambda_1}{\lambda_2 - \lambda_1} N_{1,0} (e^{-\lambda_1 t} - e^{-\lambda_2 t}) \quad (2.8)$$

Just what happens depends largely on the ratio of the half lives:

a) In the case of the parent having a much longer half-life than the daughter, $(\lambda_2 \gg \lambda_1)$, λ_1 is very small and can be allowed to approach zero in eq. (2.8). The resulting equation which gives the amount of the product nuclide as a function of time, takes the form,

$$N_2 = \frac{\lambda_1}{\lambda_2 - \lambda_1} N_{1,0} (1 - e^{-\lambda_2 t}) \quad (2.9)$$

(*) Solution of this equation is given in the appendix B in more detail.

or,

$$N_2 = K(1 - e^{-\lambda_2 t}) \quad (2.10)$$

The significance of K becomes apparent if t approaches to infinity. Then $N_2 = K$ and K is the equilibrium value that N_2 approaches as time goes on.

The form of the build-up or saturation factor $(1 - e^{-\lambda t})$ is shown in the figure (2.1)

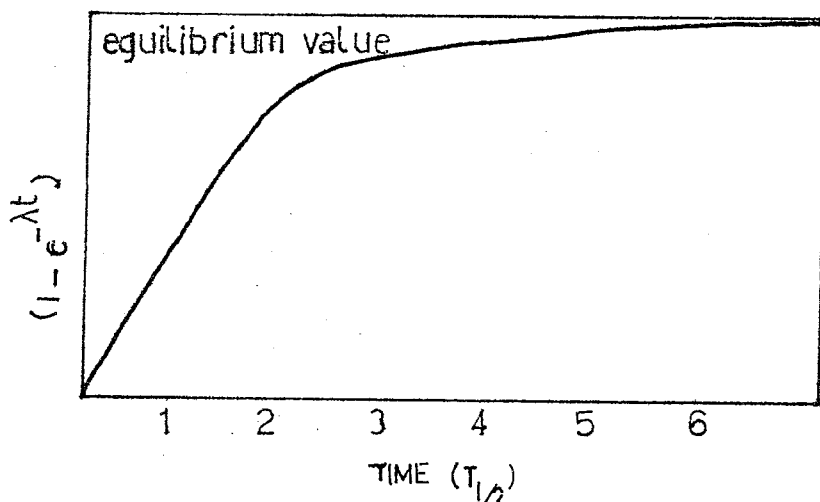


Fig. 2.1. Build up of a radioactive nuclide created at a constant rate [EL-59].

b) If $\lambda_2 > \lambda_1$, but not great enough for one to ignore the decay of the parent during the growth of the daughter, then the situation is called transient equilibrium. Fig.(2.2) shows the growth and decay of a parent and its daughter.

As shown in fig.(2.2), the amount of daughter at first increases, then reaches a maximum, and finally decreases at the same rate as the parent with which it is in transient equilibrium(*)).

(*) Detailed explanations for that subject can be supplied from many text books, such as [EL-59, Ha-69, Wi-66, Se-67, Ou-75].

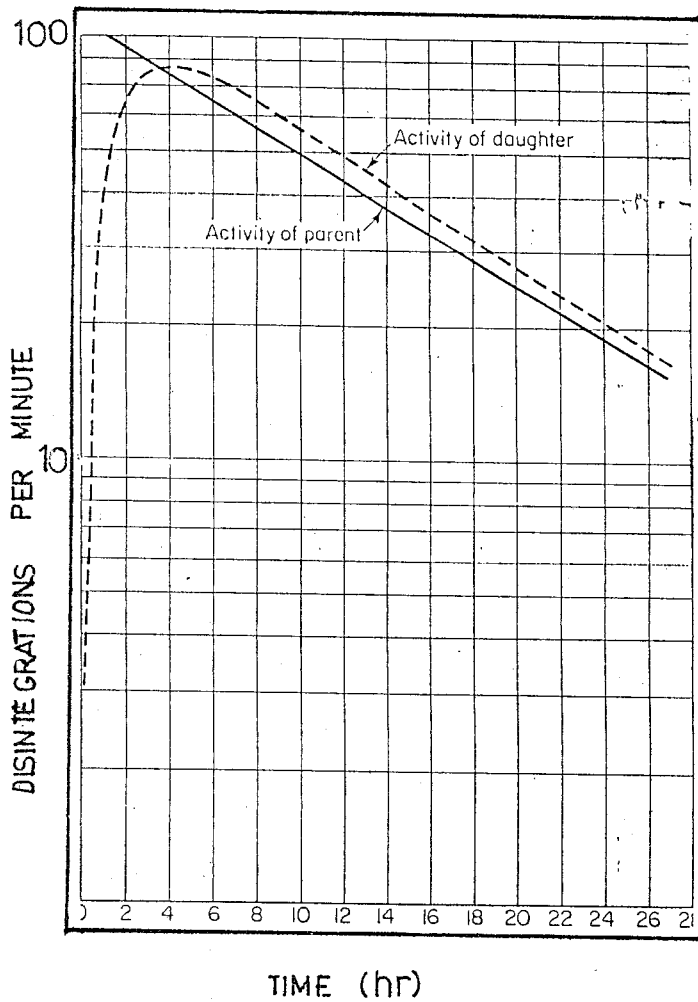


Fig. 2.2. Growth and decay of daughter in transient equilibrium with parent.

2.2. INTERACTION OF NEUTRONS WITH MATTER

Since neutrons are not charged and are much heavier than electrons, they do not interact significantly with electrons. They are also not repelled by electric fields; even very low energy (thermal) neutrons readily penetrate and react with nuclei. Therefore, neutrons are not easily detected by most counting instruments because they do not ionize. They must be detected by their secondary effects. In other words, for the detection of neutrons, a nuclear reaction must have reasonably high probability and must result in the production

of some other form of easily detectable radiation, such as charged particles or gamma-ray. Suitable reactions(*) fall into four general categories as follows:

Activation reactions: These reactions result in the production of excited nuclei with reasonable long life times. Samples of suitable materials, usually thin foils, are exposed to neutron irradiation. They are then removed from the neutron flux and their radioactivity is measured usually by β or γ counting. The most important of those reactions is the (n,γ) reaction. This process is also known as radiative capture, since one of the products of the reaction is instantaneous gamma-rays.

Activation reactions are very useful for the detection of neutrons. That subject will be studied in detail in the later subsections.

Transmutation reactions: Neutrons may be captured by nuclei and cause the prompt emission of charged particles such as protons, deuterons or alpha particles. In that case, the incident neutron must supply sufficient energy to overcome the force binding the charged particle to the compound nucleus for the reaction to take place. Besides those, sometimes two or more neutrons are emitted when a nucleus is struck by a high energy neutron. These processes here are of the $(n,2n)$ or $(n,3n)$ types. A closely related process is the (n,pn) reaction which also occurs with highly energetic incident neutrons.

Fission: When a neutron collides with a heavy nucleus, the nucleus splits into two large fragments with the release of considerable energy. This process is known as fission. The fission process provides a technique for detecting neutrons

(*) These explanations and more detailed form of them can be obtained from many text books, such as [Ch-61, La-72, Wa-62, Wi-66].

in the presence of very intense gamma radiation.

Scattering: When a moving particle collides with another particle, the kinetic energy is exchanged between them according to the laws of conservation of energy and momentum. There are two types of collisions: First, if the potential energy of the system remains unchanged, kinetic energy being conserved during the collision, the phenomenon is called elastic scattering. Second is the inelastic scattering in which one of the particles is left in an excited state after collision. Here momentum and total energy of the particles before and after collision are conserved. However, kinetic energy is not.

In each scattering event, part of the kinetic energy of the neutron is transferred to initially stationary and heavier nucleus; as a result of that they are slowed down. Then the slowed down neutrons become more effective in causing nuclear reactions.

The emission of gamma radiation, as a result of the reactions described above is due to the release of excess energy of the excited nucleus. The energy of that gamma-ray is equal to the difference between the energy levels of the excited and stable states of the nucleus. In some cases two or more gamma-rays of discrete energies are emitted in cascade. Those gammas are usually emitted within 10^{-13} sec. Then β particle, α particle and electron capture decays occur. Sometimes that excited state has a characteristic half-life of its own. Whenever that half-life is of the order of 10^{-7} to 10^{-6} sec, or longer; the excited state is called an isomer of the product. The subsequent decay of that isomer by the emission of a gamma ray is called an isomeric transition. A typical example is given for Co^{60} in the figure (2.3).

The ^{60}Ni nucleus formed by the β^- decay of ^{60}Co is rapidly converted by the emission of 2.50 MeV gamma-ray. This conversion takes place by the emission of a photon of 1.17 MeV and immediately there-after a photon with 1.33 MeV.

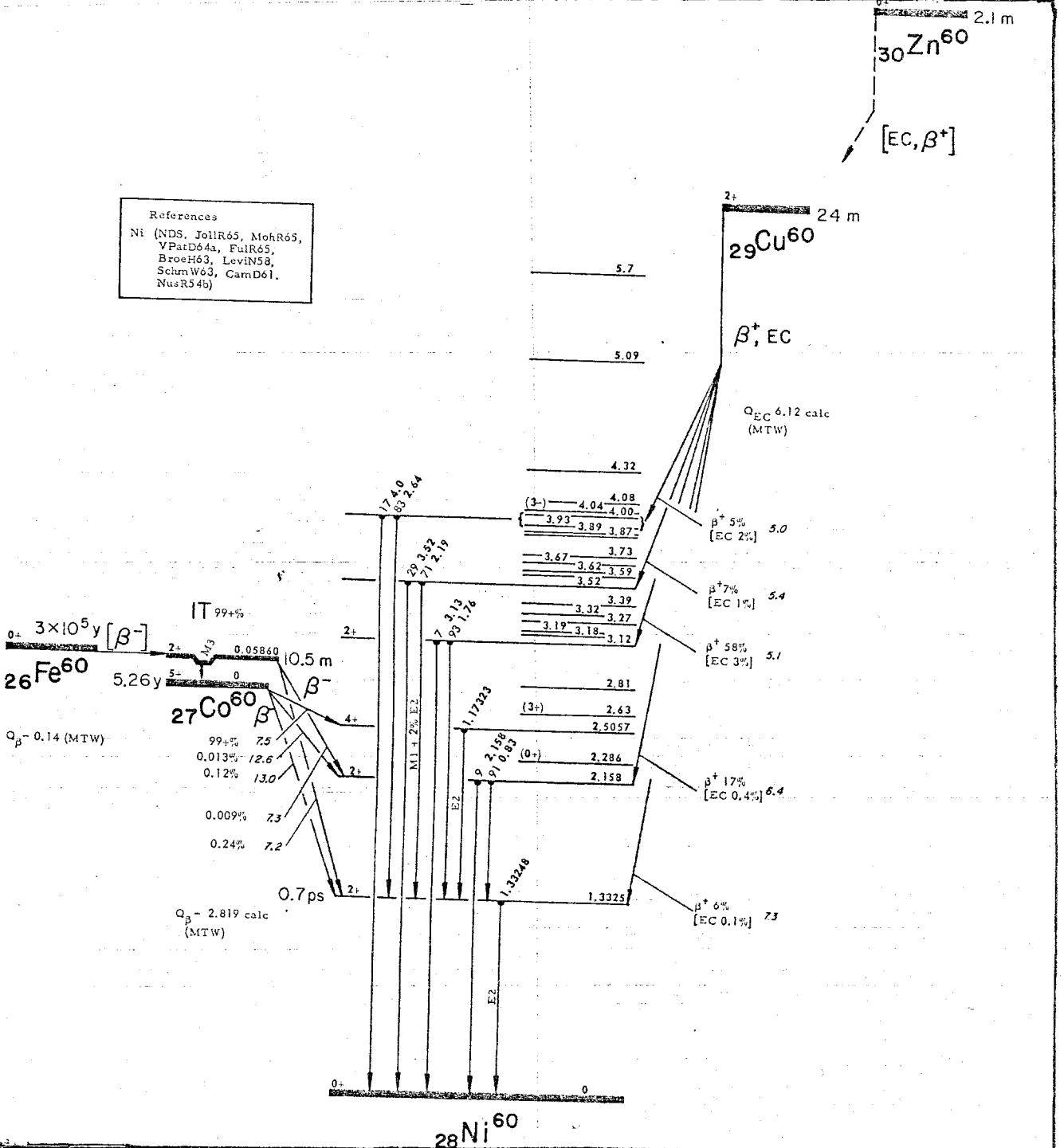


Fig. 2.3. Decay scheme of ^{60}Co [Le-68].

A more complicated one is ^{152}Eu . It has an average level spacing of about 5 keV above 100 keV. It decays to the spherical nucleus ^{152}Gd and to the deformed nucleus ^{152}Sm . Deformed and spherical states are expected close to the ground state. It is very difficult to determine the configurations of these levels in ^{152}Eu [Eg-74]. The decay scheme of ^{152}Eu can be seen in figure(2.4).

The detection of those gamma rays by suitable detectors and the analysis of that spectrum may be used to identify the radiating nuclide, its amount, etc. In the following subchapters, detection of gamma-rays, analysis of spectra obtained from unknown nuclei and, qualitative and quantitative analysis techniques will be examined in detail.

2.3. INTERACTION OF GAMMA RAYS WITH MATTER

Electromagnetic radiation of high photon energy usually originates from excited atomic nuclei and is called γ ray or X ray. These rays are usually classified according to their mode of origin, not their energy.

There are a number of processes which can cause gamma-rays to be scattered or absorbed. A catalogue of the possible processes by which the electromagnetic field of the γ -ray may interact with matter has been put in the following systematic form by Fano [Fa-53];

Kinds of interaction:

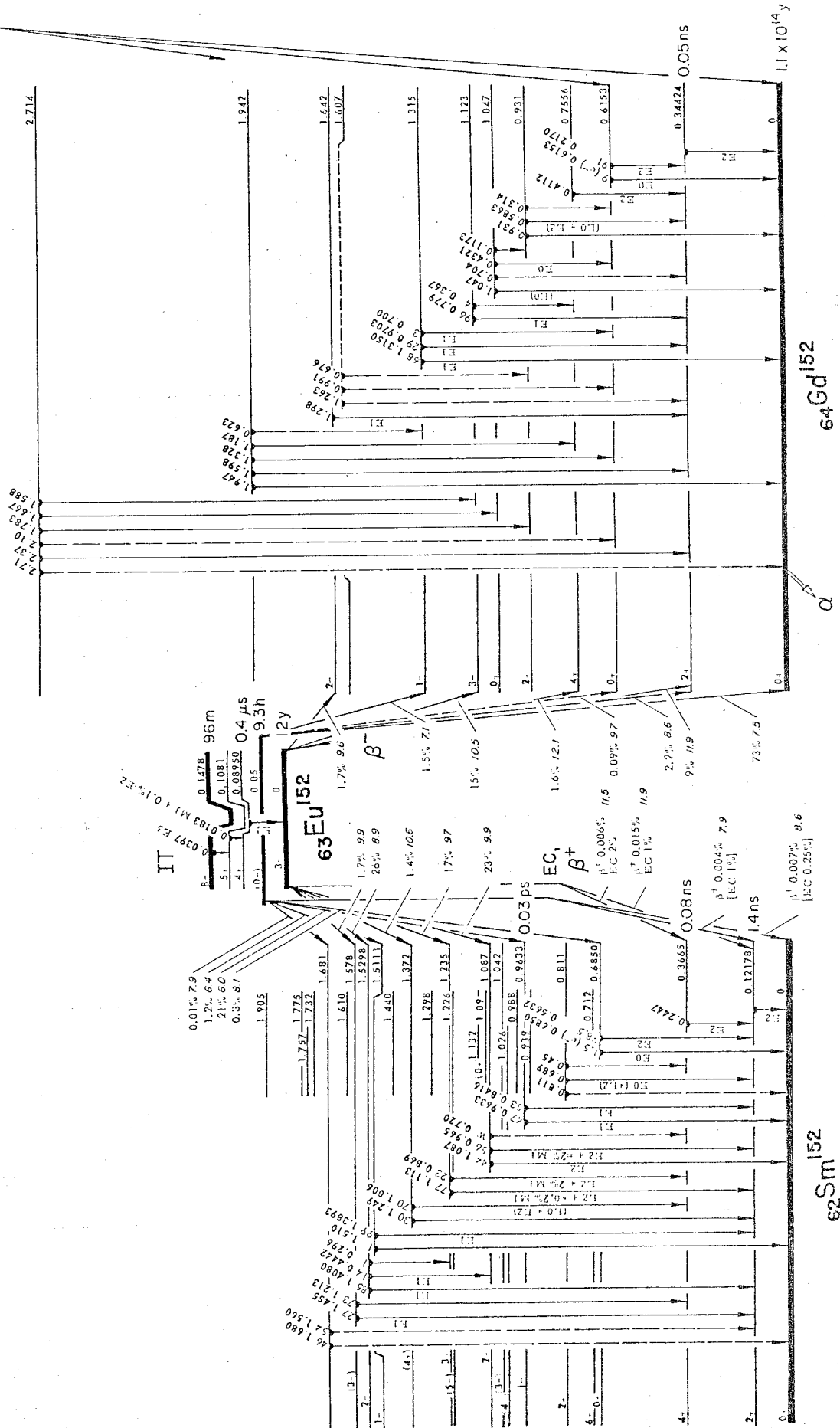
- 1- Interaction with atomic electrons
- 2- Interaction with nucleons
- 3- Interaction with the electric field surrounding nuclei or electrons

Fig. 2.4 Decay scheme of ^{152}Eu

[Le-68]

^{152}b

β^+, EC



4- Interaction with the meson field surrounding the nucleons.

Effects of interaction:

- a) Complete absorption
- b) Elastic scattering (coherent)
- c) Inelastic scattering (incoherent)

There are 12 ways of combining these interactions and effects, thus in theory there are 12 different processes by which gamma-rays can be absorbed or scattered. Many of these processes are infrequent. Three of them, photoelectric effect, compton scattering and pair production, are the most important effects. Figure(2.5) illustrates the relative importance of the three main interactions.

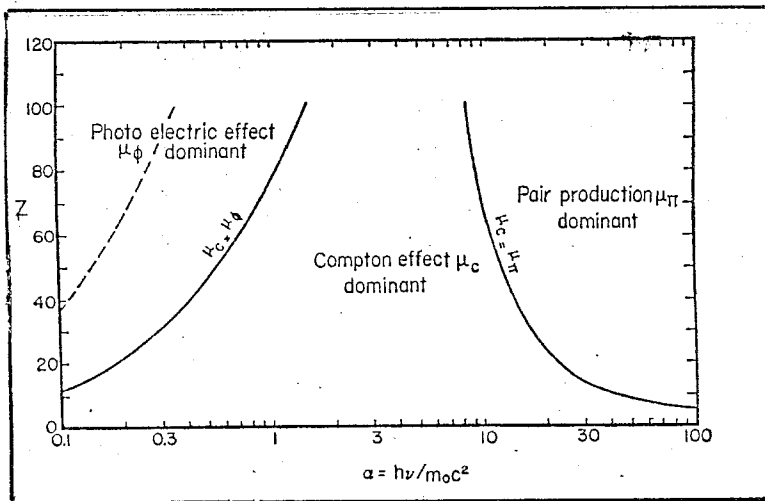


Fig. 2.5. Relative importance of the three main γ -ray interactions as a function of the absorber atomic number Z and the photon energy in the units of M_0C^2 are given (the dotted curve indicates the Rayleigh scattering) [Aj-60].

These effects will be examined in the following sub-chapters.

The minor effects which are of interest in special cases are as follows [Ev-55];

Rayleigh Scattering (1b): Gamma rays scattered by small angles resulting in only a small recoil, especially when their energy is low. The recoil is then often absorbed by a whole atom or molecule. This scattering is more likely for high Z than for low Z materials.

Thomson Scattering (2b): Thomson scattering can combine coherently with Rayleigh scattering. Because of the large mass of the nucleus, the effects are small but appear to have been detected.

Delbruck Scattering (3b): Delbruck scattering is due to virtual electron pair formation in the field of the nucleus. The effect is extremely small.

Nuclear Resonance Scattering (2c): This type of scattering involves the excitation of a nuclear level by an incident photon, with subsequent reemission of the excitation energy.

Photodisintegration of nuclei (2a): It is possible whenever the photon energy exceeds the separation energy of a neutron or proton. Except for $^9\text{Be}(\gamma, n)$ and $^2\text{H}(\gamma, n)$, these effects are generally confined to the high energy region above approximately 8 MeV.

Meson Production (4a): Meson production requires photon energies above 150 MeV.

2.3.1. Photoelectric Effect

Below energies of about 0.1 MeV the predominant mode of γ -ray interaction in all medium and high Z absorbers is the photoelectric process. In this process the photon energy is used to remove one of the electrons from an inner shell of an atom of the absorber element. Clearly this process can occur only if the incoming gamma-ray has an energy higher than the binding energy of the electron to be removed. Thus, the energy of incident photon can be determined from the kinetic energy T of the ejected electron and the binding energy E of the atomic level from which it was ejected by the use of the expression,

$$h\nu = T + E \quad (2.11)$$

where, $h\nu$ represents the energy of incident photon, ν being frequency and h being Planck's constant.

When the experimental mass attenuation coefficient μ/ρ for the photoelectric effect is plotted as a function of λ of the incident photon, a series of jumps corresponding to the binding energy of the different orbits can be observed (e.g. Fig. 2.6).

These energies are given by Mosely's law, [Se-64],

$$E = Rhc \frac{(Z-\sigma)^2}{n^2} \quad (2.12)$$

where,

$$Rhc = 13.605 \text{ eV}$$

Z is the atomic number

n is the quantum number of different electronic orbits (in the K series $n = 1$, in the L series $n = 2$, etc).

σ is the screening constant which has the approximate value 3 for the K shell, 5 for the L shell etc. Figure(2.6) shows these jumps for platinum.

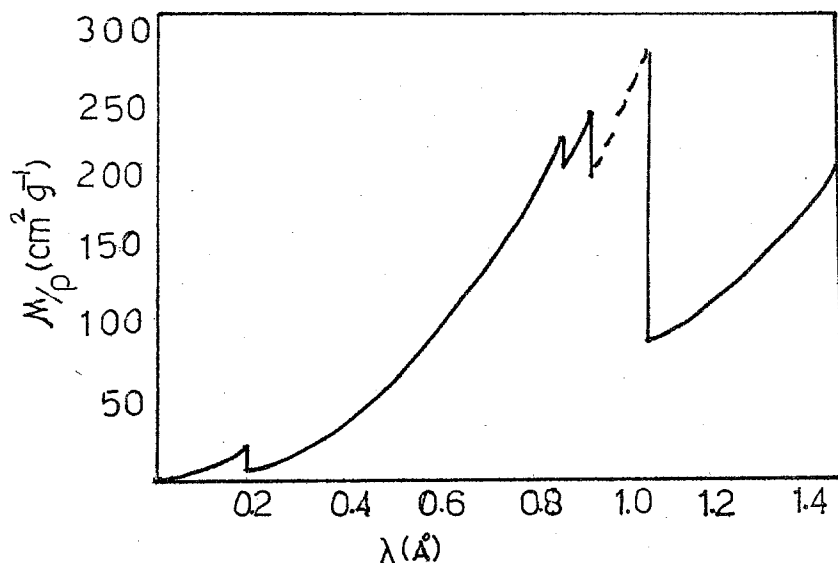


Fig. 2.6. Mass absorption coefficient of platinum as a function of the wavelength of the incident ray [Se-64].

The vacancy created by the ejection of an electron from the inner shell is filled by an outer electron falling into it, and this process may be accompanied by the emission of fluorescent radiation. Alternatively it is possible that no fluorescent radiation is emitted but that an electron from an outer shell is ejected. For instance, a vacancy in a K shell may be filled by an L electron and another L electron emitted with energy $E_k - 2E_L$. Similarly, if the K vacancy is filled by an L electron and an M electron is emitted, the ejected electron will have the kinetic energy $E_k - E_L - E_m$. These and similar processes are called AUGER processes.

2.3.2. Compton Scattering

Compton scattering is inelastic (incoherent) scattering of photons with the atomic electrons. In this type of effect, the photon energy is considered to be large enough compared to the atomic binding energy E , so that the atomic electrons may be considered as free electrons. The incident photon momentum $h\nu_0/c$ is conserved between the scattered photon and the struck electron. Except for the trivial case of zero scattering angle, the direction of the scattered photon is not parallel to the direction of the incident photon. The scattered photon must therefore have a smaller momentum and a smaller quantum energy than the incident photon. The remaining momentum and energy are imparted to the struck electron.

Assuming that the struck electron is at rest and unbound, the schematic representation of the Compton scattering is given in figure (2.7). Where, $h\nu$ represents the energy of the incident photon, $h\nu'$ is the energy of the scattered photon, θ is the angle through which the photon scatters, ψ is the angle at which the electron recoils.

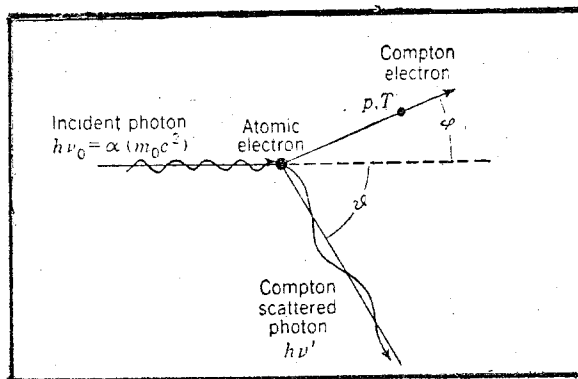


Fig. 2.7. The schematic representation of Compton scattering.

From the conservation of energy,

$$h\nu = h\nu' + T \quad (2.13)$$

where T represents the kinetic energy of scattered electron.

The energy of the scattered photon, $h\nu'$, in terms of $h\nu$, θ and α which is a dimensionless parameter and equal to $h\nu/M_0C^2$, by using the energy and momentum conservation of the system can be written as, [A_j-60].

$$h\nu' = \frac{h\nu}{1+\alpha(1-\cos\theta)} \quad (2.14)$$

It will be seen that $h\nu' = h\nu$ for $\theta = 0^\circ$ for any α and that

$$h\nu' = \frac{1}{2} M_0 C^2 \quad \text{at} \quad \theta = 180^\circ \quad \text{with} \quad \alpha \gg 1 \quad (2.15)$$

$$h\nu' = M_0 C^2 \quad \text{at} \quad \theta = 90^\circ \quad \text{with} \quad \alpha \gg 1 \quad (2.16)$$

Then the kinetic energy of the recoil electron in terms of ϕ and θ is

$$T = h\nu - h\nu' = h\nu \frac{\alpha(1-\cos\theta)}{1+\alpha(1-\cos\theta)} \quad (2.17)$$

2.3.3. Pair Production

At photon energies above $2 M_0 C^2 = 1.02 \text{ MeV}$ pair production is a possible effect. It is the transformation of a gamma ray into an electron-positron pair, also called materialization. When a recoil is absorbed by an electron, the threshold energy required by the conservation of energy and

momentum is $4 MC^2$, and there are two electrons and a positron acquiring appreciable momentum. In a cloud or bubble chamber they form a triplet; [Se-64]. Pairs can also be produced by heavy particle collision, by electron-electron collisions, in mesonic decay and by internal conversion in some gamma transitions.

The angular distribution of the electron-positron pairs is difficult to predict theoretically. If T_- and T_+ are used to denote the kinetic energies of these pairs, they can be given as

$$T_+ + T_- = h\nu - 2 M_0 C^2 \quad (2.18)$$

where $2 M_0 C^2$ is the threshold energy required to create these pairs. [Detailed information about these processes can be obtained in many books, such as Aj-60, Ev-55, Se-64]

2.4. DETECTION OF GAMMA RAYS

Gamma rays are quantized electromagnetic waves emitted as a result of spontaneous changes between energy states of excited nuclei. Gamma rays are usually observed by detecting secondary charged particles, such as photoelectrons ejected from atoms, Compton electrons resulting from interactions with unbound electrons, positron-electron pairs created in the vicinity of a nucleus which lies in the path of the gamma-ray, or photons emitted when the gamma ray interacts with a nucleus.

In many cases one may say that a gamma-ray transition has been observed even though a gamma ray does not leave the nucleus. The most predominant process alternative to gamma -

ray emission is the internal conversion emission of a bound electron. For a gamma-ray transition involving an energy greater than 1.02 MeV a positron-electron pair may be emitted from the nucleus instead of a gamma ray or a conversion electron. This process, called internal pair creation, has the same relationship to external pair production as internal conversion has to photoelectron emission. In special cases when the emission of gamma rays is completely forbidden the transition may take place either by internal conversion, or if this energy is greater than 1.02 MeV, by pair production.

Finally, it may be said that the gamma-ray transitions may be detected by the observation of charged particles external to the atom following interaction with the γ rays. It may also be detected by measuring the internal conversion electron lines emitted from the atom itself. Most of the devices and techniques used for the measurement of gamma rays, actually analyze the secondary electrons. The only exception is the diffraction crystal spectrometer which analyzes the gamma rays directly by Bragg reflection. The ideal device for the detection and measurement of gamma rays would have high efficiency, good resolution and corresponding to each γ ray present, would yield a single peak or line. Since there are many interactions of gamma rays with matter, continuous distributions as well as discrete lines are frequently produced. There are only a few special types of spectrometers which will eliminate continuous distribution and yield only one line for each gamma ray present [A_j-60]. They do so generally with a considerably smaller efficiency than for the ordinary electrons. The most common devices for γ -ray detection are the ionization chamber, proportional counter and Geiger-Müller counter. [Pr-64, En-66, Ou-75]. But the NaI (TL) scintillation counter is the most widely used device which will be examined in the following subchapter. In recent years, with

the development of semiconductor counters which have considerably good resolution, these counters have found wide use. So the second subchapter of this chapter deals with the semiconductor detectors in detail and Li drifted type is especially examined as the computer code studied in this thesis, Gretel, makes use of this type of semiconductor detector.

2.4.1. Scintillation Detectors

One of the oldest methods of nuclear radiation detection is to make use of light scintillations. Although the basic principles remain the same, the techniques used have changed greatly. These techniques are particularly useful in the measurement of the energies of β and gamma rays. Besides these, the solid or liquid phosphors which are used with the scintillation detectors can absorb the β particles or the secondary electrons produced by gammas. Consequently, measurements of the energies of these particles are possible.

Figure(2.8) is a schematic diagram of a scintillator and the electronic arrangement(*) used in a counting system.

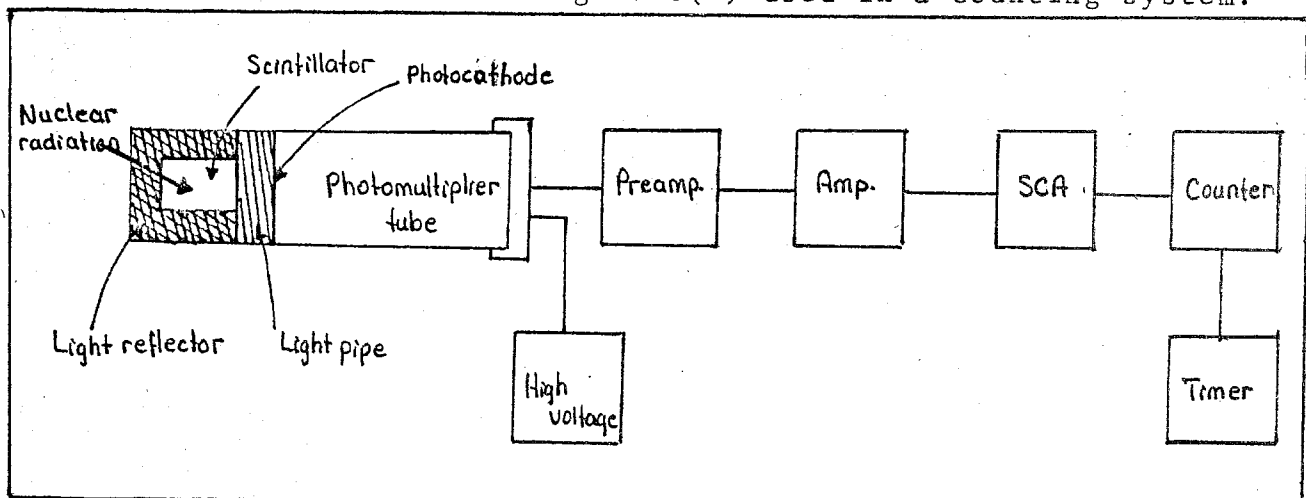


Fig. 2.8. Schematic diagram of a basic scintillation detection system [Bo-82].

(*) This electronic arrangement is taken from NE laboratory notes prepared by F.Borak.

A wide variety of scintillators is in use today. Table 2.1 gives them and some of their characteristics.

TABLE 2.1. Properties of several scintillators [Pr-64]

Materials	Density, g/cm ³	Wavelength of maximum emission, Å	Relative beta-ray pulse height	Alpha- beta ratio, %	Decay time, sec
1. Anthracene crystal.....	1.25	4,400	100	9	3×10^{-8}
2. Trans-stilbene crystal.....	1.15	4,100	60	9	$4-8 \times 10^{-9}$ (β)
3. Liquid phosphors	0.86	3,500-4,500	40-60	9	$2-8 \times 10^{-9}$
4. Plastic phosphors	1.06	3,500-4,500	28-48		$3-5 \times 10^{-9}$
5. NaI(Tl).....	3.67	4,200	210	44	3×10^{-7}
6. LiI(Eu).....	4.06	4,700	70	95	1.2×10^{-6}
7. CsI(Tl).....	4.51	4,200-5,700	55-95		1.1×10^{-6} (β) 0.43×10^{-6} (α)
8. ZnS:Ag) (powder)	4.10	4,500	200	100	$4-10 \times 10^{-8}$ (fast component) $4-10 \times 10^{-5}$ (slow component)

As can be seen in table 2.2, one of the most widely used scintillators is the NaI(TL). When approximately 0.5 % of thallium is added to sodium iodide, a crystal grown from the mixture has the property of emitting light in the visible range of wavelength upon the passage of an ionizing particle. The crystals are generally made cylindrical in shape and the container is provided with a glass window adjacent to one of the flat ends. In order to get the maximum yield of light out of the window, a layer of reflecting material is packed between the surfaces of the crystal and the container walls. To convert the light scintillation to an electronic pulse the packed crystal mounted on a photomultiplier tube by using heavy grease for optical contact with the face of the tube. Typical photomultiplier tubes have a photosensitive layer on the inside of the flat end of the tube. When photons from the scintillation event eject photoelectrons from this layer, the

electrons are attracted to a dynode assembly where multiplication takes place. A voltage divider is arranged so that each successive dynode is at a voltage of 70 to 200 volts. For each photoelectron emitted from the photosensitive surface a burst containing 10^6 or more electrons depending on the voltage, is collected at the anode. The amplitude of the final pulse depends on the initial number of photoelectrons which in turn depends on the number of photons produced by an ionizing event in the crystal [Aj-60].

2.4.2. Semiconductor Counters

The operation of a semiconductor detector is analogous to the operation of an ionization chamber. In an ionization chamber the incident radiation produces positive ions and electrons, and an electrical signal is obtained by collecting these ions. In a semiconductor counter the incident radiation produces electrons and holes; then information about the radiation is obtained by collecting them. One major difference is that only 3.5 eV has to be expended to produce an electron - hole pair in a semiconductor, while the value is approximately 30 eV in an ionization chamber [Ou-75].

To understand the operation of the semiconductor counter, it can be considered as a bar of a uniform crystal with electrodes attached at the both ends as seen in figure (2.9),

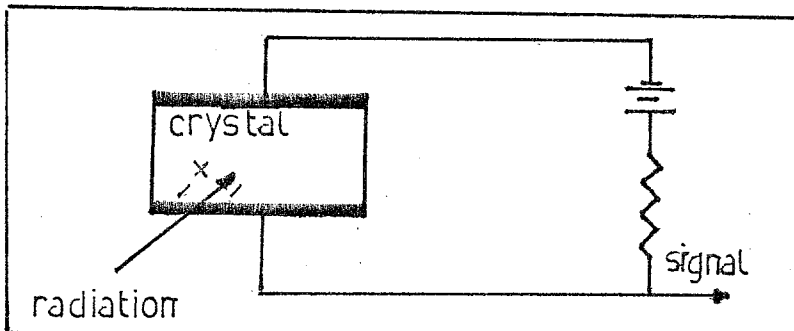


Fig. 2.9. Single crystal semiconductor counter [Ou-75]

When radiation falls on the crystal, electron-hole pairs are produced and when they are collected by the electrodes, a signal is obtained across the resistance, the height of which is proportional to the energy expended by the radiation in the material.

Energy resolution is usually an important factor in selecting a detector material. Energy resolution depends on the number of electron-hole pairs produced by the radiation. The energy required to produce an electron-hole pair is lower in a semiconductor than in an insulator. Hence the resolution of semiconductor detectors will be much better than that of detectors made of insulators.

The electron-hole pairs produced should be free to move in the detector and should be able to reach the electrodes. If the detector contains traps produced by impurities and defects, the electrons may be captured before reaching the electrode. Therefore only single crystals with the least number of traps are used for detectors. Trapping centers enhance recombination of holes and electrons, which will also reduce the pulse height.

2.4.2.1. LITHIUM DRIFTED DETECTORS(*)

Lithium is a donor atom. The diffusion coefficient is about 10^7 higher than that for phosphorus and therefore, deep diffused junctions can be prepared. First, lithium is coated onto single crystals into which it is diffused by heating. In the second step, the sample is heated and a strong reverse

(*) Detailed explanations about these detectors can be obtained from the given references, such as [De-63, Pr-64, Ou-75, So-72].

bias(*) is applied. This helps the lithium atoms to diffuse deeper into the crystal. There are three regions; the p region, the n region and an intrinsic region. Figure(2.10) shows these regions.

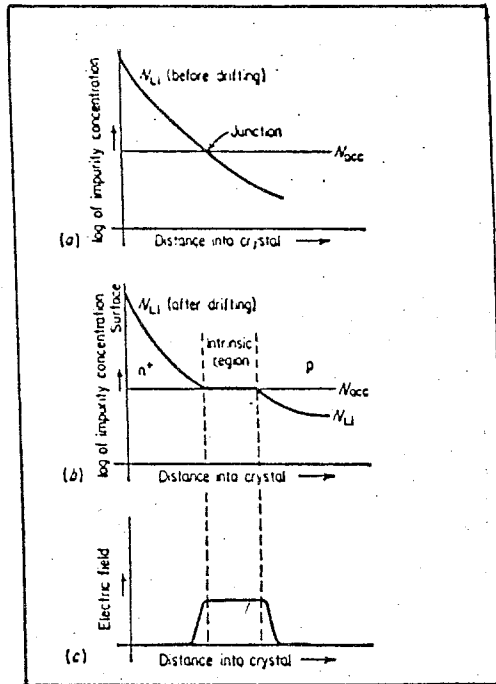
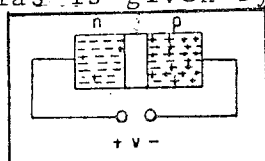


Fig. 2.10. Concentration profile for acceptor impurities (N_{acc}) and for lithium donor impurities (N_{Li}) in a diffused junction diode. (a) Before the drifting process; (b) after the drifting process; (c) electric field distribution with applied reverse bias [Pr-64].

When lithium atoms diffuse onto a surface, the donor - acceptor profile of the junction diode will be as shown in the figure 2.10 (a). If the resulting junction is reverse - biased and the temperature is raised for a short time, the donor-acceptor profile changes to that shown in fig. 2.10 (b). The donor and acceptor impurities in equal numbers compensate each other and effectively produce an intrinsic region bet-

(*) Reverse bias is given by the following figure.



ween the p and the n regions. When the material is returned to room temperature and operated as a detector with a reverse bias, the residual carriers are swept from the intrinsic region and a space charge is developed at either edge, as indicated by the electric field plot in fig. 2.10 (c).

Intrinsic regions over 1 cm thick have been produced by this technique and further developments will lead to thicker sensitive regions. But when one goes to larger and larger sensitive regions, the pulse rise time lengthens.

2.5. GAMMA RAY SPECTROMETRY

The detection of gamma rays is mostly performed by means of a NaI(TL) scintillator or with a Ge(Li) junction. As the output of these detectors is proportional to the amount of energy absorbed, integral counting as well as spectrometry is possible. The NaI(TL) and the Ge(Li) detectors are made in different shapes, according to the purposes for which they will be used such as planar, coaxial etc.

Although gamma rays are monoenergetic, their energy spectra are complex, as illustrated in Fig.(2.11),

According to the size and the absorbing properties of the detector used and depending on the energy of the gamma ray, total or partial absorption can occur.

With NaI(TL) detectors, the escape of the detector X ray by photoelectric effect can be observed at 28 keV. At higher energies the escape peak disappears under the photo-

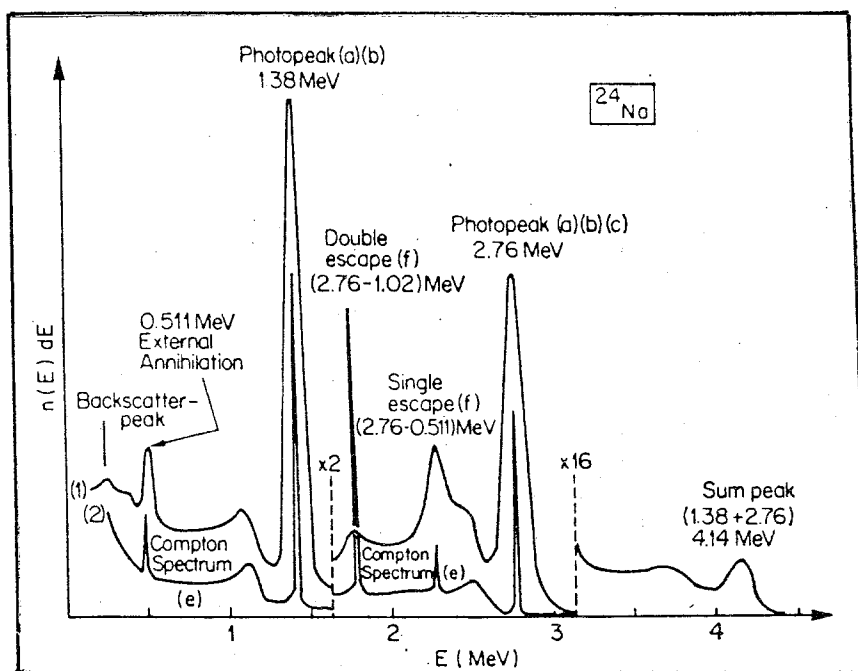


Fig. 2.11. Gamma spectrum of ^{24}Na [S0-72].

peak due to the poor detector resolution. In Ge(Li) detectors the escape peak should appear at about 10 keV below the photopeak, and should be observed up to energies of approximately 1 MeV.

When dealing with gamma rays with energies higher than 1.02 MeV single and double escape of the positron annihilation quanta produce peaks at energies respectively 0.511 and 1.02 MeV below the gamma ray energy. In this case an annihilation peak also occurs at 0.511 MeV due to positron escape and subsequent annihilation in the surrounding materials. When positron emitters are measured, the typical annihilation gamma ray of 0.511 MeV appears in the spectrum and in good geometry also the 1.02 MeV sumpeak.

In order to achieve gamma spectrometric measurements, pulse height discrimination becomes necessary, which can be

done either with a single channel or a multichannel pulse height analyzer.

The resolution of a gamma spectrometer varies with the gamma ray energy and is normally given for the 0.662 MeV peak of ^{137}Cs . With a NaI(TL) detector, the resolution R is expressed in percent using the expression, [SO-72],

$$R\% = \frac{(\text{FWHM}) \times 100}{E} \quad (2.19)$$

where, FWHM is full width at half maximum of the photopeak and E is the peak energy.

When dealing with Ge(Li) detectors, the resolution is normally expressed as the FWHM in keV of the 0.662 MeV ^{137}Cs . Typical resolution of these detectors are between 5 and 2 keV.

The detector resolution R_D is energy dependent, and can be expressed as follows:

$$R_D^2 = 2.355 F E \epsilon \quad (2.20)$$

where, R_D and the gamma energy E are given in keV, ϵ is the energy needed to create an electron-hole pair and is approximately equal to 3.10^{-3} keV for Ge, for Ge(Li) detectors, F is the Fano factor which is variance to yield factor (The true value of fano factor for Si and Ge is still unknown. However it is assumed as 0.1 for both Si and Ge [Or-80]).

Then the resolution of the spectrometer R is composed of the electronic line broadening, for which the main factor is noise produced by the detector and the amplifying instruments, and of the detector resolution as,

$$R^2 = R_E^2 + R_D^2 \quad (2.21)$$

The energy of the unknown gamma ray can be determined by the pulse height of the photopeak maximum. The calibration of the spectrometer is performed by means of isotopes emitting gamma rays of known energies. Sets of calibration sources are commercially available or can be prepared by irradiation of the stable isotopes. As a Ge(Li) detector shows a very good resolution, gamma ray energies can be determined within a few tenths of a keV. Table 2.2 shows some calibration sources for gamma-ray spectrometry.

Table 2.2. Calibration sources for gamma-ray spectrometry
[S0-72]

Isotope	Energy (keV)	Isotope	Energy (keV)
²⁴¹ Am	59.568 ± 0.017	⁹⁵ Nb	765.83 ± 0.07
¹³¹ I	80.166 ± 0.009	⁵⁴ Mn	834.84 ± 0.07
¹⁵³ Gd	97.43 ± 0.02	⁸⁸ Y	898.01 ± 0.07
¹⁶⁰ Gd	103.18 ± 0.02	²⁰⁷ Pb	1063.82 ± 0.28
¹⁷⁷ Lu	112.97	⁶⁰ Co	1173.13 ± 0.04
¹⁴¹ Ce	145.44 ± 0.05	⁶⁰ Co	1332.39 ± 0.05
¹³⁹ Ce	165.84 ± 0.03	²⁴ Na	1368.40 ± 0.04
		⁶⁰ Co(D.E.)	1576.9 ± 0.32
¹⁷⁷ Lu	208.36	²⁰⁸ Tl(ThC'')(D.E.)	1592.3 ± 0.13
²⁰³ Hg	279.12 ± 0.05	²⁴ Na (D.E.)	1731.6 ± 0.16
¹³¹ I	364.47 ± 0.005	⁸⁸ Y	1836.1 ± 0.07
		⁶⁰ Co	2035.2 ± 0.50
		⁶⁰ Co(D.E.)	2232.2 ± 0.87
¹⁹⁸ Au	411.776 ± 0.01	²⁴ Na (S.E.)	2242.6 ± 0.14
		⁶⁰ Co	2598.9 ± 0.30
Annihilation	511.006 ± 0.02	²⁰⁸ Tl (ThC'')	2614.3 ± 0.09
²⁰⁷ Pb	569.65 ± 0.10	²⁴ Na	2753.6 ± 0.12
¹³⁷ Cs	661.59 ± 0.07	⁶⁰ Co	3202.4 ± 0.65
		⁶⁰ Co	3254.2 ± 0.65
		⁶⁰ Co	3273.6 ± 0.4
		⁶⁰ Co	3452.5 ± 0.75

S.E. = single escape D.E. = double escape.

Quantitative data of the gamma ray intensity can be obtained from the area or the height of the photopeak. The photopeak can be described as a Gaussian error curve, with an area S, given by

$$S = 1.06 (\text{FWHM})h \quad (2.22)$$

where, h is the peak height.

When peaks are superimposed on Compton continuum of higher energy or when poor counting statistics are obtained, it is difficult to determine the peak height. There are some other methods to calculate the peak areas, such as the method of Covell (which will be examined in chapter 3 in detail), spectrum stripping, etc.

Linearity of a detector is another important subject. This property of semiconductor detector is excellent, and deviations from linearity in existing spectrometers can be traced back to amplification and analysis equipment except for very low energies, [HO-71].

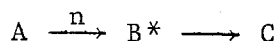
2.6. QUANTITATIVE ANALYSIS

2.6.1. Basic Principles and Concepts

The technique of activation analysis is based on the formation of radioactive nuclides as a result of reactions between nuclear particles and the isotopes of the trace elements of interest. These isotopes are transformed into different isotopes of either the same or a different element. There are many nuclear reactions for this purpose. In the great majority of cases two-particle reactions are utilized; one particle being a reactant and the other being a product. The particles for general use as reactants are neutrons, protons, deuterons, alpha particles and photons.

As an example, when a target material is exposed to

neutrons, photons or charged particles, the resulting nuclear reaction and the product decay may be represented by,



The stable isotope A being bombarded is transmuted to radionuclide B* which decays with its characteristic decay constant λ_B to the stable isotope C. (or unstable isotope C* which will also decay with λ_C , etc.). These subjects are examined in 2.1.1 and 2.1.2 in detail.

The rate of formation, R_F , of a particular activation product B in a given sample is proportional to the intensity of the flux or beam of incident particles, to the concentration of the target nuclide in the sample, and to the cross-section for the nuclear reaction. Thus, for the case of neutron irradiation R_F is given as,

$$R_F = \phi N_A \sigma V \quad (2.23)$$

where $N_A = Avf\rho/M_A$ or for unit volume, $N_A = A_vmf/M_A$ and R_F is assumed as constant (the case of $R_F \neq$ constant, was examined in subchapter 2.1.2). Then,

$$R_F = \frac{\phi m A_v f \sigma}{M_A} \quad (2.24)$$

where,

N_A = the number of target atoms

ϕ = the neutron flux (n/cm²-sec)

m = the mass of the trace element in the specimen
(gm)

M_A = the atomic weight of the trace element.

f = the fractional isotope abundance of the target nuclide

A_v = Avagadro's number (atom/gm. atom)

σ = the reaction cross-section (cm^2/atom)

The decay rate expressed in atomic disintegrations per second of the product radionuclide in the specimen was given by the equation 2.3, as

$$dN/dt = -\lambda N$$

Therefore, the rate of change of the quantity of the activation product in the sample during the irradiation is given as,

$$\frac{dN_B}{dt} = R_F - N_B \lambda_B \quad \text{or,}$$

$$\frac{dN_B}{dt} = \frac{\phi m A_v f \sigma}{M_A} - N_B \lambda_B \quad (2.25)$$

The solution of this differential equation will be given in the appendix B. The result is,

$$N_B = \frac{\phi m A_v f \sigma}{\lambda_B M_A} (1 - e^{-\lambda_B t}) \quad (2.26)$$

The number of disintegrations per second of the product represented as the activity after an irradiation time t sec is,

$$A = N_B \lambda_B = \frac{\phi m A_v f \sigma}{M_A} (1 - e^{-\lambda_B t}) \quad (2.27)$$

The equation (2.27) can be used to estimate the activation levels for the various elements in the sample under varying conditions of flux, irradiation time and sample size. This equation is rearranged for M and is given as, [KO-60],

$$M(\text{gm}) = \frac{M_A A}{\phi A_v f \sigma (1 - e^{-\lambda_B t})} \quad (2.28)$$

$$M(\text{gm}) = \frac{M_A A}{6.02 \times 10^{23} \phi f \sigma (1 - e^{-\lambda_B t})} \quad (2.29)$$

The maximum or saturation activity is attained when the irradiation time is long compared to the half-life of the product and the rate of formation is equal to the rate of disintegration $[1 - e^{-\lambda_B t}]$ approaches unity and,

$$A_s = N_B \lambda_B = \frac{\phi M A_v f \sigma}{M_A} \quad (2.30)$$

Assuming the sample has been irradiated to saturation and no significant decay has occurred in the interval between the end of bombardment and counting time t, a formula can be obtained to calculate the weight of the element by reducing the eg. (2.30) to,

$$M(\text{gm}) = \frac{M_A A}{6.02 \times 10^{23} \phi f \sigma} \quad (2.31)$$

Besides, the sensitivity expressed in terms of the minimum weight of a target element that will produce the minimum detectable activity after a certain bombarding time t will be given by,

$$\text{Sensitivity in micrograms} = \frac{M_A A \times 10^6}{6.02 \times 10^{23} \phi f \sigma (1 - e^{-\lambda B t})} \quad (2.32)$$

2.6.2. Methods of Analysis

When determining a large number of elements in one sample, the comparison of samples is very troublesome and time consuming. Besides, it is often impossible to irradiate a large number of comparison samples together with the unknown samples because of the restricted volumes of irradiation containers. Moreover, when performing multielement determinations, the presence of one or several elements in the sample is often only realized during the evaluation of the gamma spectra. When no comparison standards for these elements were irradiated a semiquantitative guess is all that can be derived concerning the concentrations.

Several techniques have been put forward to eliminate the above mentioned shortcomings, such as absolute, single or multicomparator methods [H0-71].

Absolute Method: This direct method performed by Girardi, [Gi-64], based on the determination of weights by means of the activation relationship is generally considered as a semiquantitative technique. It deals with absolute gamma-ray counting and direct calculation of weights from an estimation of the activation formula. The reliability of that method depends on the accuracy with which the absolute of the thermal cross-section, the resonance integral and the decay schemes are known. But most of these are not sufficiently well known to allow their use for accurate absolute calculations.

Comparator Method: Many sources of error can be eliminated by irradiating the sample with a standard of similar composition. When the composition of a substance is unknown, a preliminary irradiation can be used as a qualitative determination, and a suitable standard can be fabricated for simultaneous irradiation with the sample. The sample and standard should be approximately the same weight, shape and thickness. Thus, effects due to variations in flux and self attenuation during the irradiation will not influence the analysis.

If the activities are measured in an identical manner, comparing the counting rates in the sample and the standard will give the weight of the element in the sample by the following simple relation, [Ad-70],

$$\begin{array}{l} \text{Weight of the} \\ \text{element in} \\ \text{sample} \end{array} = \frac{(\text{Activity in sample}) (\text{Weight of element in standard})}{(\text{Activity in standard})}$$

Single Comparator(*): The use of Ge(Li) detectors allows the simultaneous determination of an important number of elements after a single irradiation. The classical comparator method requires one standard for every element to be determined; but this method uses only one comparator consisting of known amounts of elements to be determined.

2.7. USE OF COMPUTERS IN NAA

With the development of modern methods of radionuclide measurement, the ability to record vast amounts of data reliably and automatically has led to the use of computers. Many computer programs have been developed up to now, to interpret and evaluate that data. Some of these programs have been pre-

(*) This method will be discussed in appendix C in detail.

pared for big computers; some others are for small ones, which are particularly for the routine processing of spectrometric data.

The large computers allow, of course, the most sophisticated and complete analysis of that data. The computing programs can easily be written and the hardware problems for data collection and transfer to the large computer system have been solved for several years. This data analysis system keeps its full validity, particularly when large batches of samples must be analysed in a similar way [Be-73].

On the other hand, the small computers are used at the laboratory site for both data acquisition and analysis, requiring the small amounts of memory locations. However, the programs prepared for these computers could not define the location of photopeaks and the computation of the peak areas as reliably as those prepared for the big ones [Ci-81].

As a result of this short discussion, it can be concluded that; the choice of a data-handling system is a rather cumbersome task since the money involved is important and the information available on the real performances of the various systems are often incomplete.

Modern gamma-spectrometers collect their data simultaneously in a large number of energy channels. When NaI(TL) scintillators are used, a few hundred energy channels were sufficient for adequate representation of the spectrum. But, the introduction of the Ge(Li) type detector with its excellent energy resolution, pushed the number of energy channels required for an adequate spectral response, up into the thousands. For the relevant information from the enormous amount of digital data comprised in a Ge(Li) type spectrum,

"data-reduction system" has been used. The purpose of such a system is to recognize valuable data and to reject useless information.

The following five steps can be distinguished in the course of a data-reduction procedure, as in Gretel which will be discussed in chapter 3, [H0-71];

- a) the preliminary treatment of all data
- b) the location of possible photopeaks
- c) the determination of the peak limits
- d) the application of statistical and peak shape tests,
- e) the evaluation of exact peak centroid location,
determination of net peak area with corresponding
standard deviation and detection limit.

a) In general, the preliminary treatment of data involves the application of smoothing techniques and the computation of the first or higher derivatives of the spectral data. Smoothing, as will be discussed in 3.3.2, is used to minimize the influence of the statistical counting fluctuations. First or second derivatives which can be calculated according to the least squares method, detect the location of the photopeaks [So-72, Sa-64].

b) For the detection of exact photopeak location, the calculation of the maxima that belong to true photopeaks is the second problem. When a preliminary data treatment has produced a first derivative, then the maxima correspond to points where this function goes from a positive value to a negative one.

c) The most critical point in the data-reduction system is the determination of the left and right boundaries of the

photopeaks. This will determine the separation that is to be made between the peak area and underlying continuum. This will also determine both accuracy and precision of any quantitative interpretation of the spectra.

d) Once a maximum has been detected and its boundaries have been established, some further tests are required to make sure that a real photopeak has been found. There are different tests used for that purpose. Some of them used in Gretel will be studied in 3.3.3.

e) Once a peak has been detected and accepted, the spectral data contributing to the photopeak should be interpreted in terms of energy and intensity of the corresponding gamma radiation. This is the classical problem of NAA and it is usually the last step of many computer programs.

Besides these, another important subject of some programs is the analysis of detection limits for the peaks actually found or for those not detected but expected at specified location.

Up to now some analysts have used minimum data-reduction without any qualitative interpretation. Others used a very extensive data reduction but only a limited amount of qualitative analysis, mainly to interpret the standard spectra in order to define the elements to be determined [Ho-71].

The program examined in this thesis is one of these programs which involves data-reduction, interpretation of the activation process and standard spectra quantitatively, which will be examined in the following chapter in detail.

3. GRETTEL

A COMPUTER PROGRAM FOR GAMMA RAY SPECTROMETRY WITH Ge(Li) DETECTORS

3.1. INTRODUCTION OF THE COMPUTER CODE

Gretel performs the quantitative analysis of gamma-ray spectra obtained by Ge(Li) detectors, using special "oriented libraries" which are prepared for each particular problem (Table 3.1 shows such an oriented library).

The scheme of the Gretel program is the following: An energy calibration source is counted as the first spectrum of the series; then the unknown neutron activated samples are successively counted. The counting equipment consists of a 4096 channel analyzer for high resolution spectrometry with Ge(Li) detectors, equipped with sample changers to allow the automated acquisition of gamma-ray data from successive specimen. The data coming from the analyzer, are first punched on paper tapes, then transferred on to magnetic tapes. These tapes are then read by one of the subroutines (TAPE, TAPEAS or TAPEGA), according to the punching code adopted; such as IBM, ASCII or ASCII with no parity. These routines transform the data in a suitable format, either by the functions CONV or CONVB.

The analysis is carried out by the single comparator method as will be seen in the appendix C.

TABLE 3.1- EXAMPLE OF "ORIENTED LIBRARY" FOR LONG LIVED RADIOISOTOPES

[Gu-74]

Radioisotope	Energy (keV)	Resolution (keV) FWHM	Decay Constant (minutes)	Specific Activity ⁽¹⁾ (cpm)	Molecular Weight
Ag 110 M	657.8	3.0	1.85×10^{-6}	6.911×10^8	107.87
Ca 47	489.5	2.9	1.01×10^{-4}	5.761×10^{13}	40.08
Ca 47	807.4	3.1	1.01×10^{-4}	9.274×10^{13}	40.08
Ca 47	1296.4	4.2	1.01×10^{-4}	1.13×10^{14}	40.08
Co 60	1173.2	4.1	2.497×10^{-7}	2.612×10^7	58.9
Co 60	1332.4	4.2	2.497×10^{-7}	2.893×10^7	58.9
Cs 134	604.7	3.0	5.98×10^{-7}	4.1×10^7	132.91
Cs 134	795.0	3.1	5.98×10^{-7}	5.683×10^7	132.91
Hg 203	279.1	2.5	1.035×10^{-5}	1.018×10^9	200.61

(1) Calculated from irradiation of one μg of element in a thermal flux of 10^{13} n/cm².sec at saturation and at decay time equal zero.

The program is essentially composed of a MAIN part which controls the various subroutines, also involving those mentioned above, thirteen subroutines and three functions.

These routines which detect and evaluate peak areas perform the following operations:

- Local smoothing of the spectrum for minimizing the influence of the statistical counting fluctuations.
- First derivative of the smoothed spectrum.
- Detection of peak location according to the change of sign of the first derivative.
- Computation of the net area of each peak found.
- Possibility of detection and computation of double peak.

f) Computation of the minimum counting rates that could have had a photopeak to be detected over the existing background in case no peaks are found in the interval indicated by the oriented library [Gu-74, Gu-67].

All those subjects will be discussed in the following subchapters in more detailed form. Description of the routines will also be given in appendix G.

Results in ppm or sensitivity limits for the elements not detected, are obtained in the form of digital lists.

At first, the original raw spectrometric data are collected on punched tapes which are then processed at the Scientific Information Processing Center, Ispra-Italy, by an IBM 370/165 computer with fortran IV level G and assembly languages. It also has possibility of obtaining analog outputs as results given in digital lists, using a CALCOMP unit. One request, a graphical representation of the spectrum could be produced, showing the way in which the spectrum has been processed. The code has been updated by the author for UNIVAC-1106, by deleting all the plotting subroutines from the main program because a Univac plotting unit was not available in BU. The updated form of the code is recorded on a new magnetic tape for further studies by BU-NE Department. This tape(*) contains four files:

1. file: UNIVAC-1106 version of GRETEL source program containing main and subprograms.
2. file: sample problem input data
3. file: spectra for sample problem
4. file: sample problem output data

(*) This tape covering the UNIVAC version of Gretel and input and output data is labelled as "SAGLAM" for further studies by the NE Department.

3.2. THEORY OF NAA USED IN GRETEL

In this chapter, the theory of neutron activation analysis, NAA, used in Gretel is given in a summary form as follows:

The neutron activation analysis in subchapter 1.3 and the basic principles of this method in 2.6 were examined in detail. The equation (2.27) derived in the same subchapter 2.6 was given

$$A = N_B \lambda_B = \frac{\Phi m A_v f \sigma}{M_A} (1 - e^{-\lambda_B t})$$

where, A represents the amount of radioactivity of the daughter nuclei N_B which is present at the end of the activation time t. Since the activity also decays exponentially with time, according to the equation (2.5), the activity A_d that will be detected, will depend on the decay time τ after irradiation,

$$A_d = A e^{-\lambda_B \tau} = \frac{\Phi m A_v f \sigma}{M_A} (1 - e^{-\lambda_B t}) e^{-\lambda_B \tau} \quad (3.1)$$

This equation can be rearranged for the weight of the element as,

$$g(\mu\text{gm}) = \frac{M_A A_d \times 10^6}{6.02 \times 10^{23} \Phi f \sigma (1 - e^{-\lambda_B t}) e^{-\lambda_B \tau}} \quad (3.2)$$

Besides the basic activation formulas given above the analysis of the spectra with computers will be discussed in the following chapters in the detailed form.

3.3. ANALYSIS OF GAMMA-RAY SPECTRA WITH DIGITAL COMPUTERS

3.3.1. Introduction

With the development of modern methods of radionuclide measurement, as discussed in subchapter 2.7, the ability to record vast amounts of data automatically has led to the use of computers. Many computer programs have been derived up to now, to interpret that data.

Gretel is one of these programs prepared for gamma-ray spectra obtained with high resolution Ge(Li) detectors. The program is used for the quantitative analysis of trace elements from gamma-spectra of neutron activated materials. As mentioned before, it is composed of three main parts as follows:

First part introduces the smoothing procedure by taking 5, 7 or 9 point segments of the spectrum. This procedure is used to minimize the influence of the statistical counting fluctuations. The smoothed spectrum is then examined to define the location of a peak where the sign of the derivative changes from a positive value to a negative one. This subject will be discussed in the subchapter 3.3.2 in more detail.

In the second part, the peaks are searched and located exactly. To avoid the statistical fluctuations from being kept as analytical peaks, the absolute values of the maximum and minimum of the derivative must be larger than a confidence band chosen suitably. The criteria for the choice of confidence band will be explained in the subchapter 3.3.3. Besides this procedure, some other tests for recognizing the peaks as single or double will also be introduced in the same subchapter.

Third part examines the computation methods of peak

areas. There are two different methods used in the program.

1) Computation of the areas of the peaks observed over the background: If the peaks found are above the Compton continuum, their left and right limits are examined according to Cove1, [Co-59], in the subsection 3.3.4.1.

2) Searching of detection limits: As mentioned in the 3.1 briefly, in case the peak is not detected in the interval given by the library, the minimum counting rates due to the hidden peaks are calculated by a particular subroutine, subroutine SENSIT. That procedure will be given in the 3.3.4.2 in more detail.

Finally, the results in ppm or sensitivity limits for the elements not detected are obtained in the form of a digital list. The errors due to statistical counting fluctuations, as percentages, are also estimated in the program. That data is also given in the appendix K.

3.3.2. Smoothing and Differentiating Procedures

The majority of published literature on the uses of computers in activation analysis relates to the unscrambling of gamma-ray spectra. This may be achieved in a variety of ways, such as spectrum stripping, the method of least squares fitting, vector analysis and iterative methods, etc [Gi-65]. One of these ways is the spectrum smoothing procedure.

The primary output of any experiment is information which measures the phenomenon under observation. The random errors which are indistinguishable from this information are

characteristically described as noise. It is important for an analyst, to remove as much of this noise as possible without unduly degrading the underlying information. There are some mathematical methods for handling such problems in the processing of analytical data.

One of the simplest ways to smooth the fluctuating data is by a moving average based on the concept of a convolute and of a convolution function(*) [Sa-64].

In summarized form, it can be said that the method of data convolution is to fit a polynomial to such data by use of the method of least squares. Assuming that f 's represent any set of convoluting integers, and S_n is the smoothed value in the n th channel, the mathematical description of this process is given by [Sa-64, Gu-74],

$$S_n = \frac{1}{N} \sum_{i=-m}^{+m} (f_i S_{n+i}) \quad (3.3)$$

where the index n indicates that $(2n+1)$ points are used for that procedure. For the calculation of these convolution coefficients, a set of experimental points is to be fitted to a curve, as given in the eq. (3.3), and then the square of the differences between the computed values and the observed values are to be minimum for the total of the observations.

There is also an associated normalizing or scaling factor. These coefficients and factors are given in the tabular form [Sa-64]. The table 3.2 gives the convoluting integers and normalizing factors.

(*) Convolution method will be discussed in the appendix D.

TABLE 3.2 [Sa-64]

CONVOLUTES	SMOOTHING		QUADRATIC		CUBIC		A20	A30			
POINTS	25	23	21	19	17	15	13	11	9	7	5
-12	-253										
-11	-138	-42									
-10	-33	-21	-171								
-09	62	-2	-76	-136							
-08	147	15	9	-51	-21						
-07	222	30	84	24	-6	-78					
-06	287	43	149	89	7	-13	-11				
-05	322	54	204	144	18	42	0	-36			
-04	387	63	249	189	27	87	9	9	-21		
-03	422	70	284	224	34	122	16	44	14	-2	
-02	447	75	309	249	39	147	21	69	39	3	-3
-01	462	78	324	264	42	162	24	84	54	6	12
00	467	79	329	269	43	167	25	89	59	7	17
01	462	78	324	264	42	162	24	84	54	6	12
02	447	75	309	249	39	147	21	69	39	3	-3
03	422	70	284	224	34	122	16	44	14	-2	
04	387	63	249	189	27	87	9	9	-21		
05	322	54	204	144	18	42	0	-36			
06	287	43	149	89	7	-13	-11				
07	222	30	84	24	-6	-78					
08	147	15	9	-51	-21						
09	62	-2	-76	-136							
10	-33	-21	-171								
11	-138	-42									
12	-253										
NORM	5175	8059	3059	2261	323	1105	143	429	231	21	35

Increase in the number of data points used causes better smoothing but decreases the peak amplitudes at the same time. The minimum distortion occurs when the polynomial accurately describes the analytical data. The best results are obtained when the data are digitized at high densities, i.e., points very close together and the number of points used in the convolution is chosen to be small enough so that no more than one inflection in the observed data is included in any convolution interval. As an example, equation (3.4) can be illustrated for 5 points as, [Gi-65, Gu-74].

$$S_n = \frac{1}{35} \left[-3S_{n-2} + 12S_{n-1} + 17S_n + 12S_{n+1} - 3S_{n+2} \right] \quad (3.4)$$

Another important subject in the use of this method is the degree of polynomial. High degree polynomials represent less smoothing than the others. This is especially important when there are double peaks found. Therefore, it is shown that, quadratic and cubic functions are the best for that convolution process. [Ci-81].

This type of data processing, as far as computers are concerned requires a relatively small amount of programming and relatively little use of the computer memory.

As mentioned in the previous subchapters, the first selection of possible peaks is done by considering the change of sign of the first derivative from a positive value to a negative value. The zero point can be considered as the location of a peak.

Therefore after smoothing procedure, this smoothed spectrum is differentiated by the following formula [Gu-74],

$$d_n = \frac{S_{n-1} + S_{n+1}}{2} \quad (3.5)$$

where d_n is the differentiated value of the content of the n th channel. S_n 's represent the smoothed values in the channels $(n-1)$ and $(n+1)$.

These smoothing and derivation processes are carried out by the subroutines SMOOS and DERIV which are given in the appendices G, I.

3.3.3. Search of the Peaks

After the smoothing procedure, the second problem con-

sidered by Gretel is the search of peak locations. As mentioned in the previous subchapter, a first selection of possible peaks is made by considering the change of sign of the first derivative from a positive value to a negative value. Therefore, one needs the values of maxima and minima of the spectrum. Many authors use different methods, [Ho-71], to obtain these values with or without preliminary data treatment as explained in the subchapter 2.7. The most simplified system of all is used by Anders [An-69], by just taking any set of consecutive data points that are higher than their neighboring points as a possible peak.

A similar method is used in Gretel which can be explained briefly as follows:

A possible maximum in channel n , is given by the condition below [Gu-67],

$$C_{n-2} < C_n - P \sqrt{C_n} > C_{n+2} \quad (3.6)$$

when its content C_n satisfies that condition. Where P is chosen empirically (presently $P = 1$).

The minima on both sides of the maxima are chosen as the channels $(n-k)$ and $(n+k')$ for which:

$$C_{n-k-1} \geq C_{n-k} - P \sqrt{C_{n-k}} \quad (3.7)$$

and,

$$C_{n+k'+1} \geq C_{n+k'} - P \sqrt{C_{n+k'}} \quad (3.8)$$

The point n is considered as the location of a peak.

For the determination of the exact photopeak locations, besides the conditions stated, two conditions must be satisfied simultaneously:

1. The absolute values of the maximum and the minimum of the derivative must be larger than a confidence band suitably chosen. That band is calculated by multiplying standard deviations of data in each channel with a factor chosen, [Gu-74]. Data is taken from the smoothed spectrum instead of the original one. As a result of that process, the counting fluctuations are reduced and statistical fluctuations are kept as analytical peaks are avoided.

2. The ratio between maximum and minimum values of the spectrum must be less than or equal to 1.5. This also avoids unsymmetrical peaks being considered.

After having made these tests, the center of a peak where the sign of the first derivative changes from a positive value to a negative value, is determined by linear interpolation by taking the negative and positive differentiated data in the nearest channels to the center.

In the cases in which these two conditions are not fulfilled, a further examination of the shape of the derivative is performed to ascertain if a double peak is present in the considered interval. For two successive peaks to be considered as a doublet, the distance between the summit of each peak should be less than $(2.0 \times \text{FWHM})$ [Gu-67, Gu-74].

According to the tests performed on the ratios between the maxima (AMAX) and the minima (AMIN) of the derivatives, three different cases are possible according to Guzzi, [Gu-74]. These are shown in the figure(3.1).

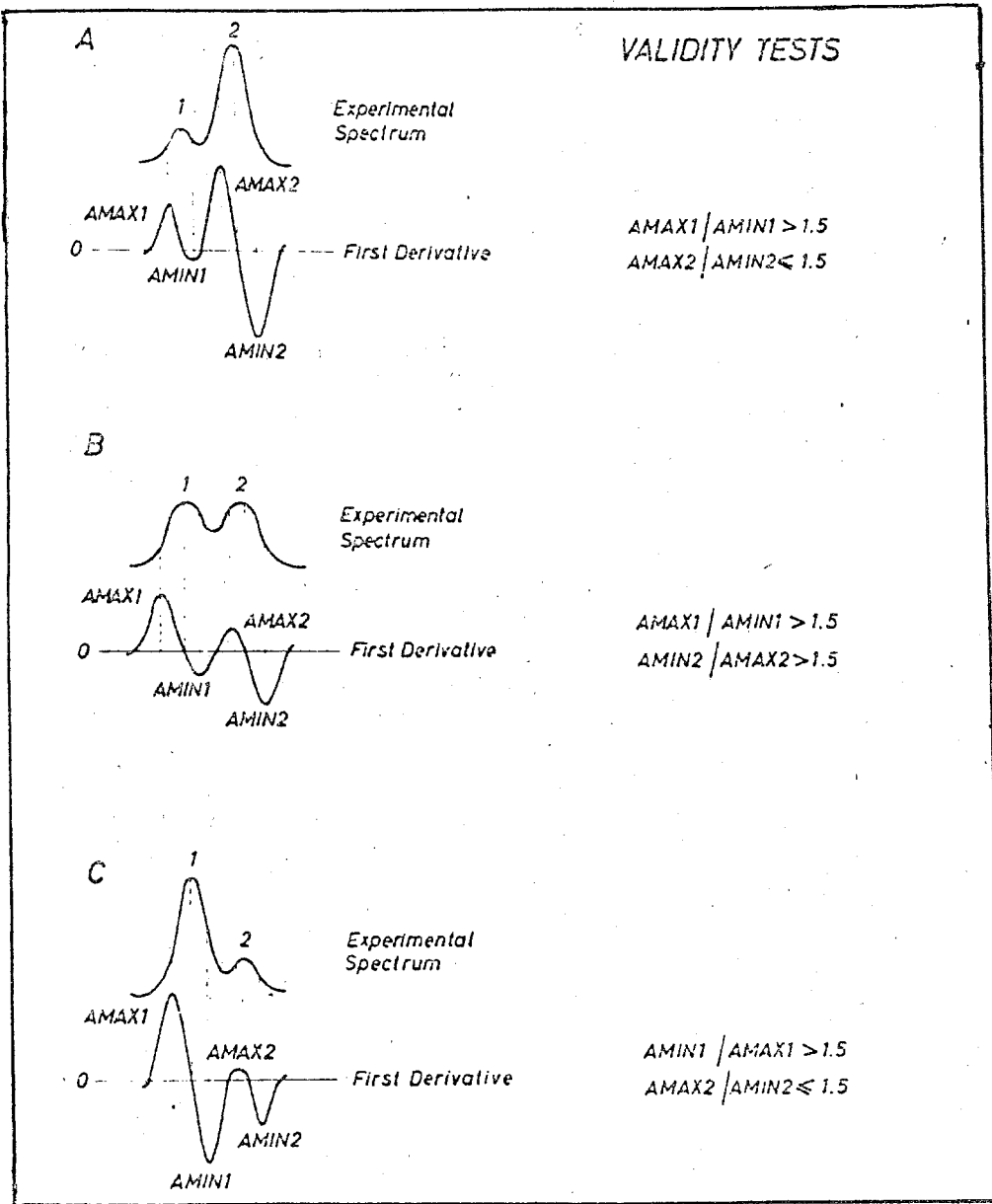


Figure 3.1

where,

A. $AMAX1 / AMIN1 > 1.5$ $AMAX2 / AMIN2 \leq 1.5$

Peak 1 is considered as belonging to a doublet, peak 2 is treated as a single peak.

B. $AMAX1 / AMIN1 > 1.5$ $AMIN2 / AMAX2 > 1.5$

Both 1 and 2 are treated as a doublet

C. $AMIN1 / AMAX1 > 1.5$ $AMAX2 / AMIN2 \leq 1.5$

Peak 1 is considered as single peak, peak 2 is treated as being a doublet.

On the other hand, the other most critical point in searching of peak locations is the determination of the left and right boundaries of the photopeaks. This will indeed determine the separation that is to be made between the peak area and underlying continuum. Many authors follow different ways for the determination of the peak boundaries [Ho-71]. That used in Gretel can be given in a brief form as follows:

The below figure, fig.(3.2), shows a small photopeak which is assumed as Gaussian, superimposed on a high constant

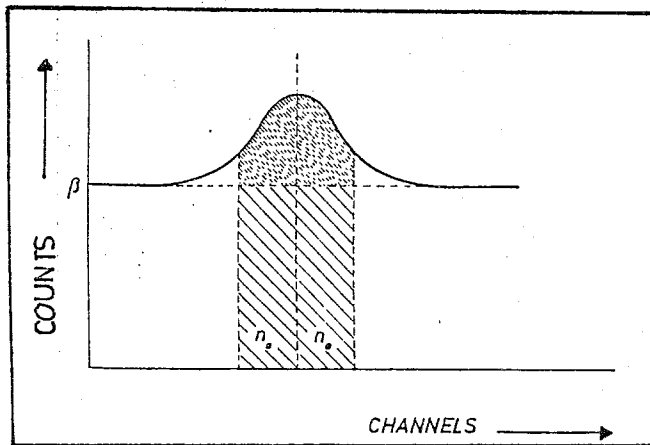


Fig.3.2. Schematic representation of a photopeak [Pa-66].

background such as it would be obtained after a very long measurement. The area of the background under the photopeak as seen in the above figure is given by J.Pauly et al. [Pa-66], by the relation, $(2n_0 = 1.2 \times \text{FWHM})$. Where FWHM is the full width at half maximum. If the background counting rate is small, as often happens when a lithium drifted germanium detector is used, the optimum value of n_0 increases slightly.

But approximate value is taken as $2n_0 = 1.2 \times \text{FWHM}$ for all values of the background [Pa-66, Cu-74].

Under this consideration the left and right limits of the peaks are determined in the program by the relations given below;

$$H_1 = n - 1.2 \times \text{FWHM} \quad (3.9)$$

$$H_2 = n + 1.2 \times \text{FWHM} \quad (3.10)$$

where, H_1 and H_2 represent left and right boundaries of a peak. FWHM is the energy resolution.

All these procedures described above are carried out by the subroutines ANALYS and PKF.

3.3.4. Computation of Peak Areas

One of the classical problems of qualitative and quantitative interpretation of gamma spectra for NAA is the calculation of peak areas found by some tests described before.

The procedure employed by Gretel can be explained in two subsections as follows:

1. Computation of the areas of the peaks observed over the existing background. These may be single or double peaks.

2. Computation of the minimum counting rates in case no peaks are found in the given interval.

3.3.4.1. Computation of the Areas of Observed Peaks

a) Computation of the area S of a single photopeak:

There are several methods, in the literature, used for the separation of a peak from its underlying continuum. Seven different methods, [Co-59, Yu-68, Qu-69], were especially studied up to now. Five of them use a straight line to separate the peak from the underlying continuum. Two use a nonlinear procedure, [Ho-71]. One of them which is known as Covell's method, [Co-59], is more common than the others. The method used in Gretel is derived from that method and it can be summarized as follows:

This method describes a procedure which utilizes the digital nature of pulse height distribution data to apply statistical methods of data reduction, [Co-59]. As the response of a pulse height analyzer is stated in counts per unit increment of amplitude, individual channel responses may be represented graphically as rectangles whose areas bear direct proportionality to the number of counts observed in the channel. A graphical representation is given in the figure (3.3).

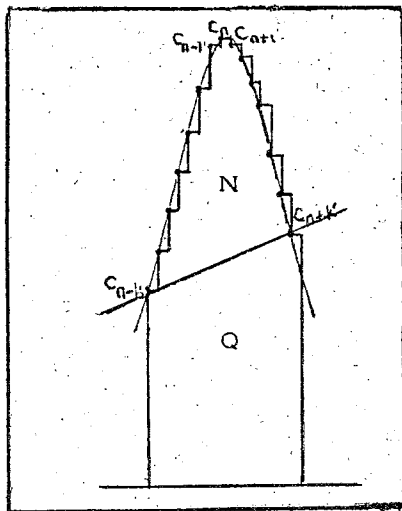


Fig.3.3. Fraction, N , of total counts contained in peak channel, a_2 and n channels on either side.

Here, the response of the channel containing the greatest number of counts is defined as C_n and succeeding channel responses, progressing down the low amplitude side of the peak as $C_{n-1} \dots C_{n-k}$, similarly channel responses on the high amplitude side of the peak are defined as $C_{n+1} \dots C_{n+k}$, (where C_{n-k} and C_{n+k} are defined by the intercept points of the rising and falling traces of the curve).

If the contributions from the C_n channel and the channels on either side are added, the sum is a value representing the total counts contained in these channels. This value is graphically represented as the area P in the above figure. If this area P, were divided by a line connecting the ordinate values of C_{n-k} , C_{n+k} as seen in the figure (3.3), the area above this line can be defined as follows:

$$N = P - Q \quad (3.11)$$

where, N may be used under a wide variety of observational conditions to estimate the total peak area.

Although this method due to Covell, [Co-59], is the most common one, it is especially used for NaI(TL) detector peak area computations. For the Ge(Li) detector peak area computations, many authors use different methods, [Ho-71].

The method used in Gretel is the one which is based on smoothing and differentiating procedures of the spectra. Yule, [Yu-68], gives the principles of that method which is derived from the Covell's method. This method can be explained briefly as follows:

The method known as the smoothed first derivative method employs the simplified least squares data convolution

technique(*) to generate, from the observed spectrum, a new spectrum which has essentially all the statistical scatter removed. The new spectrum defines a smooth curve through the observed data. The data convolution techniques can also generate a spectrum which is the smoothed first derivative of the observed data. The total peak area is then computed from the equation given by P.H. Yule as, [Yu-68]

$$\text{Total peak area} = \sum_{n=n-k}^{n+k'} C_n \quad (3.12)$$

where, C_n is the number of counts in channel n , and the boundary channels are $(n-k)$, $(n+k')$. The base area is the area of the trapezoid below the peak, similar to the figure 3.3.

$$\text{Base Area} = \frac{C_{n-k} + C_{n+k'}}{2} (k-k'-1) \quad (3.13)$$

and the net peak area is given by the difference between them, [Yu-68, Gu-74]

$$\text{Net peak area} = \text{Total Peak Area} - \text{Base Area}$$

or,

$$S = \sum_{n=n-k}^{n+k'} C_n - \frac{C_{n-k} + C_{n+k'}}{2} (k-k'-1) \quad (3.14)$$

where, the quantity $(k-k'-1)$ is the number of channels included in the peak or, in other words, it represents the width of the peak.

(*) This subject will be discussed in the appendix D as mentioned in subchapter 3.2.

The associated errors in the evaluation of the peak areas are also calculated and given as digital lists.

b) Computation of the area of a double peak:

For the computation of the area of a double peak several assumptions are made by the authors. One that is used in Gretel is given as: The area of a doublet is approximated by considering the left half of the first peak and the right half of the second peak as being symmetrical with respect to the remaining portions of each peak. Then these halves are doubled to obtain the total area. The underlying continuum is subtracted as explained in the previous subchapter. Standard

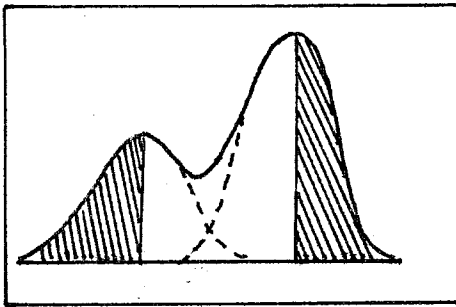


Fig.3.4. Computation of the area of a doublet.

deviation is approximated by doubling the standard deviation calculated for each half of each peak of the doublet.

All these operations are performed by the routines SURF and SURF2.

3.3.4.2. Search of Detection Limits

In the case where no peaks are found by the tests described in the subsection 3.3.4.1, in the indicated interval, the minimum counting rates that could have had a photo-

peak over the existing background, are also calculated in Gretel. This procedure is performed by the subroutine SENSIT. Theory of that calculation can be explained briefly as follows:

As known, many properties of materials depend on their impurity content. There always exists a lower limit of concentration beyond which the effect of an impurity becomes negligible. Therefore, in an analytical study it is usually sufficient to demonstrate that the concentration is smaller than a detection limit chosen to be low enough, so that the effects on the properties of the substance can be neglected.

As a result of this idea, in many problems it is important to know the concentration, or an upper limit for it for as large a number of elements as possible [Pa-66].

In NAA, the gamma spectra obtained from the experiments contain a small number of visible photopeaks which can be used to determine the exact concentration of the corresponding impurities. The other elements produce peaks too small to appear above the background due to their activities and it is then possible only to evaluate "an upper limit of the concentration, which is the sensitivity or detection limit for that particular gamma-spectrum. The operation for the evaluation of detection limits can be summarized as follows: If the element to be determined does not give a visible photopeak in the gamma-ray spectrum, whether a chemical separation has been applied or not, a detection limit is deduced by considering in the zone in which the photopeak should occur.

For this procedure, a gamma sensitivity spectrum is derived from the experimental gamma-ray spectrum by means of a mathematical treatment which is based on the statistical

laws. This is done as; taking a detection limit for each spectrometer channel, the whole set of detection limits constitute the sensitivity spectrum. This spectrum is one in which the channel contents represent the minimum detectable activity of an isotope having a peak centered on the channel. The principles on which the derivation of the gamma sensitivity spectrum is based differ according to the configuration of the part of the spectrum which is considered. A typical gamma spectrum obtained with a NaI(Tl) scintillation detector is given in the figure (3.5).

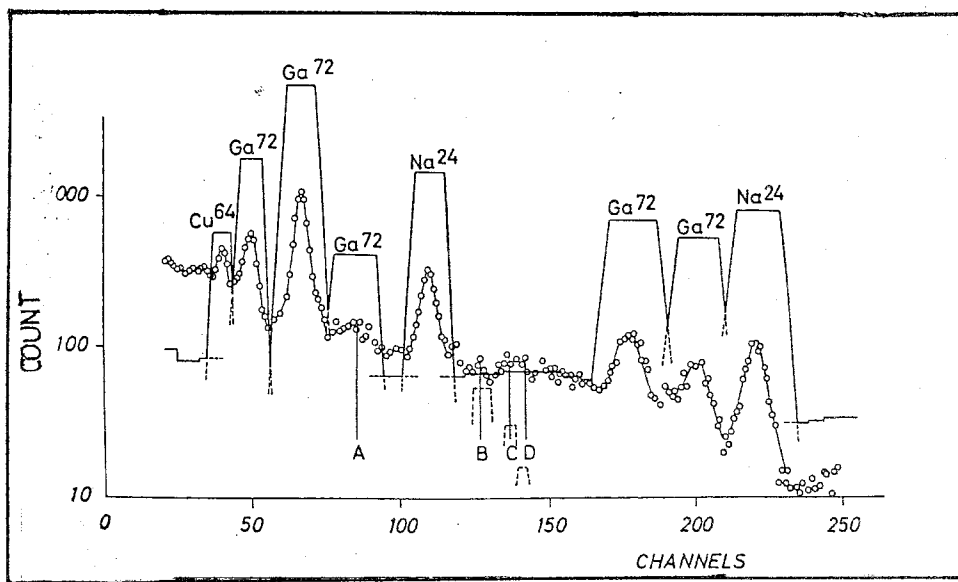


Fig.3.5. Derivation of the sensitivity spectrum from an original one. [Pa-66].

(The detection limits corresponding to the spectrum are seen in the above figure. The part of the curve corresponding to the photopeaks have a form which retains the aspect of the original peak, but the top is flattened. Between these peaks the region corresponding to the background shows a detection limit which varies rather smoothly as a function of the channel number).

As seen, it consists of rather smoothly varying continuous distribution on which several gamma peaks of different radioisotopes are superimposed. For the regions of the spectrum which are far from the visible photopeaks, and where the difference in the counting rates of different isotopes, for adjacent channels is small, the detection limit is calculated from the counting statistics.

In the zone where the counting rates vary smoothly between successive channels, the minimum detectable activity is obtained by considering the activity in the spectrum to be a background and calculating an upper limit of the statistical fluctuations as given below:

If v_1 is the counting rate of the background, it is given by [Se-64, D1-57, Se-59],

$$v_1 = \frac{n_1}{t_1} \pm \frac{n_1^{1/2}}{t_1} \quad (3.15)$$

Similarly, for the observed effect, , it is given as,

$$v = \frac{n}{t} \pm \frac{n^{1/2}}{t} \quad (3.16)$$

where, n and n_1 represent counts; t and t_1 represent counting times. The expressions following the $\pm$ sign are always standard deviations.

For the net activity the below equation is given,

$$\text{Net activity} = v - v_1 \quad (3.17)$$

Since the minimum detectable activity is obtained by taking the activity to be background due to this method described above, it can be $v_1 = v - v_1 \Rightarrow v = 2v_1$ Then using

the formula of propagation of errors,

$\sigma(v) = \left[\sigma^2(v_1) + \sigma^2(v) \right]^{1/2}$, one obtains the standard deviation of the measurement as,

$$\sigma(v) = \left[\frac{n}{t^2} + 2 \frac{n_1}{t_1^2} \right]^{1/2} = \left[\frac{v}{t} + 2 \frac{v_1}{t_1} \right]^{1/2} \quad (3.18)$$

Assuming the counting times t and t_1 are the same, and considering a confidence level required for the problem, $\sigma(v)$ is given by, [Pa-66]

$$\sigma(v) = k \left[\frac{2v_1 + v}{t} \right]^{1/2} \quad (3.19)$$

where the constant k can be taken as 2 corresponding to a confidence level of 95% [Pa-66, Gu-74]

The minimum detectable effect is then derived from this equation by assuming equal to its standard deviations, which finally gives,

$$v = \left(\frac{1}{2} k^2/t \right) \left[1 + (1 + 8tv_1/k^2)^{1/2} \right] \quad (3.20)$$

From the curve representing the gamma detection limits, in fig.3.5, and from the value given by equation (3.20), the minimum detectable weights of the elements are calculated. First, from fig.3.5, it is easy to derive the minimum detectable absolute activity by dividing by the efficiency of the counter. By means of these limits, the minimum detectable weights for each element are then derived by comparing the results to the specific activity which is calculated from the given experimental conditions. Since an element can give several radioisotopes upon irradiation, e.g. ^{72}Ga , a certain

number of different values are obtained for the minimum detectable weight of the element. Among these values, the smallest one is selected for that element.

3.4. MODIFICATION OF THE CODE

As mentioned in the previous subchapter 3.1, the computer code examined in this thesis was prepared originally by using an IBM 370/165 computer with fortran IV level G and assembly languages of the computer used, at the Scientific Information Processing Center, Ispira-Italy. It was then adopted for use with the UNIVAC-1106 at BU. However, since a Calcomp plotting unit was not available in BU and the assembly language depends on the computer used, it was impossible to modify it in the original form to UNIVAC-1106. Therefore, all the plotting subroutines (assembler) which were given as two files, were deleted from the source program. The remaining part of the code which still consists of a main part and thirteen subroutines, was run by using the UNIVAC-1106. This listing is given in the appendix I.

The modified form of the code then was recorded on a new tape for further studies by BU-NE Department.

Some examples of the updating programs(*) used during the procedure mentioned above, are given as follows:

The first program is used to read the original tape and the second is used to record the content of that tape to the main storage of a UNIVAC-1106 computer used in BU. First,

(*) These updating programs are prepared according to EXEC level 36R2 (1978) manual, [EXEC-78], SPERRY UNIVAC, Series 1100, Executive System, Vol.2.

main program and input data were recorded on the same file, and they were then loaded separately into three different files; one being the program file, second being the input data for sample problem, and the third one as spectra for sample problem. To rerun the main program by adding these data files, many corrections were necessary. Programs 3,4,5 and 6 given in the following pages, represent some examples for updating procedure. Program 6 is also used for the rerun of the main program. Finally, program 7 is used for recording all these processes and output data on a new tape and program 8 is for getting a listing of that tape.

Program 1 Reads the tape

@ RUN,
@ ASG, TJ T., 6C9, D80218
@ COB, ISE

IDENTIFICATION DIVISION.

PROGRAM-ID. TEYP-OKUMA.

ENVIRONMENT DIVISION.

CONFIGURATION SECTION.

SOURCE-COMPUTER. U-1106.

OBJECT-COMPUTER. U-1106.

INPUT-OUTPUT SECTION.

FILE-CONTROL.

SELECT T ASSIGN TO UNISERVO T.

DATA DIVISION.

FILE SECTION.

F D T

LABEL RECORDS OMITTED

RECORDING MODE IS F.

O 1 T-REC PIC × (80).

PROCEDURE DIVISION.

AC. OPEN INPUT T WITH NO REWIND.

OPEN OUTPUT F.

OKU. READ T AT END GO TO SON.

DISPLAY T-REC UPON PRINTER.

GO TO OKU.

SON. CLOSE T WITH NO REWIND.

STOP RUN.

@ MAP, IN

@ ×QT

@ FIN

Program 2 Records the tape on the mass-storage

@RUN,
@ASG, TJ T.,6C9, D80218
@DELETE, C GAMMA.
@ASG, UP GAMMA.
@USE F., GAMMA.
@COB, ISE
 IDENTIFICATION DIVISION.
 PROGRAM-ID. TEYP-OKUMA.
 ENVIRONMENT DIVISION.
 CONFIGURATION SECTION.
 SOURCE-COMPUTER. U-1106.
 OBJECT-COMPUTER. U-1106.
 INPUT-OUTPUT SECTION.
 FILE-CONTROL.
 SELECT T ASSIGN TO UNISERVO T.
 SELECT F ASSIGN TO MASS-STORAGE F.
 DATA DIVISION.
 FILE SECTION.
 FD T
 LABEL RECORDS OMITTED
 RECORDING MODE IS F.
 01 T-REC PIC X(80).
 FD F
 RECORDING MODE IS SDF.
 01 F-REC PIC X(80).
 PROCEDURE DIVISION.
 AC. OPEN INPUT T WITH NO REWIND.
 OPEN OUTPUT F.
 OKU. READ T AT END GO TO SON.
 MOVE T-REC TO F-REC.
 WRITE F-REC
 GO TO OKU.
 SON. CLOSE T WITH NO REWIND.
 CLOSE F.
 STOP RUN.

@MAP, IN
@XQT
@DATA, L SAGLAM.
@FIN

Program 3 Updating of main program

@RUN,
@ASG, UP DATA1.
@ASG, A SAGLAM.
@COPY SAGLAM., DATA1.
@ASG, A GAMMA.
@FTN, US GAMMA. MAIN
-183,183
-184
IF (NPK.LE.1) GO TO 333
@PREP GAMMA.
@MAP, I GAMMA.MAIN
IN GAMMA. MAIN
LIB GAMMA..
@XQT GAMMA. MAIN
@ADD DATA 1.
@FIN

Program 4 Updating of the data file

```
@RUN,  
@ASG,A      DATA 1.  
@ASG,UP     DATA 1 (+1).  
@DATA,L     DATA 1., DATA (+1).  
-10,10  
      21  
@END  
@FIN
```

Program 5 Evaluation of a new program and updating of it

```
@RUN,  
@ASG,A     GAMMA.  
@ASG,A     DATA 1.  
@ASG,A     DATA 2.  
@USE    II, DATA 2.  
@ASG,UP    SPECTRA., F ///1000 , UD105  
@COPY      GAMMA., SPECTRA.  
@DELETE, SR    .SPET  
@DELETE, SR    .FIGURE  
@FTN,S            .KAL  
@FTN,S            .SMOOS
```

```
      .  
      .  
      .
```

```
@MAP,I   SPECTRA.MAIN
      IN   SPECTRA.MAIN
      LIB   SPECTRA.
@XQT      SPECTRA.MAIN
@ADD      DATA 1.
@ADD      DATA 2.
@FIN
```

Program 6 Rerun of the main program with data

```
@RUN,
@ASG,A    SPECTRA.
@ASG,A    TKR.
@USE      11, TKR.
@ASG,A    INPUTX.
@ASG,UP    INPUTX(+1).
@DATA,L    INPUTX., INPUTX(+1).
-10,10
      22
      SCX 1119.0    3.12 5.675E-06 3.900E+00
@END
@XQT      SPECTRA. MAIN
@ADD      TKR.
@ADD      INPUTX.
@FIN
```

Program 7 Records to a new tape from the mass-storage

```
@RUN,  
@ASG,A    SPECTRA.  
@ASG,T    T.,16N,SAGLAM  
@MOVE    T.,1  
@COPY,GM    SPECTRA.,T.  
@FIN
```

Program 8 Taking the listing of the tape

```
@RUN,  
@ASG,T    T.,16N,SAGLAM  
@UNIVAC*TTTC.TTTC T.  
@FIN
```

4. COMPARISON OF GRETTEL WITH ANOTHER CODE, CORGAM

The importance and use of digital computers in NAA was mentioned in the previous subchapters. It was also mentioned that there are many computer codes dealing with this problem. Each of these codes have their own techniques in analysing, in unfolding complex gamma-ray spectra. Some of them study the estimating of intensity coefficients, as Corgam. There are many methods for estimating these coefficients. These methods are stripping, least squares, iterative methods, vector analysis and stepwise regression [Ec-68, Le-65].

One of these codes mentioned above, prepared by NEA-CPL in the experimental data processing field is Corgam. Corgam was studied before as a subject of another thesis introduced at NE Department. Since these two codes, Corgam and Gretel, deal with NAA, but their logics for the approaching to the problems are quite different from each other, a short comparison of these two computer codes will be given in this chapter. Therefore first, Corgam will be introduced in a summary form in the following subchapter.

4.1. INTRODUCTION OF CORGAM

Corgam, a correlation algorithm for gamma-ray spectra, allows the reference gamma-ray spectra to be correlated with a complex gamma-ray spectrum [Ec-69]. This correlation algorithm is based on a least squares procedure which is a logical method for determining the intensity coefficients for NAA.

The analysis used by Corgam is described as the least squares procedure for fitting a dependent variable, e.g., a complex spectrum, to a set of independent variables, e.g., a set of reference spectra.

Corgam sets up a mathematical model in order to make an analysis by using least squares procedure. The formulation of Corgam can be summarized as follows:

The gamma-ray spectrum of a complex sample containing several isotopes (m in number) may be represented by [Ec-69],

$$Y_i = \sum_{j=1}^m B_{ij} \quad i = 1, 2, \dots, n \quad (4.1)$$

where,

Y_i = the counts registered in channel i of a multi-channel analyzer.

B_{ij} = the counts registered in channel i from the j^{th} constituent isotope.

The reference gamma-ray spectra of each of the m constituent isotopes are related to the complex gamma-ray spectrum by,

$$Y_i = \sum_{j=1}^m X_{ij} B_j + e_j \quad i = 1, 2, \dots, m \quad (4.2)$$

where,

B_j = the intensity coefficient for isotope, j , i.e., a coefficient proportional to the quantity of isotope j in the complex sample.

X_{ij} = the counts registered in channel i from the j^{th} isotope of known quantity.

e_i = the deviation at point i of the estimated value of Y_i , $\sum X_{ij} B_j$, from the observed value Y_i .

There are some difficulties in applying least squares technique to the unfolding of gamma-ray spectra. One of them is that some calculated intensity coefficients may have negative sign. For physical reasons, they must be non-negative. Besides, the least squares method may yield positive intensity coefficients for all reference spectra, but some of these may not be significant according to the tests applied. The problem is to determine which reference spectra are not significant.

In order to overcome these difficulties, some statistical tests are applied such as t value for the null hypothesis which is calculated for each intensity coefficient. These values are used in the correlation algorithm to eliminate those reference spectra for which the calculated intensity coefficients are negative or non-significant.

As a summary, it can be stated that Corgam,

a) compensates for electronic instabilities in the multichannel analyzer,

b) corrects for the background, subtracts it from the complex spectrum,

c) normalizes the data to a fixed neutron flux level,

d) represents a mathematical model in matrix form for reference spectra and allows a choice of weighting factors,

e) allows a choice of methods for the calculation of standard deviations.

f) it corrects for decay time,

g) may also consider time dependent spectra for those isotopes which have isomers.

Finally, Corgam is a novel method of unfolding complex gamma-ray spectra because the variance calculations include contributions from both the complex and reference spectra. Also, only reference spectra that have intensity coefficients which are significant at a preselected level are retained in the output.

4.2. COMPARISON OF THE TWO CODES

As mentioned in the previous chapters, Gretel and Corgam are two computer programs prepared for NAA problems. However, their mathematical approach to these problems are quite different. One of the reasons for these differences is that they are prepared to analyze data from different kinds of detectors; Gretel uses data from a Ge(Li) detector, Corgam from a NaI(Tl) scintillator. Therefore, their data handling procedures are different from each other. Besides that, many authors apply their own techniques to unfold gamma-ray spectra. For the comparison of these two codes, this subchapter gives the main specifications in the following form:

CORGAM is a computer code of unfolding complex gamma-ray spectra by using reference spectra. In the analysis, intensity coefficients of the isotope spectra which are significant at a preselected level are retained in the output and the variance calculations include contributions both from the reference and the complex spectra [Bö-82, Ec-69]. Its specifications can be given in a brief form as follows:

1. Corgam is based on estimating the intensity coefficient. It uses the least squares procedure for fitting a dependent variable, i.e., a complex spectrum, to a set of independent variables, i.e., a set of reference spectra [Ec-68].

2. Corgam sets up a mathematical model and by using the least squares procedure, correlates the reference gamma-ray spectra with a complex one.

Example: The program multiplies any spectrum of a reference isotope with a coefficient and similarly other spectra with other coefficients. These coefficients are defined as intensity coefficients. Then, adding the contents of these spectra channel by channel, a third spectrum is obtained. Subtracting the complex spectrum due to an unknown sample from this third one, their difference is obtained. By using the least squares procedure the program tries to calculate these coefficients for the minimization of that difference. Those intensity coefficients are proportional to the quantity of isotope in the complex spectrum. Finally, the quantity of the isotope in the complex can be given in terms of ppm.

3. The reference spectra and the complex one required for Corgam are collected for the fixed counting time, using the scintillation detectors.

4. Before the analysis of the spectra, Corgam compensates for electronic instabilities in the spectrometer system used, because these instabilities can cause shifts in these spectra. To compensate for these inconsistencies, the gamma ray photopeaks of the spectra are fit to a Gaussian function using the least squares method.

5- Corgam corrects for the background and decay times. It normalizes the data to a fixed neutron flux level.

6- Corgam in the standard form is restricted to a problem of 400 channels and 15 reference spectra.

7- Corgam requires less than 1 minute of execution time with the restrictions mentioned above, on the computer IBM 360/50.

8- Corgam also considers time dependent spectra for those isotopes which have isomers.

9- The program was recorded originally for IBM 360/50 using fortran IV level G. It was then updated by Bögrün, [Bö-82], for UNIVAC-1106. The updated form of the code was recorded on to a magnetic tape for further use in BU-NE department.

10- An intermediate disk storage is required for Corgam. The source deck contains 1303 cards.

GRETEL is a computer program which is set up for routine batchwise processing of spectrometric data [Gu-74]. The characteristics of Gretel can be given as follows:

1- Gretel performs the quantitative analysis of gamma-ray spectra using special oriented libraries which are prepared for each particular problem.

2- The gamma-ray spectra needed for Gretel, are obtained by Ge(Li) detectors.

3- For the analysis of the spectra, a preliminary data handling procedure required was discussed in detail in the

chapter 3. As a summary, this procedure can be given as: First, smoothing of the spectra is done. Then, the peak locations are searched by using some tests to avoid the statistical fluctuations being kept as analytical peaks. Computation of the peak areas is the third step of the program. Two different methods are used for that purpose. Since Gretel examines any spectrum peak by peak, for the visible peaks a method derived by Yule, [Yu-69], and for hidden peaks new detection limits are determined.

4- In Gretel, compton continuum at the fixed level is subtracted from the computed areas of the photopeaks. Also, decay time corrections are made.

5- The restrictions for Gretel in the standard form, are given as; 50 elements as the maximum number for the oriented library which corresponds the reference isotopes in Corgam, and 4096 channels as the maximum number of the spectrometer used.

6- For a typical analysis, on the average, of 30 spectra of 2048 channels (a calibration spectrum plus 29 unknowns) with a library of 30 elements the running time is about 30 seconds.

7- The memory requirements are not much for Gretel.

8- Gretel was recorded originally for IBM 370/165 using fortran IV level G. The code was then updated by the author for UNIVAC-1106, by deleting all the plotting subroutines since there was not available a Univac plotting unit. The modified form of the code is recorded on a new magnetic tape for further studies by BU-NE department.

As a result of these specifications and others given in 4.1, one can conclude that Corgam is a more comprehensive program prepared for the activation analysis by using a NaI(Tl) scintillation detector and a small channel region in the spectrometer used. On the other hand, Gretel is one of first ones prepared for NAA problems by using a Ge(Li) detector which has a better energy resolution than a scintillator. The program also gives analog outputs with a Calcomp unit and it performs the quantitative analysis.

5. DISCUSSION AND CONCLUSIONS

The basic principles, the use and importance of neutron activation analysis, NAA, was discussed in the previous subchapters. It was also mentioned that the determination of the microconstituents of an unknown sample was possible down to the ppm (parts per million) level or even to the ppb (parts per billion) level. As also discussed in the previous chapter, determination of a large number of elements in one sample and the treatment of the comparison samples are very troublesome and time consuming. Therefore, instrumental activation analysis has been improved as one of the major techniques for the control and measurement of the trace element distribution in a large number of samples. The speed of this technique is based to a large extent on the availability of nuclear reactors with a high and reliable neutron flux. It also depends on the development of the high resolution gamma-ray spectrometers with the computer hardware and software. The computer programs for NAA problems are usually prepared for the spectra obtained especially by two types of detectors, such as NaI(Tl) scintillators and Li drifted Ge or Si semiconductors. Recently, semiconductor detectors with excellent energy resolution pushed the number of energy channels required for an adequate spectral response up into the thousands.

As an example, a computer code for NAA based on a GE(Li) spectra was studied in this thesis as a part of it. The program which is set up for routine batchwise processing of spectrometric data, was first prepared in 1966 as one of

the earliest codes developed for Ge(Li) spectra. This program has then been revised and rewritten to follow the continuous development of the analytical instrumentation, such as multi-channel analyzer, calcomp unit of the computer used, etc. The new program, called Gretel, performs the quantitative analysis of gamma-ray spectra obtained by a Ge(Li) detector which has a better energy resolution than a scintillator. The data-reduction system required for Ge(Li) spectra is used in Gretel. This system involves the smoothing of the spectra, differentiation of these data, the determination of peak locations and peak boundaries, the computation of peak areas and the associated errors. The program gives the results in terms of ppm in the form of digital lists. In addition, with some modifications in the main program, analog outputs can also be obtained by means of a Calcomp unit if it is available. Also, Gretel performs the search of the small peaks not detected over the background and gives the minimum detectable activity.

However, one can not conclude that Gretel is an integral program. Because the program fails in the identification of isotopes when there are interferences coming from different isotopes, e.g., Sc and Zn. To avoid this problem, the preparation of the libraries for each kind of problem, which is a responsibility of the analyst, is very important. Some suggestions can be made to correct this shortcoming. First, for the preparation of any library, "analytical peaks" should be chosen. These peaks are those given the best sensitivity and the highest chance of being free from interference within a range of energy which depends on the resolution of the detector. Second, when there are interfering isotopes, some modifications can be applied to the main program such as control peaks. Introducing other peaks due to interfering isotopes, one can check whether these isotopes are present or not. Third, the energy range in which the interfering peaks are

searched could be restricted considerably by using a pre-amplifier with a high resolution power. Fourth, since the single Ge(Li) measurements are not adequate so the coincidence measurements may be offered for analysis of gamma-ray spectra free of interference. However, coincidence technique is not much successful for Ge(Li) detectors because the reduced detector efficiency seems to be a serious obstacle. •

The other program, Corgam, described briefly in the Chapter 4 to make a comparison, examines the spectra obtained by a NaI(Tl) scintillation detector. As mentioned in the above subchapters, Corgam has a more comprehensive mathematical model; so it can be stated that it is one of the better programs for the resolution of gamma-ray spectra with the correlation algorithm. Besides, the program can be applied to other kinds of analysis problems, such as the analysis of a spectrum in the ultraviolet light region, etc. On the other hand, Corgam is restricted to a problem of 400 channels and 15 reference spectra due to the level of accuracy desired. Also, the computer memory requirements are large for the correlation algorithm.

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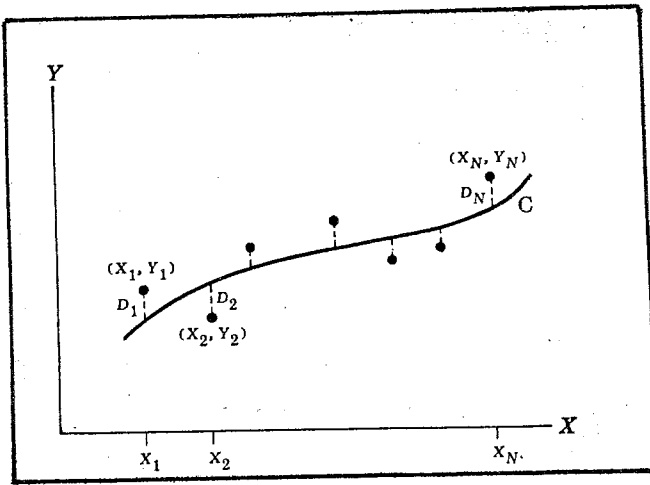
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APPENDIX A

THE METHOD OF LEAST SQUARES

A measure of the fit for the curve C given in the figure, to the given data is provided by the quantities $D_1^2 + D_2^2 + \dots + D_N^2$. Those D's are denoted as deviations or errors.



If D's are small, it is said that the fit is good, if they're large, the fit is bad.

Of all the curves approximating a given set of data points, the curve having the property that $D_1^2 + D_2^2 + \dots + D_N^2$ is a minimum is called a best fitting curve in the sense of least squares. It is also called a least square curve. Thus a line having this property is called a least square line; a parabola with this property is called a least square parabola, etc. [Mu-61].

The least square line: The least square line approximating the set of points $(x_1, y_1), (x_2, y_2), \dots, (x_n, y_n)$ has the equation, [Ec-68],

$$y_i = A + Bx_i + e_i \quad i = 1, 2, \dots, n \quad (1)$$

where e_i is the deviation at point i of the estimated value, from its observed value y_i , A and B are the true values of the intercept and slope respectively. The estimators a and b for A and B , can be obtained by minimizing the sum of squared errors. If each data point has a different, statistical, accuracy, then each residual is multiplied by a proper weight, w_i . Therefore the sum of squared errors is given by, [Ec-68, So-41],

$$S = \sum_{i=1}^n e_i^2 = \sum_{i=1}^n [\bar{y}_i - (A+Bx_i)]^2 w_i \quad (2)$$

For this problem w_i are equal to 1. Then,

$$S = \sum_{i=1}^n [\bar{y}_i - (A+Bx_i)]^2 \quad (3)$$

S is minimized by differentiating with respect to A and B , and equating the results to zero. Then a and b the estimators are substituted for A and B . The resulting equations are,

$$\sum_{i=1}^n [\bar{y}_i - (a+bx_i)] = 0 \quad (4)$$

$$\sum_{i=1}^n \{x_i [\bar{y}_i - (a+bx_i)]\} = 0 \quad (5)$$

Equations (4) and (5) are rearranged into a form called the normal equations [Mu-61],

$$an + b \sum_{i=1}^n x_i = \sum_{i=1}^n y_i \quad (6)$$

$$a \sum_{i=1}^n x_i + b \sum_{i=1}^n x_i^2 = \sum_{i=1}^n x_i y_i \quad (7)$$

Solution of these equations for a and b, results in,

$$b = \frac{\sum_{i=1}^n x_i y_i - ((\sum_{i=1}^n x_i)(\sum_{i=1}^n y_i))/n}{\sum_{i=1}^n x_i^2 - (\sum_{i=1}^n x_i)^2/n} = \frac{\sum_{i=1}^n (x_i - \bar{x})(y_i - \bar{y})}{\sum_{i=1}^n (x_i - \bar{x})^2} \quad (8)$$

$$a = \bar{y} - b\bar{x} \quad (9)$$

where

$$\bar{y} = (\sum_{i=1}^n y_i)/n \quad \text{and,} \quad (10)$$

$$\bar{x} = (\sum_{i=1}^n x_i)/n \quad (11)$$

Then, equation (1) can also be rewritten in terms of the estimators b and mean value $\bar{x}$, $\bar{y}$ through the use of equation (9) as,

$$y_i = \bar{y} + b(x_i - \bar{x}) \quad (12)$$

APPENDIX B

ISOTOPE BUILD-UP AND DECAY

Considering N_1 , the number of parent and N_2 , the number of daughter nucleus as functions of decay constants and time, it is possible to write the decay rate of the parent nuclei as,

$$\frac{dN_1}{dt} = -\lambda_1 N_1 \quad (1)$$

or

$$\frac{dN_1}{dt} = \lambda_1 N_{1,0} e^{-\lambda_1 t} \quad (2)$$

where,

$N_{1,0}$: number of parent nuclei initially present
 λ_1 : decay constant of parent

The analogous equation for the daughter nuclei is given by,

$$\frac{dN_2}{dt} = R_T - \lambda_2 N_2 \quad (3)$$

Since R_T represents the rate of formation of the daughter, it is also given as,

$$\frac{dN_2}{dt} = \lambda_1 N_{1,0} e^{-\lambda_1 t} - \lambda_2 N_2 \quad (47)$$

which is a 1st order linear differential equation.

For the solution of that equation, choosing $y = N_2$,
 $x = t$, $P(x) = \lambda_2$ and the integrating factor as,

$$e^{\int P(x) dx} \Rightarrow e^{\int \lambda_2 dt} = e^{\lambda_2 t}$$

then, multiplying both sides of the original differential equation, eg.(4), with this factor, it is obtained as, [Wy-66]

$$e^{\lambda_2 t} \frac{dN_2}{dt} + e^{\lambda_2 t} \lambda_2 N_2 = \lambda_1 N_{1,0} e^{\lambda_2 t} e^{-\lambda_1 t}$$

where the sum of two terms on the left side are equal to the term $\frac{d}{dt} (e^{\lambda_2 t} N_2)$.

Then,

$$\frac{d}{dt} (N_2 e^{\lambda_2 t}) = \lambda_1 N_{1,0} e^{(\lambda_2 - \lambda_1)t}$$

Integrating both sides from $t=0$ to $t=t$ and rearranging it for N_2 ,

$$N_2 = \frac{\lambda_1}{\lambda_2 - \lambda_1} N_{1,0} (e^{-\lambda_1 t} - e^{-\lambda_2 t}) + C e^{-\lambda_2 t} \quad (5)$$

If the amount of daughter is initially zero, i.e., at $t = 0$, $N_2 = N_{2,0} = 0$, then the second term of the equation (5) at the right side which represents the decay of daughter vanishes.

Finally, one obtains

$$N_2 = \frac{\lambda_1}{\lambda_2 - \lambda_1} N_{1,0} (e^{-\lambda_1 t} - e^{-\lambda_2 t}) \quad (6)$$

APPENDIX C

SINGLE COMPARATOR METHOD

Neutron activation analysis usually consists of 5 steps, such as 1. irradiation, 2. purification, 3. measurement, 4. comparison, 5. verification [Ad-70].

For the first step, suitable amounts of the sample and standard are weighted and irradiated simultaneously for a predetermined time. The sample and standard sizes should be chosen so that the intensity of radiation emitted from the desired radioelement will be approximately the same in both sources. There are three methods for the analysis such as a) absolute method, b) comparator or relative method, c) single comparator method. Single comparator method is used in Gretel so it will be described briefly in the following:

The use of germanium-lithium drifted detectors allows the simultaneous determination of an important number of elements after a single irradiation. The classical comparator method requiring one comparator for each element to be determined hence becomes less appropriate for routine work. These operations are time-consuming and can also introduce sources of error. On the other hand, the absolute method which performs direct calculation of weights of the isotopes from an estimation of the activation formula, suffers from the inaccurate knowledge of several constants.

In the single comparator method, the known weights of the elements are irradiated together with a suitable flux monitor under specified conditions (A known weight of a reference element is taken as flux monitor). The unknown sample is then irradiated with a similar flux monitor which is used

as a single comparator for different elements [Gu-65]. This method is suitable especially for routine or automated determinations; mainly when dealing with isotopes with a very short half-life, when a large number of elements is determined or when standards and samples could not be irradiated together.

There is also a very similar method which utilizes elements present in known quantities in the matrix itself as internal standard (as in Gretel), with the advantage of correcting flux perturbations within the matrix.

Principle of the method can be summarized as follows:

The weight, w , of an irradiated element is related to the photopeak counting rate, A_p , of the radioisotope measured by the relation, [Gu-65],

$$w = \frac{A_p}{\phi S D} \frac{M}{\sigma N E a \delta} \quad (1)$$

where,

M : atomic weight of the irradiated element

δ : isotopic abundance of the target nuclide.

σ : effective activation cross section for the neutron energy spectrum used

N : Avagadro's number

E : efficiency of the detector for the gamma ray measured

a : gamma-ray abundance in decay scheme

ϕ : neutron flux

S : saturation factor $(1 - e^{-\lambda T})$ depending on irradiation time T and decay constant λ

D : decay factor $(e^{-\lambda t})$, where t is decay time

When the neutron flux is measured by irradiating a known weight of a reference element (neutron flux monitor) and measuring the induced radioactivity by gamma-ray spectrometry and using the other known quantities due to this element, the same relation holds:

$$\phi = \frac{A_p^*}{S^* D^* W^*} \frac{M^*}{\sigma^* \delta^* N E^* a^*} \quad (2)$$

where the asteriks refer to the neutron flux monitor which is also used as single comparator for the elements to be determined in a sample. By substituting ϕ taken from eg.2 in eg.1 one obtains,

$$W = k \frac{A_p}{A_p^*} \frac{S^* D^*}{S D} W^* \quad (3)$$

where

$$k = \left(\frac{\sigma^* \delta^* N E^* a^*}{M^*} \right) / \left(\frac{\sigma \delta N E a}{M} \right) \quad (4)$$

APPENDIX D

CONVOLUTION METHOD

In much experimental work, the information may be obtained in the form of a two-column table of numbers, e.g. A vs. B. There are some methods for handling of this type of data [Sa-64].

One of the simplest ways to smooth fluctuating data is by a moving average. In this procedure, one takes a fixed number of points, adds their ordinates together, and divides by the number of points to obtain the average ordinate at the center abscissa of the group. Next, the point at one end of the group is dropped, the next point at the other end added, and the process is repeated.

Table (1) illustrates how the moving average might be obtained. According to this table, the set of numbers on the right are the data or ordinate values, those on the left, the abscissa information. The outlined block in the center may be considered to be a separate piece of paper on which a new set of abscissa numbers are written ranging from -2 through 0 to +2. The C's at the right represent the convoluting integers. For the moving average, each C is numerically equal to one. To perform a convolution of the ordinate numbers in the table of data with a set of convoluting interges C_i , each number in the block is multiplied by the corresponding number in the table of data, the resulting products are added and this sum is divided by five. The set of ones is the convoluting function and the number by which we divide, in this case 5, is the normalizing factor. To get the next point in the moving average, the center block is slid down one line and process repeated.

1800.0			705	Table 1. Convolution operation.
1799.8			712	
1799.6			717	
1799.4			718	
1799.2	$ x_0 - 2$	C_{-2}	721	Abscissa point on the left, tabular data on the right. Operation is the multiplication of the data points by the corresponding C_i, summation of the products and division by a normalizer, resulting an a single convolute at point x_0.
1799.0	$ x_0 - 1$	C_{-1}	722	
1798.8	$ x_0$	C_0	725	
1798.6	$ x_0 + 1$	C_1	730	
1798.4	$ x_0 + 2$	C_2	735	
1798.2			736	
1798.0			741	
1797.8			746	
1797.6			750	

The concept of convolution can be generalized beyond a simple moving average. In the general case the C 's represent any set of convoluting integers. There is an associated normalizing factor. The mathematical description of the process described above, is given as;

$$Y_j = \frac{\sum_{i=-m}^{+m} C_i Y_{j+i}}{N} \quad (1)$$

The index j represents the running index of the ordinate data in the original data table. For the moving average, each C_i is equal to one and N is the number of convoluting integers. However, for many types of data, the set of all i 's which yields the average, is not particularly useful. For example, on going through a sharp peak, the average would tend to degrade the end of the peak. There are different types of smoothing functions. Figure (1) illustrates these functions.

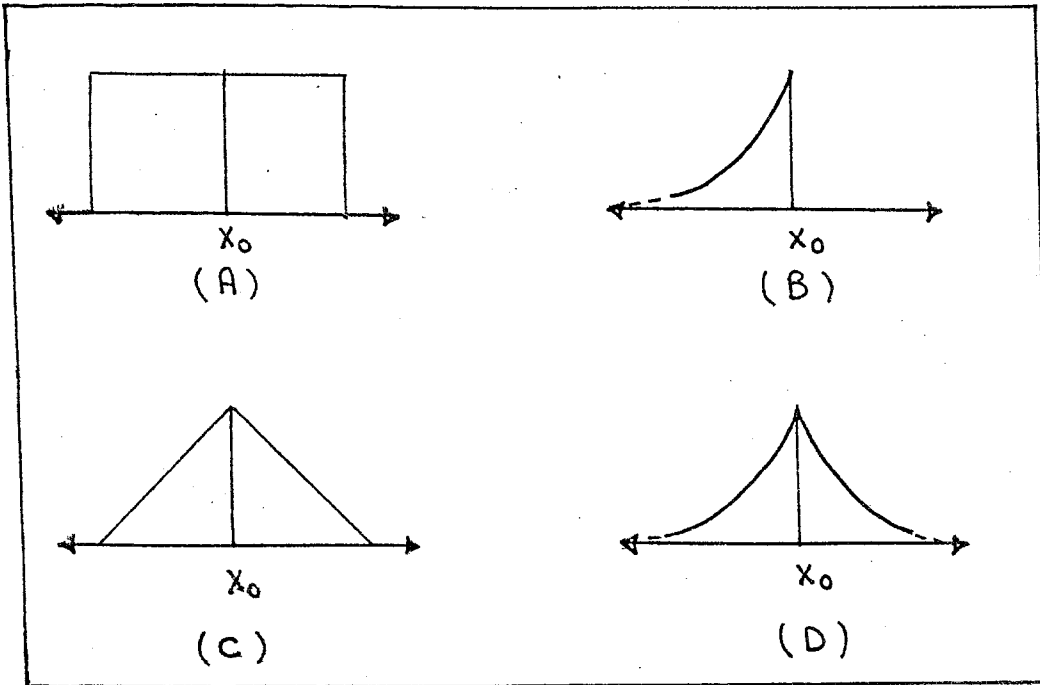


Fig.1- Various convolute functions

A. Moving average. B. Exponential function. C. Symmetrical triangular function. D. Symmetrical exponential function.

The function in the figure 1.B is an exponential set which simulates the exponential law. The usual spectrum from a spectrophotometer is the resultant of two convolutions of the actual spectrum of the material; first with a function representing the slit function of the instrument which is given in the figure 1.C, and then this first convolute spectrum is further convoluted with a function representing the time constant of the instrument.

APPENDIX E

LISTING OF COMMON VARIABLES USED IN GRETAL

The common variables used in Gretal can be given in the following form:

ELI	: Symbol of the radioelement in calibration library
ENI	: Energy of this isotope in keV
ICAN	: Channel in which the calibration peak falls
INT	: Interval for the search of the calibration peak (in channels)
ELEM	: Symbol of the radioelement in oriented library
EN	: Energy of this isotope in keV
FWT	: Full width at half maximum in keV (oriented lib.)
XLAM	: Decay constant of the radiosotope in minutes (oriented lib.)
AA	: Specific activity (oriented lib.)
COUNT	: Counts in each channel
SPSM	: The smoothed values in each channel
DER	: Differentiated values of the contents in each channel
ISTAR	: First channel of the spectrum
IEND	: Last channel of the spectrum
INDEX	: Index indicating the kind of the peak found
JMAX	: Number of peaks found by PKF
FMEDIO	: Average value of the geometrical correction factor
AREA	: Area of the photopeak calculated by SURF
ERR	: Error calculated during the evaluation of peak areas
FWHM	: Full width at half maximum
LS	: Left limit of the peak found
CENTR	: Center of the peak found by PKF

LD : Right limit of the peak found

PMOL : Molecular weight of the element given in the oriented library

DECAY : Decay constant of the radioisotope (calibration lib.)

VALA : Specific activity for a μ gram sample irradiated in a thermal flux of 10^{13} neutrons/cm²sec (calib.lib.)

PESO : Weight of the element in grams (calib. lib.)

TDECI : Decay times of radioisotopes calculated by TDEC

TIRR : Irradiation duration in minutes

TCON : Counting duration in minutes

ALFA : Slope of the calibration line

BETA : Intercept of the calibration line

FLUX : Neutron flux during irradiation

WEIGHT : Weight of the sample in grams

STAND(6,K):Area zero/microgram of the element

NELST : Number of standards in standard library

IREL : A switch indicating the way which is followed by the program

NIRR : Irradiation number

MO : The beginning of the measurement being year, month, day, hour and minute

ANOME : The name of the problem involved

ASAMPL : Identification of the sample

EL4 : Symbol of the radioisotope found by FIRST for the calibration line

CKK : Center of the peak found by FIRST

VKK : Energy of the peak found by FIRST

VMAX : Counts due to the peak found by FIRST

NPK : Number of the peaks in the calibration lib.

EL3 : Symbol of the radioisotope found by SENSIT

VS : Counts due to that peak

P2 : Energy of that peak in keV

CS : Center of that peak

VS2 : Right limit of that peak
VD2 : Left limit of that peak
CLS2 : Channel contents at the right limit of the peak
CLD2 : Channel contents at the left limit of the peak
ISENS : Number of peaks found by SENSIT
EL2 : Symbol of the radioisotope found by PKF
VC : Counts due to the single peak found by PKF
PI : Energy of that peak in keV
VS1 : Right limit of that peak
VD1 : Left limit of that peak
CC : Center of that peak
CLSI : Channel contents at the right limit of that peak
CLDI : Channel contents at the left limit of that peak
IPICK : Number of single peaks found by PKF
FDEC : Decay factor calculated as $e^{\lambda t_{dec}}$ by ANALYS
FSAT : Saturation factor calculated as $(1 - e^{-\lambda t_{irr}})$ by ANALYS
DENS : Density of the sample
EL2D : Symbol of the radioisotope found by PKF
VCD : Counts due to double peaks found by PKF
PID : Energy of the doublet in keV
CCD : Center of one of the doublet
VID : Left limit of that peak (or right limit according to peak)
CLID : Channel contents at the limit of that peak
IPCD : Number of doublet
EL2REL : Symbol of the radioisotope found by PKF when there is a standard library
EL3REL : Symbol of the radioisotope found by SENSIT when there is a standard library
EL2RDRE : Symbol of the radioisotope which gives a doublet when there is a standard library

APPENDIX F

LISTING OF THE INPUT DATA AND INPUT OPTIONS

Input for Gretel consists of:

i) IME: One card for the computation of decay times (in minutes) of the radioisotopes after irradiation in the reactor. (It indicates the days in each month of the year)

FORMAT (12I6)

ii) LIBST: One card indicating the number of peaks of the calibration library.

FORMAT (I6)

iii) One card for each gamma peak of the calibration library. This card contains:

ELI	: Symbol of the radioelement
ENI	: Energy in keV
ICAN	: Channel in which the calibration peak falls
INT	: Interval for the search of the calibration peak (in channels)
DECAY	: Decay constant of the isotope to which the peak belongs ($\lambda = 0.693/T_{1/2}$ in minutes)
VALA	: Specific activity for a g sample irradiated in a thermal flux of 10^{13} neutrons/cm ² sec
PESO	: Weight of the element in grams

where the formats are given as,

FORMAT (A6, F7.2, 2I4, 2E10.4, F12.8)

iv) NLIB: One card indicating the number of peaks of the "oriented analytical library".

FORMAT (I6)

v) One card for each gamma peak of the above mentioned library.

This card contains:

ELEM : Symbol of the radioelement of the oriented library
EN : Energy in keV
FWT : Full width at half maximum in keV
XLAM : Decay constant of the radioisotope in minutes
AA : Specific activity
PMOL : Molecular weight of the element

Where the format is given as,

FORMAT (A8, 2F10.3, 2E10.4, F10.5)

vi) One card giving irradiation and problem characteristics:

FLUX : Neutron flux during irradiation in neutrons/
cm² sec
TIRR : Irradiation duration in minutes
IO : End of the irradiation stating its year, month
day, hour and minute
NIRR : Irradiation number
ANOME : The name of the problem involved
DENS : Density of the sample, to compute atoms/cm³
for each element

Put 1 if not requested.

FORMAT (E10.4, F10.3, 5I4, I5, 5A4, F7.3)

vii) One card with a set of figures and switches allowing the program to flow in different ways:

NSPT : Number of spectra to elaborate
NCH : Number of channels for each spectrum. In the case in which a partial data-out has been performed,

NCH = 2048

NSMO : Number of points to be used for the smoothing of the spectra

NSMO = 5 five point smoothing

NSMO = 7 seven point smoothing

NSMO = 9 nine point smoothing

IPDO : Indicates if a partial data-out has been used:

IPDO = 0 the whole spectrum has been punched

IPDO = 1 partial data-out

IASCI : Indicates the perforation code used:

IASCI = 1 IBM code

IASCI = 2 ASCII code

IASCI = 3 ASCII code, no parity

where the format is given as,

FORMAT (6I6)

viii) In the case of partial data-out (IPDO = 1) one card for each measurement, containing:

MO : year, month, day, hour and minute of the beginning of the measurement

ASAMPL : identification of the sample

WEIGHT : weight of the sample in grams

TCON : counting duration in minutes

ISTAR : first channel of the partial spectrum

IEND : last channel

ITR : switch indicating whether to elaborate the spectrum or not

ITR = 1 elaboration

ITR = 0 no elaboration

IS : switch indicating whether the area of the peaks must be computed with the original spectrum or with the smoothed one

IS = 0 original spectrum

IS = 1 smoothed spectrum

Where the format is given as,

FORMAT (5I4, 3A4, F10.6, F10.3, 2I5, 3I2)

ix) According to the values of IASCI, one of the following cards is read for each measurement:

a) IASCI = 1 or 2; the card contains:

MO, ASAMPL, WEIGHT, ITR, IS

FORMAT (5I4, 3A4, F10.6, f0X, 2I2)

b) IASCI = 3; the card contains:

MO, ASAMPL, WEIGHT, TCON, ITR, IS

FORMAT (5I4, 3A4, F10.6, F10.3, 10X, 2I2)

APPENDIX G

DESCRIPTION OF GRETEL

Gretel, as mentioned in the chapter 3.1, performs the quantitative analysis of gamma-ray spectra obtained by Ge(Li) detectors. The program is composed of a MAIN part which controls the various subroutines and allows the reading of the input data and library from punched cards, a number of subroutines and three functions.

As also mentioned in the same subchapter and in the 3.4, some modifications due to elimination of some subroutines belong to plotting procedure were required on original Gretel in order to adopt it to UNIVAC 1106. After having eliminated these routines, the program reran. These thirteen subprograms called by the main program are given in the following pages in detail.

The gamma-ray spectra of up to 4096 channels obtained by the multichannel analyzer, are first punched on paper tapes and read by one of the routines, TAPE, TAPEAS or TAPEGA, according to the punching code adopted, such as IBM, ASCII or ASCII with no parity. The routines transform the data in a suitable format, either by the function CONV or CONVB.

The routines used in Gretel, except the five mentioned above, are listed in the order called:

- 1- Function TDEC
- 2- Subroutine SMOOS
- 3- Subroutine DERIV
- 4- Subroutine FIRST
- 5- Subroutine KAL

- 6- Subroutine ANALYS
- 7- Subroutine PKF
- 8- Subroutine SURF
- 9- Subroutine SURF2
- 10- Subroutine SENSIT
- 11- Subroutine PRINT

The following is the description listing of the main and subprograms called by the main part. These descriptions are given with their options.

MAIN PROGRAM

The main part of the program reads the data of the calibration source, the oriented or standard libraries and some other input parameters, such as weight of samples and neutron flux. It also controls the various subroutines and functions.

It contains a number of common statements at the beginning of the program. These are described in the appendix E.

1. Function TDEC

```
FUNCTION TDEC (MO, IO, IME)  
  DIMENSION MO(5), IO(5), IME(12)
```

Where MO, IO and IME are defined in the appendix F. MO and IO are in terms of year, month, day, hour and minute so they are given as MO(5), IO(5).

This function calculates the decay times of isotopes in minutes.

2. Subroutine SMOOS

SUBROUTINE SMOOS (NSMO)

NSMO : number of points to be used for the smoothing
of the spectra

NSMO = 5 five point smoothing

NSMO = 7 seven point smoothing

NSMO = 9 nine point smoothing

This routine performs the smoothing of the spectra on five, seven or nine points in order to reduce the influence of the statistical fluctuations.

3. Subroutine DERIV

SUBROUTINE DERIV (NSMO)

NSMO : it is explained in the above routine

It gives the first derivative of the smoothed spectrum. The shape of this spectrum is then analyzed in each given interval.

4. Subroutine FIRST

SUBROUTINE FIRST (LIBST, NCH, III, JS)

DIMENSION CK(10), VK(10), FCOR(10)

LIBST : number of peaks of the calibration library

NCH : number of channel of the spectrum

III : number of spectra to elaborate

JS : switch indicating whether the area of the peaks must be computed with the original spectrum or with the smoothed one

JS = 0 original spectrum

JS = 1 smoothed spectrum

CK : centers of the peaks found
VK : energies of these peaks
FCOR : correction factor for counting geometry to
the specific activities of the elements
belonging to the oriented library and
finally the calibration line equation.

It is given as, $\frac{\text{exp. specific activity}}{\text{theoretical specific activity}}$

Subroutine FIRST treats the first spectrum of the series that is calibration spectrum. There is a certain number of known peaks in that spectrum. The calibration spectrum is always the first one in the set.

5. Subroutine KAL

SUBROUTINE KAL (N,C,E,A,B)
DIMENSION C(10), E(10)

N : number of measurement
C : number of channel used as abscissa for the calibration line
E : gamma-ray energies used as ordinate for the calibration line
A : slope of the calibration line
B : intercept of that line

Subroutine KAL calculates the calibration by least squares fitting, according to the data given in the calibration library.

The maximum number of the calibration peaks is 10, so C and E are given as C(10) and E(10).

6. Subroutine ANALYS

SUBROUTINE ANALYS(NLIB,NCH,III,JS)

NLIB : number of peaks of the oriented library which
is prepared for each particular problem.
NCH : as defined in FIRST
III : as defined in FIRST
JS : as defined in FIRST

Subroutine ANALYS analyses all the spectra except that which belongs to calibration library, according to the information given in the oriented library in each interval calculated from the FWHM value given in that library. The program also performs the calculation of the concentration of the elements present in the sample in terms of ppm.

7. Subroutine PKF

SUBROUTINE PKF (JJJ,III,IRIG)

DIMENSION SIGMA(50)

DIMENSION IND(50)

JJJ : number of channels of the spectrum
III : number of spectrum to be analyzed
IRIG : number of write statements used
SIGMA : standard deviation for each element
IND(JMAX): indexes indicating the kinds of the peaks
found by the routine.

The search of the peaks based on the change of the sign of the first derivative of the spectrum is carried out by this routine, PKF. It also searches the possible double peaks in the spectrum. The mathematical tests used for that purpose are given in the subchapter 3.3.3 in detail.

The maximum number of elements of the oriented library is given as 50 so it is used as dimensions for SIGMA and IND in the routine.

8. Subroutine SURF

```
SUBROUTINE SURF (NCH,JS)
```

```
DIMENSION POUNT (4096)
```

```
NCH      : defined in FIRST
```

```
JS       : defined in FIRST
```

```
POUNT    : number of counts detected in each channel
```

This routine computes the area of the peaks found by the routine PKF. The method for the computation of areas is discussed in the subchapter 3.3.4.

9. Subroutine SURF2

```
SUBROUTINE SURF2 (IRIG)
```

```
DIMENSION SIGMA (50)
```

```
IRIG     : defined in PKF
```

```
SIGMA    : defined in the same routine.
```

Subroutine SURF2 computes the areas of doublet when they are recognized by the routine PKF.

10. Subroutine SENSIT

```
SUBROUTINE SENSIT (RAD)
```

```
RAD      : minimum detectable effect observed.
```

If no peaks are observed in the interval indicated by the oriented library, the minimum counting rates that could have had a photopeak over the existing background are calculated by this routine, SENSIT.

11. Subroutine PRINT

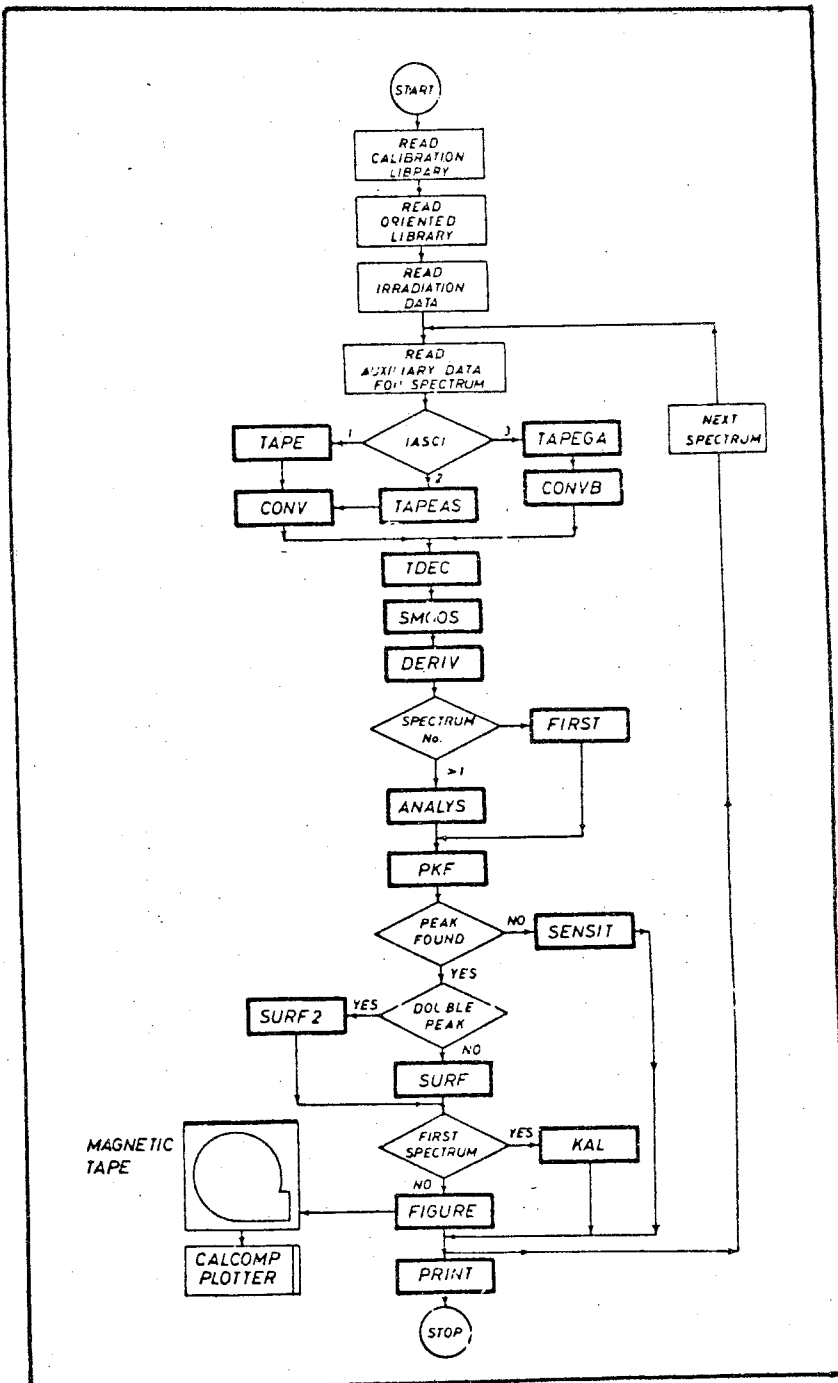
SUBROUTINE PRINT (IWR, IRIG)

IWR : an integer indicating what are printed
according to the routines

IRIG : it is defined in PKF

This routine is used for the printing of the results
obtained.

APPENDIX H



Flow chart for the GRETEL code

APPENDIX I

LISTING OF THE UNIVAC VERSION OF
GRETEL

CON QUESTO PROGRAMMA SI TROVANO I PICCHI DOPPI

REAL*8 EL1,ELEM
REAL*8 EL2,EL3,EL4,EL2D

COMMON EL1(10),EN1(10),ICAN(10),INT(10)
COMMON ELE(50),EN(50),FWT(50),XLAM(50),AA(50)
COMMON COUNT(4096),SPSM(4096),DER(4096)
COMMON ISTAR,IFID,INDEX,JMAX,FMEDIO
COMMON AREA(50),FPR(50),FWHM(50)
COMMON LS(50),CENTR(50),LD(50),PMOL(50)
COMMON DECAY(10),VALA(10),RESO(10)
COMMON TDEC1,TIRR,TCON,ALFA,BETA,FLUX,WEIGHT
COMMON STAND(6,50),NELST,IREL

COMMON/ET1/ITRR,IO(5),MO(5),ANOME(5),ASAMPL(3)
COMMON /ET4/ FL4(10),CKK(10),VKK(10),VMAX(10),MPK
COMMON /ET5/ FL3(50),VS(50),P2(50),CS(50),VS2(50),VD2(50),CLS2(50)
1,CLD2(50),ISEHS
COMMON /ET6/ EL2(50),VC(50),P1(50),CC(50),VS1(50),VD1(50),CLS1(50)
1,CLD1(50),IPICK
COMMON/ET8/EDEC,FSAT,DFENS
COMMON/ET9/EL2D(50),VCD(50),P1D(50),CCD(50),V1D(50),CL1D(50),IPICD
COMMON/ET10/EL2REL(50),EL3REL(50),FL2DPE(50)

DIMENSION IVE(12)
DATA STAN/,STAN,/,

FLUX=0.
DENS=0.
DO 777 J=1,50
AA(J)=0.
PMOL(J)=0.

777 CONTINUE
READ(5,999) (IVE(I),I=1,12)

READ(5,1000) LIBST

DO 9 K=1,LIBST

9 READ(5,998) EL1(K),EN1(K),ICAN(K),INT(K),DECAY(K),VALA(K),RESO(K)

READ(5,1000) NSPT,NCH,NSMO,IPDO,IASCI,IREL

IF(IREL) 200,200,201

201 READ(5,1000) NELST

STAND(1,K)=SIMBOLO ELEMENTO (4 LETTERE)
STAND(2,K)=ENERGIA PICCO IN KEV
STAND(3,K)=RISOLUZIONE RIVELATORE IN KEV
STAND(4,K)=COSTANTE DECADIMENTO 0.693/TDEC IN MINUTI

. GRE0002
. GRE0003
. GRE0004
. GRE0005
. GRE0006
. GRE0007
. GRE0008
. GRE0009
. GRE0010
. GRE0011
. GRE0012
. GRE0013
. GRE0014
. GRE0015
. GRE0016
. GRE0017
. GRE0018
. GRE0019
. GRE0020
. GRE0021
. GRE0022
. GRE0023
. GRE0024
. GRE0025
. GRE0026
. GRE0027
. GRE0028
. GRE0029
. GRE0030

. GRE0033
. GRE0034
. GRE0035
. GRE0036
. GRE0037
. GRE0038
. GRE0039

. GRE0043
. GRE0045
. GRE0046
. GRE0047
. GRE0048
. GRE0049
. GRE0050
. GRE0051
. GRE0052
. GRE0053
. GRE0054

53.	C	STAND(5,K)=PRM DI ELEMENTO DELLO STANDARD	GRE0055
54.	C	STAND(6,K)=ARFAPICCO AL TEMPO ZERO/MICROGRAMMI ELEMENTO (CALCOLA-	GRE0056
55.	C	TO DAL PROGRAMMA E MESSO IN LIBRERIA DELLO STANDARD)	GRE0057
56.	C	*****	GRE0058
57.	C	*****	GRE0059
58.	C	DO 202 K=1,NELST	GRE0060
59.	C	202 READ (5,1030) STAND(1,K),STAND(2,K),STAND(3,K),STAND(4,K),STAND(5,	GRE0061
60.	C	1K)	GRE0062
61.	C	READ(5,1031) (IQ(K),K=1,5),NIRR,TIPR,(ANOME(K),K=1,5)	GRE0063
62.	C	GO TO 500	GRE0064
63.	C	200 READ (5,1040) NLIB	GRE0065
64.	C	DO 10 J=1,NLIB	GRE0066
65.	C	10 READ (5,1001) ELEM(J),EN(J),FWT(J),XLAM(J),AA(J),PMOL(J)	GRE0067
66.	C	READ (5,1003) FLUX,TIRR,(IQ(K),K=1,5),NIRR,(ANOME(K),K=1,5),DENS	GRE0068
67.	C		GRE0069
68.	C	WRITE(6,1014) NSPT	GRE0070
69.	C	IF(IASCI-2) 70,80,60	GRE0071
70.	C	60 WRITE (6,1012) IASCI	GRE0072
71.	C	GO TO 90	GRE0073
72.	C	70 WRITE(6,1015) IASCI	GRE0074
73.	C	GO TO 90	GRE0075
74.	C	80 WRITE(6,1016) IASCI	GRE0076
75.	C	90 WRITE (6,1013)	GRE0077
76.	C	DO 50 K=1,NLIB,2	GRE0078
77.	C	K1=K+1	GRE0079
78.	C	50 WRITE (6,1019) ELEM(K),EN(K),FWT(K),XLAM(K),AA(K),PMOL(K),ELEM(K1)	GRE0080
79.	C	1,EN(K1),FWT(K1),XLAM(K1),AA(K1),PMOL(K1))	GRE0081
80.	C		GRE0082
81.	C		GRE0083
82.	C	500 TCON=1,	GRE0084
83.	C	DO 1 I=1,NSPT	GRE0085
84.	C	I1=I	GRE0086
85.	C	IF(IPDO.EQ.1) GO TO 444	GRE0087
86.	C	ISTAR=1	GRE0088
87.	C	TEND=NCH	GRE0089
88.	C	IF(IASCI-2) 20,30,33	GRE0090
89.	C		GRE0091
90.	C	*** LETTURA CON CODICE ASCII NO PARITY ***	GRE0092
91.	C		GRE0093
92.	C	33 READ(5,1008) (MO(K),K=1,5),(ASAMPL(K),K=1,3),WEIGHT,TMIS,ITR,ICALC,	GRE0094
93.	C	1JS	GRE0095
94.	C	IF(ASAMPL(1).NE.STAN) GO TO 7770	GRE0096
95.	C	DO 7771 IJK=1,NELST	GRE0097
96.	C	7771 STAND(5,IJK)=WEIGHT*STAND(5,IJK)	GRE0098
97.	C	7770 CALL TAPEAS(NCH,COUNT)	GRE0099
98.	C	TCON=TMIS	GRE0100
99.	C	GO TO 40	GRE0101
100.	C		GRE0102
101.	C	*** LETTURA CON CODICE IBM ***	GRE0103
102.	C		GRE0104
103.	C	20 READ(5,1004) (MO(K),K=1,5),(ASAMPL(K),K=1,3),WEIGHT,ITR,ICALC,JS	GRE0105
104.	C	IF(ASAMPL(1).NE.STAN) GO TO 7772	GRE0106
105.	C	DO 7773 IJK=1,NELST	GRE0107
106.	C	7773 STAND(5,IJK)=WEIGHT*STAND(5,IJK)	GRE0108
107.	C	7772 CALL TAPEAS(NCH,COUNT)	GRE0109
108.	C	IF(COUNT(1).LE.1.) GO TO 40	GRE0110
109.	C	TCON=COUNT(1)/240.	GRE0111
110.	C	GO TO 40	GRE0112
111.	C		GRE0113
112.	C	*** LETTURA CON CODICE ASCII ***	GRE0114
113.	C		GRE0115
114.	C	30 READ(5,1004) (MO(K),K=1,5),(ASAMPL(K),K=1,3),WEIGHT,ITR,ICALC,JS	GRE0116
115.	C	IF(ASAMPL(1).NE.STAN) GO TO 7774	GRE0117
116.	C	DO 7775 IJK=1,NELST	GRE0118
117.	C	7775 STAND(5,IJK)=WEIGHT*STAND(5,IJK)	GRE0119
118.	C	7774 CALL TAPEAS(NCH,COUNT)	GRE0120
119.	C	IF(COUNT(1).LE.1.) GO TO 40	GRE0121
120.	C	TCON=COUNT(1)/60.	GRE0122
121.	C	GO TO 40	GRE0123
122.	C		GRE0124
123.	C	444 GO TO 40	GRE0125
124.	C	1,ITR,ICALC,JS	GRE0126
125.	C	IF(ASAMPL(1).NE.STAN) GO TO 7776	GRE0127
126.	C	DO 7777 IJK=1,NELST	GRE0128
127.	C	7777 STAND(5,IJK)=WEIGHT*STAND(5,IJK)	GRE0129
128.	C	7776 NCH=TEND-ISTAR+1	GRE0130
129.	C	IF(IASCI-2) 445,446,447	GRE0131
130.	C		GRE0132
131.	C	*** LETTURA CON CODICE ASCII NO PARITY ***	GRE0133
132.	C		GRE0134
133.	C	447 CALL TAPEAS(NCH,COUNT)	GRE0135
134.	C		GRE0136


```

217. 1016 FORMAT(1H+,45X,,PERFORATI CON CODICE ASCII,,) . GPF0224
218. 1017 FORMAT(1H0,5X,,NON CI SONO PICCHI DI CALIBRAZIONE,,) . GPF0225
219. 1018 FORMAT(1H0,5X,,LIBERTA,///,ELEM,,5X,,EN(KEV) RES(KEV) DEC(MIN) . GPF0226
220. 1,,3X,,VALORE A,,3X,,P.MOL,,5X,,ELEM,,5X,,EN(KEV) PFS(KEV) DEC(M . GPF0227
221. 2IN,,3X,,VALORE A,,3X,,P.MOL,///) . GPF0228
222. 1019 FORMAT(1H0,A3,2X,F7.2,F7.3,3X,1PE9.3,2X,E9.3,2X,0PF7.3,5X,A8,2X,F . GPF0229
223. 17.2,F7.3,3X,1PE9.3,2X,E9.3,2X,0PF7.3) . GPF0230
224. 1020 FORMAT(1H1,SPETTRO N,,15,,CANALI,,15,,TCON,F10.3,,MINUTI . GPF0231
225. 1 DATA,,5I4,,DECAD,,F12.3,,MIN,,14X,,CALCOMP,/) . GPF0232
226. 1030 FORMAT(A4,2E7.2,2E9.3) . GPF0233
227. 1031 FORMAT(5I4,15,F10.3,5A4) . GPF0234
228. GO TO 334 . GPF0235
229. 333 WRITE (6,1017) . GPF0236
230. 334 STOP . GPF0237
231. END . GPF0238

```

IN 405 IBANK 330 DBANK 14898 COMMON

```

1. FUNCTION TDEC(M0,I0,IME)
2. DIMENSION M0(5),I0(5),IME(12)
3. I1=M0(2)
4. I2=I0(2)
5. TDEC=(( (M0(1)-I0(1))*365+IME(I1)-IME(I2)+M0(3)-I0(3))*24+M0(4)-I0(
6. 14)) *60+M0(5)-I0(5)
7. RETURN
8. END

```

TN 52 IBANK 18 DBANK

```

1.  SPECTRA,KAL
2.  SUBROUTINE KAL(N,C,E,A,B)
3.  DIMENSION C(10),E(10)
4.  C
5.  FN=FN
6.  C1=0.
7.  E1=0.
8.  C2=0.
9.  CE=0.
10. DO 1 J=1,N
11.   C1=C1+C(J)
12.   E1=E1+E(J)
13.   C2=C2+C(J)**2
14.   CE=CE+C(J)*E(J)
15. 1 CONTINUE
16.  DEN=FN*C2-C1**2
17.  A=(FN*CE-C1*E1)/DEN
18.  B=(C2*E1-C1*CE)/DEN
19.  RETURN
20. END

```

ETN 85 IBANK 26_DBANK

```

1. SUBROUTINE TAPE(NPT,Y)
2. DIMENSION Y(4096)
3. INTEGER A(64)
4. J=0
5. 2 CONTINUE
6. IF (J.EQ.NPT) GO TO 3
7. READ(11,1,END=4) (A(I),I=1,60)
8. 1 FORMAT (6X,60A1)
9. J=J+1
10. Y(J)=CONV(A,1,5)
11. IF(Y(J).NE.0.) GO TO 2
12. Y(J)=1.
13. GO TO 2
14. 3 NPT=J
15. RETURN
16. 4 WRITE(6,5)
17. 5 FORMAT (1H1,**, END OF FILE ON PAPER TAPE **, )
18. STOP
19. END

```

FTN 48 IBANK 103 DBANK

```

      S. SPECTRA.TAPEAS
      BR1 *06/16/82-22:12(7,)
          1. SUBROUTINE TAPEA'(UNIT,Y)
          2. DIMENSION Y(4096)
          3. INTEGER A(64)
          4. INTEGER UNO/,1,/-
          5. J=0

```

```

6.      22 CONTINUE
7.      IF (J.EQ.NPT) GO TO 3
8.      READ (11,1,END=4) K,(A(I),I=1,60)
9.      1 FORMAT (3X,I3,60A1)
10.     IF (K-57) 2,101,102
11.     2 IF (K-1) 22,22,103
12.     103 DO 100 I=1,60
13.     A(I)=UNO
14.     100 CONTINUE
15.     101 DO 5 K=1,8
16.     J=J+1
17.     L1=K*7-6
18.     L2=L1+5
19.     Y(J)=CONV(A,L1,L2)
20.     IF (Y(J).NE.0.) GO TO 5
21.     Y(J)=1.
22.     5 CONTINUE
23.     GOTO 22
24.     102 DO 6 K=1,8
25.     J=J+1
26.     L1=K*7-5
27.     L2=L1+5
28.     Y(J)=CONV(A,L1,L2)
29.     IF (Y(J).NE.0.) GO TO 6
30.     Y(J)=1.
31.     6 CONTINUE
32.     GOTO 22
33.     3 NPT=J
34.     RETURN
35.     4 WRITE (6,15)
36.     15 FORMAT (1H1,**END OF FILE ON PAPER TAPE**)
37.     STOP
38.     END

```

```

TN 117 IBANK 116 DBANK
SPECTRA.TAPEGA
*06/16/82-22:12(3,)
SUBROUTINE TAPEGA(NPT,Y)
1.  DIMENSION Y(4096)
2.  INTEGER A(64)
3.  INTEGER UNO/,1, /
4.  J=0
5.  22 CONTINUE
6.  IF (J.EQ.NPT) GO TO 3
7.  READ (11,1,END=4) K,(A(I),I=1,64)
8.  1 FORMAT (3X,I3,10X,64A1)
9.  IF (K-71) 2,101,101
10.  2 IF (K-12) 22,22,103
11.  103 DO 100 I=1,64
12.  A(I)=UNO
13.  100 CONTINUE
14.  101 DO 5 K=1,8
15.  J=J+1
16.  L1=K*8-7
17.  L2=L1+7
18.  Y(J)=CONVB(A,L1,L2)
19.  IF (Y(J).NE.0.) GO TO 5
20.  Y(J)=1.
21.  5 CONTINUE
22.  GO TO 22
23.  3 NPT=J
24.  RETURN
25.  4 WRITE (6,15)
26.  15 FORMAT (1H1,**END OF FILE ON PAPER TAPE**)
27.  STOP
28.  END
29.

```

```

30.  FUNCTION CONV(A,L1,L2)
31.  C
32.  INTEGER A(1),TAB(11)
33.  DATA TAB /,0,1,2,3,4,5,6,7,8,9, /
34.  *
35.  CONV=0.
36.  DO 2 I=L1,L2
37.  IF (A(I).EQ.TAB(11)) A(I)=TAB(1)

```

GPF0577
 GPF0578
 GPF0580
 GPF0581
 GPF0582
 GPF0583
 GPF0584
 GPF0585
 GPF0586
 GPF0587
 GPF0588
 GPF0589
 GPF0590
 GPF0591
 GPF0592
 GPF0593
 GPF0594
 GPF0595
 GPF0596
 GPF0597
 GPF0598
 GPF0599

GPF0615
 GPF0616
 GPF0617
 GPF0618
 GPF0619
 GPF0620
 GPF0621
 GPF0622
 GPF0623
 GPF0624
 GPF0625
 GPF0626
 GPF0627
 GPF0628
 GPF0629

GPF0630
 GPF0631
 GPF0632
 GPF0633
 GPF0634
 GPF0635
 GPF0636
 GPF0637
 GPF0638
 GPF0639

```

38. DO 1 J=1,10
39. IF(A(I).EQ.TAB(J)) GO TO 3
40. 1 CONTINUE
41. CONV=0.
42. RETURN
43. 3 FJ=J
44. CONV=CONV*10.+FJ-1.
45. 2 CONTINUE
46. RETURN
47. END

```

. GPF0648
 . GPF0649
 . GPF0650
 . GPF0651
 . GPF0652
 . GPF0653
 . GPF0654
 . GPF0655
 . GPF0656
 . GPF0657

```

48. C FUNCTION CONVR(A,L1,L2)
49.
50. INTEGER A(1),TAB(11)
51. DATA TAB /,0 ,,,1 ,,,2 ,,,3 ,,,4 ,,,5 ,,,6 ,,,7 ,,,8 ,,
52. *,
53. * CONVR=0.
54. DO 2 I=L1,L2
55. IF(A(I).EQ.TAB(11)) RETURN
56. DO 1 J=1,10
57. IF(A(I).EQ.TAB(J)) GO TO 3
58. 1 CONTINUE
59. CONVR=0
60. RETURN
61. 3 FJ=J
62. CONVR=CONVR*10.+FJ-1.
63. 2 CONTINUE
64. RETURN
65. END

```

. GPF0658
 . GPF0659
 . GPF0663
 . GPF0664
 . GPF0665
 . GPF0666
 . GPF0667
 . GPF0668
 . GPF0669
 . GPF0670
 . GPF0671
 . GPF0672
 . GPF0673
 . GPF0674
 . GPF0675

N 216 IBANK 164 DBANK
 SPECTRA.SMOOS

1 *06/16/82-22:12(2,)
 SUBROUTINE SMOOS(NSMO)

```

1. C
2. REAL*8 EL1,ELEM
3. COMMON EL1(10),EN1(10),ICAM(10),INT(10)
4. COMMON ELEM(50),FN(50),FWT(50),XLAM(50),AA(50)
5. COMMON COUNT(4096),SPSM(4096),DER(4096)
6. COMMON ISTAR,IEND,INDEX,JMAX,FMEDJO
7. COMMON AREA(50),FRP(50),FWHM(50)
8. COMMON LS(50),CENTR(50),LD(50),PMOL(50)
9. COMMON DECAY(10),VALA(10),RESO(10)
10. COMMON TDEC1,TIRP,TCON,ALFA,BETA,FLUX,WEIGHT
11.
12. C
13. C
14. IF(NSMO-7) 30,40,50
15. C
16. C *** SMOOTHING A 9 PUNTI ***
17. C
18. 50 N1=ISTAR+4
19. N2=IEND-4
20. C0=50.
21. C1=54.
22. C2=39.
23. C3=14.
24. C4=-21.
25. CNORM=231.
26. GO TO 60
27. C
28. C *** SMOOTHING A 7 PUNTI ***
29. C
30. 40 N1=ISTAR+3
31. N2=IEND-3
32. C0=7.
33. C1=6.
34. C2=3.
35. C3=-2.
36. C4=0.
37. CNORM=21.
38. GO TO 60
39. C
40. C *** SMOOTHING A 5 PUNTI ***
41. C
42. 30 N1=ISTAR+2

```

. GPF0678
 . GPF0679
 . GPF0680
 . GPF0681
 . GPF0682
 . GPF0683
 . GPF0684
 . GPF0685
 . GPF0686
 . GPF0687
 . GPF0688
 . GPF0689
 . GPF0690
 . GPF0691
 . GPF0692
 . GPF0693
 . GPF0694
 . GPF0695
 . GPF0696
 . GPF0697
 . GPF0698
 . GPF0699
 . GPF0700
 . GPF0701
 . GPF0702
 . GPF0703
 . GPF0704
 . GPF0705
 . GPF0706
 . GPF0707
 . GPF0708
 . GPF0709
 . GPF0710
 . GPF0711
 . GPF0712
 . GPF0713
 . GPF0714
 . GPF0715
 . GPF0716
 . GPF0717
 . GPF0718

```

43.      N2=IEND-2
44.      C0=17.
45.      C1=12.
46.      C2=-3.
47.      C3=0.
48.      C4=0.
49.      CNORM=35.
50. C
51. 60 DO 1 K=N1,N2
52.   SPSM(K)=C0*COUNT(K)+C1*(COUNT(K+1)+COUNT(K-1))+C2*(COUNT(K+2)+COUNT(K-2))+C3*(COUNT(K+3)+COUNT(K-3))+C4*(COUNT(K+4)+COUNT(K-4))
53.   1T(K-2))
54.   1 SPSM(K)=SPSM(K)/CNORM
55.   RETURN
56.   END

N 105 IBANK 39 DBANK 13030 COMMON
SPECTRA.DERIV
1 *06/16/82-22:12(2,)
1. SUBROUTINE DERIV(NSMO)
2. REAL*8 EL1,ELEM
3. COMMON EL1(10),EN1(10),ICAN(10),INT(10)
4. COMMON ELEM(50),FN(50),FWT(50),XLAM(50),AA(50)
5. COMMON COUNT(4096),SPSM(4096),DER(4096)
6. COMMON ISTAR,IEND,INDEX,JMAX,FMEDIO
7. COMMON AREA(50),FRR(50),FWHM(50)
8. COMMON LS(50),CENTR(50),LD(50),PMOL(50)
9. COMMON DECAY(10),VALA(10),PESO(10)
10. COMMON TDEC1,TIRR,TCON,ALFA,BETA,FLUX,WEIGHT
11. C
12. IF(NSMO-7) 10,20,30
13. 30 IS=ISTAR+5
14.   IE=IEND-5
15.   GO TO 40
16. 20 IS=ISTAR+4
17.   IE=IEND-4
18.   GO TO 40
19. 10 IS=ISTAR+3
20.   IE=IEND-3
21. C
22. 40 DO 1 I=IS,IE
23.   IP=I+1
24.   IN=I-1
25.   DER(I)=(SPSM(IP)-SPSM(IN))/2.
26.   1 CONTINUE
27.   RETURN
28.   END

N 56 IBANK 15 DBANK 13030 COMMON
SPECTRA.FIRST
1 *06/16/82-22:12(2,)
1. SUBROUTINE FIRST(LIRST,NCH,III,JS)
2. REAL*8 EL1,ELEM
3. REAL*8 EL4
4. COMMON EL1(10),EN1(10),ICAN(10),INT(10)
5. COMMON ELEM(50),FN(50),FWT(50),XLAM(50),AA(50)
6. COMMON COUNT(4096),SPSM(4096),DER(4096)
7. COMMON ISTAR,IEND,INDEX,JMAX,FMEDIO
8. COMMON AREA(50),FRR(50),FWHM(50)
9. COMMON LS(50),CENTR(50),LD(50),PMOL(50)
10. COMMON DECAY(10),VALA(10),PESO(10)
11. COMMON TDEC1,TIRR,TCON,ALFA,BETA,FLUX,WEIGHT
12. COMMON STAND(6,50),NELST,IPEL
13. COMMON /ET4/ EL4(10),CKK(10),VKK(10),VMAX(10),NPK
14. C
15. DIMENSION CK(10),VK(10),FCOR(10)
16. C
17. NPK=0
18. JUJ=0
19. CALL PKF(JUJ,III,IRIG)
20. IF(JMAX.LE.1) RETURN
21. CALL SURF(NCH,JS)
22. C
23. WRITE (6,1998)
24. WRITE (6,1999)
25. DO 3 J=1,JMAX
26.   DO 4 I=1,LIRST
27.     IDELTA=ICAN(I)-INT(I)
28.     KDELTA=ICAN(I)+INT(I)
29.     IF(CENTR(J)-IDELTA) 4,6,6
30.     IF(CENTR(J)-KDELTA) 7,7,4
31.     6 NPK=NPK+1
32.     CK(NPK)=CENTR(J)

```

```

33. VK(NPK)=EN1(I)
34. EL1(NPK)=EL1(I)
35. FDEC=EXP(DECAY(I)*TDEC1)
36. FSAT=1.-EXP(-DECAY(I)*TIRR)
37. CPM=AREA(J)/TCON
38. AZERO=CPM*FDEC
39. ASPER=0.
40. FCOR(NPK)=0.
41. IF (IREL.EQ.1) GO TO 100
42. ASPER=FLUX*FSAT*PESO(I)*1.E+06/AZERO
43. FCOR(NPK)=ASPER/VALA(I)
44. C
45. ***** I SEGUENTI VALORI VANNO NEL COMMON /ET4/ E SERVONO PER ,FIGUR
46. C
47. 100 IK=CK(NPK)+0.5
48. VMAX(NPK)=COUNT(IK)
49. EL4(NPK)=EL1(NPK)
50. CKK(NPK)=CK(NPK)
51. VKK(NPK)=VK(NPK)
52. C
53. WRITE (6,2000) NPK,EL1(NPK),VK(NPK),PESO(I),CK(NPK),CPM,AZERO,ASPE
54. 1R,VALA(I),FCOR(NPK)
55. 4 CONTINUE
56. 3 CONTINUE
57. IF (NPK.LE.1) RETURN
58. C
59. CALL KAL(NPK,CK,VK,ALFA,BETA)
60. FMEDIO=0.
61. DO 5 K=1,NPK
62. 5 FMEDIO=FMEDIO+FCOR(K)
63. FMEDIO=FMEDIO/NPK
64. WRITE (6,2002)
65. WRITE (6,2001) NPK,FMEDIO
66. C
67. 1998 FORMAT (1H0)
68. 1999 FORMAT (1H0,3X,,N,,,4X,,ELFM,,7X,,FN,,6X,,PESO (G),,5X,,CENTR,,5X,,
69. 1,C P M,,4X,,AREA ZERO,,3X,,A SPER,,5X,,A TEOP,,4X,,FATT CORR,////)
70. 2000 FORMAT (1H,,15,2X,A8,2X,F7.2,2X,F11.8,3X,F7.2,5(2X,1PE9.3))
71. 2001 FORMAT (1H,,5X,,VALORE MEDIO DEL FATTORE CORRETTIVO SU ,,14,, PICO
72. 1H1,,3X,,1PE9.3////)
73. 2002 FORMAT (1H0,93X,-----,////)
74. RETURN
75. END
N 198 IBANK 213 DBANK 13383 COMMON
S SPECTRA ANALYS
1 *06/16/82-22:12(2,)
1 SUBROUTINE ANALYS(NLIB,NCH,III,JS)
2 INTEGER H1,H2
3 REAL*8 EL1,ELEM
4 REAL*8 EL2,EL3,EL2D
5 C
6 COMMON EL1(10),EN1(10),ICAN(10),INT(10)
7 COMMON ELEM(50),EN(50),FWT(50),XLAM(50),AA(50)
8 COMMON COUNT(4096),SPSM(4096),DER(4096)
9 COMMON ISTAR,IEND,INDEX,JMAX,FMEDIO
10 COMMON AREA(50),ERR(50),FWHM(50)
11 COMMON LS(50),CENTR(50),LD(50),PMOL(50)
12 COMMON DECAY(10),VALA(10),PESO(10)
13 COMMON TDEC1,TIRR,TCON,ALFA,BETA,FLUX,WEIGHT
14 COMMON STAND(6,50),NELST,IPEL
15 COMMON FWHM
16 C
17 COMMON /ET1/ NIRR,IO(5),MO(5),ANOME(5),ASAMPL(3)
18 COMMON/ET2/PPM,ATOMI,PERC,AKEV,CPM,AREA0,UGM,LARG,H1,H2,J1
19 COMMON /ET5/ EL3(50),VS(50),P2(50),CS(50),VS2(50),VD2(50),CLS2(50)
20 1,CLD2(50),ISENS
21 COMMON /ET6/ EL2(50),VC(50),P1(50),CC(50),VS1(50),VD1(50),CLS1(50)
22 1,CLD1(50),IPICK
23 COMMON/ET7/C1,C2,JDOUR,KKK,IND2
24 COMMON/ET8/FDEC,FSAT,DFNS
25 COMMON/ET9/EL2D(50),VCD(50),P1D(50),CCD(50),V1D(50),CL1D(50),IPICD
26 COMMON/ET10/EL2REL(50),EL3REL(50),EL2DRE(50)
27 DATA STAN,STAN,/,/,/
28 DATA BLANK,/,/,/
29 C
30 IPICD=0
31 IPICK=0
32 ISENS=0
33 IRI=0
34 FA=2.
35 FB=1.2

```

CALL PRINT(3,IRIG)

```
C
200 IF (IREL.EQ.1) NLIR=NELST
    IF (ASAMPL(1).NE.STAN) GO TO 204
    DO 201 K=1,NELST
        XLAM(K)=STAND(4,K)
        EN(K)=STAND(2,K)
    201 FWT(K)=STAND(3,K)
    204 DO 3 J=1,NLIR
        JUJ=J
        J1=J
        FDEC=EXP(XLAM(J)*TDEC1)
        FSAT=1.-EXP(-XLAM(J)*TIRR)
        ENCH=(EN(J)-BETA)/ALFA
        FWHT=FWT(J)/ALFA
        ISTAR=(ENCH-FWHT*FA)
        IEND=(ENCH+FWHT*FA)
        H1=(ENCH-FWHT*FB)-0.5
        H2=(ENCH+FWHT*FB)+0.5
        L=1
        IF (ISTAR.GT.L.AND.IEND.LT.NCH) GO TO 7
```

CALL PRINT(4,IRIG)

```
C
C
GO TO 3
7 CALL PKF(JUJ,I11,IRIG)
IF (JMAX.GT.0) GO TO 9
IF (JDOUR.GT.0) GO TO 3
IF (ASAMPL(1).NE.STAN) GO TO 10
STAND(6,J)=BLANK
PPM=BLANK
GO TO 301
10 CALL SENSIT(RAD)
SENS=RAD/TCON
AREA0=SENS*FDEC
IF (ASAMPL(1).EQ.STAN) STAND(6,J)=AREA0/STAND(5,J)
IF (IREL.EQ.1) GO TO 300
UGM=AA(J)*AREA0/(FLUX*FSAT)
PPM=UGM/WEIGHT
ATOMI=6.02E17*UGM*DENS/(PMOL(J)*WEIGHT)
GO TO 301
```

```
C
300 IF (STAND(6,J).NE.BLANK) GO TO 299
PPM=BLANK
GO TO 301
```

```
299 PPM=(1./STAND(6,J))*(AREA0/WEIGHT)
301 CONTINUE
```

***** I SEGUENTI VALORI VANNO NEL COMMON /ET5/ E SERVONO PER ,FIGUR

```
C
IIN=(H1+(H2-H1)/2)
ISENS=ISENS+1
VS(ISENS)=COUNT(IIN)
P2(ISENS)=EN(J)
CS(ISENS)=IIN+0.5
IF (IREL.EQ.1) GO TO 400
EL3(ISENS)=ELEM(J)
GO TO 401
400 EL3REL(ISENS)=STAND(1,J)
401 VS2(ISENS)=H1+0.5
    VD2(ISENS)=H2+0.5
    CLS2(ISENS)=COUNT(H1)
    CLD2(ISENS)=COUNT(H2)
```

CALL PRINT(5,IRIG)

```
C
GO TO 3
9 CALL SURF(NCH,JS)
DO 1 K=1,JMAX
    IF (AREA(K).GT.0) GO TO 8
    GO TO 10
8 CPM=AREA(K)/TCON
    ERR1=ERR(K)/TCON
    PERC=ERR1*100./CPM
    AREA0=CPM*FDEC
    IF (ASAMPL(1).EQ.STAN) STAND(6,J)=AREA0/STAND(5,J)
    IF (IREL.EQ.1) GO TO 302
    UGM=AA(J)*AREA0/(FLUX*FSAT)
    PPM=UGM/WEIGHT
    ATOMI=6.02E17*UGM*DENS/(PMOL(J)*WEIGHT)
    GO TO 303
302 IF (STAND(6,J).NE.BLANK) GO TO 298
    PPM=BLANK
```

```

118. 298 GO TO 303
119. 303 PPM=(1./STAND(6,J))*(AREA0/WEIGHT)
120. CONTINUE
121. AKEV=CENTR(K)*ALFA+BETA
122. ENDIFF=EN(J)-AKEV
123. H1=LS(K)
124. H2=LD(K)
125. LARG=LD(K)-LS(K)+1
126. TK=CENTR(K)+0.5
127. C C
128. C ***** I SEGUENTI VALORI VANNO NEL COMMON /ET6/ E SERVONO PER FIGUR
129. C
130. IPICK=IPICK+1
131. IF(IREL.EQ.1) GO TO 402
132. EL2(IPICK)=ELEM(J)
133. GO TO 403
134. 402 EL2REL(IPICK)=STAND(1,J)
135. 403 CONTINUE
136. VC(IPICK)=COUNT(TK)
137. P1(IPICK)=AKEV
138. CC(IPICK)=CENTR(K)
139. VS1(IPICK)=LS(K)
140. VD1(IPICK)=LD(K)
141. CLS1(IPICK)=COUNT(H1)
142. CLD1(IPICK)=COUNT(H2)
143. C
144. IF(PERC.GT.50.) GO TO 40
145. C
146. CALL PRINT(9,IRIG)
147. C
148. GO TO 1
149. C
150. 40 CALL PRINT(10,IRIG)
151. 1 CONTINUE
152. 3 CONTINUE
153. C
154. IF(ASAMPL(1).NE.STAN) GO TO 202
155. WRITE(6,2001)
156. DO 3000 K=1,NELST
157. IF(STAND(6,K).NE.BLANK) GO TO 297
158. WRITE(6,2002) STAND(1,K),STAND(2,K),STAND(3,K),STAND(4,K),STAND(5
159. 1,K),STAND(6,K)
160. GO TO 3000
161. 297 WRITE(6,2000) STAND(1,K),STAND(2,K),STAND(3,K),STAND(4,K),STAND(5
162. 1,K),STAND(6,K)
163. 3000 CONTINUE
164. C
165. 2000 FORMAT(1H0,A4,6X,F7.2,F7.3,3X,1PE9.3,2(2X,E9.3))
166. 2001 FORMAT(1H1,11BRERIA, METODO, RELATIVO,///,ELEM,,5X,,EN(KEV) RES
167. 1(KEV) DEC(MIN),,5X,,UGM,,6X,,CPMO/UGM,/)
168. 2002 FORMAT(1H0,A4,6X,F7.2,F7.3,3X,1PE9.3,2X,E9.3,7X,A4)
169. 202 RETURN
170. END
IN 485 IBANK 211 DBANK 14774 COMMON
SPECTRA.PKF
1 *06/16/82-22:12(2,)
1 SUBROUTINE PKF(JJJ,III,IRIG)
2 REAL*8 EL1,ELEM,EL2D
3 INTEGER H1,H2
4 COMMON EL1(10),EN1(10),ICAN(10),INT(10)
5 COMMON ELEM(50),EN(50),FWT(50),XLAM(50),AA(50)
6 COMMON COUNT(4096),SPSM(4096),DER(4096)
7 COMMON ISTAR,IFND,INDEX,JMAX,EMED10
8 COMMON AREA(50),ERR(50),FWHM(50)
9 COMMON LS(50),CENTR(50),LD(50),PMOL(50)
10 COMMON DECAY(10),VALA(10),PES0(10)
11 COMMON TDEC1,TIRP,TCON,ALFA,BETA,FLUX,WEIGHT
12 COMMON STAND(6,50),NELST,IREL
13 COMMON FWHT
14 COMMON/ET2/PPM,ATOMI,PERC,AKEV,CPM,AREA0,UGM,LARG,H1,H2,J1
15 COMMON/ET3/ENERG(50)
16 COMMON/ET7/C1,C2,JDOUR,KKK,IND2
17 COMMON/ET8/EDEC,FSAT,DENS
18 COMMON/ET9/EL2D(50),VCD(50),P1D(50),CCD(50),V1D(50),CL1D(50),TPICD
19 DIMENSION SIGMA(50)
20 DIMENSION IND(50)
21 C
22 C
23 J1=JJJ
24 JDOUR=0
25 ITOUR=0

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26. ITOUPE=0
27. 12 CONF=1.5
28. I=ISTAR
29. JMAX=0
30. 10 IL=I
31. ISW=-1
32. 1 I=I+1
33. IF(I.GE.IEND) GOTO 6
34. IF(DEF(I)) 3,3,2
35. 2 IF(ISW.GT.0) GO TO 1
36. ISW=1
37. GO TO 1
38. 3 IF(ISW.LT.0) GO TO 1
39. ISW=-1
40. AMAX=-10.E+30
41. AMIN=+10.E+30
42. J=I-1
43. 7 IF(DEF(J).LE.AMAX) GO TO 4
44. AMAX=DEF(J)
45. JL=J
46. J=J-1
47. GO TO 7
48. 4 CONTINUE
49. J=I
50. 8 IF(DEF(J).GE.AMIN) GO TO 5
51. AMIN=DEF(J)
52. JR=J
53. J=J+1
54. GO TO 8
55. 5 CONTINUE
56. A1=SPSM(JL+1)+SPSM(JL-1)
57. A2=SPSM(JR+1)+SPSM(JR-1)
58. IF(A1.LE.0.OR.A2.LE.0) GO TO 10
59. SDEV1=CONF*SQRT(A1)/2.0
60. SDEV2=CONF*SQRT(A2)/2.0
61. AMIN=ABS(AMIN)
62. IF(AMIN.EQ.0.OR.AMAX.EQ.0) GOTO 10
63. VALMAX=AMAX1(AMAX,AMIN)
64. VALMIN=AMIN1(AMAX,AMIN)
65. RATIO=VALMAX/VALMIN
66. IF(AMAX.LE.SDEV1.AND.AMIN.LE.SDEV2) GO TO 10
67. JMAX=JMAX+1
68. C
69. C POUR CHAQUE PIC RECHERCHE DE L'INDICE IND(JMAX)
70. C
71. C IF((AMAX.GT.SDEV1.AND.AMIN.GT.SDEV2).AND.(RATIO.LE.1.5)) GOTO 97
72. C
73. C IF(AMAX.GT.SDEV1.AND.AMAX.GE.AMIN) GOTO 98
74. C
75. C IF(AMIN.GT.SDEV2.AND.AMIN.GE.AMAX) GOTO 99
76. C
77. 97 IND(JMAX)=0
78. GOTO 100
79. C
80. 98 IND(JMAX)=1
81. GOTO 100
82. C
83. 99 IND(JMAX)=2
84. C
85. C CALCUL DES PARAMETRES
86. C
87. C 100 FL=JL+(DEF(JL-1)-DEF(JL+1))*0.5/(DEF(JL+1)-2.*DEF(JL)+DEF(JL-1))
88. FR=JR+(DEF(JR-1)-DEF(JR+1))*0.5/(DEF(JR+1)-2.*DEF(JR)+DEF(JR-1))
89. IM=I-1
90. CENTR(JMAX)=IM-DEF(IM)/(DEF(I)-DEF(IM))
91. CENTR1=CENTR(JMAX)
92. SIGMA(JMAX)=(FR-FL)/2.0
93. SIGMA1=SIGMA(JMAX)
94. FWHM(JMAX)=SIGMA(JMAX)*2.355
95. FWHM1=FWHM(JMAX)
96. INDICE=IND(JMAX)
97. I=JR
98. C
99. C VALABLE POUR LE PREMIER SPECTRE DE CALIBRATION
100. C
101. C EXCLUSION DES PICS DOUBLETS
102. C
103. IF(III-1) 102,103,102
104. IF(IND(JMAX).EQ.1.OR.IND(JMAX).EQ.2) JMAX=JMAX-1
105. IF(JMAX.LT.0) JMAX=0
106. GOTO 10
107. 102 ENERG(JMAX)=CENTR(JMAX)*ALFA+BETA
ENERG1=ENERG(JMAX)

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 GRE1200

108.		IF (IND(JMAX).EQ.2) GOTO 400	GRE1117
109.		IF (IND(JMAX).EQ.1) GO TO 900	GRE1118
110.		IF (JMAX.EQ.1) GO TO 11	GRE1119
111.	C		GRE1120
112.	C	RECHERCHE SI LE PIC AVEC IND(JMAX)=0 FAIT PARTIE D'UN DOUBLET	GRE1121
113.			GRE1122
114.		IF (IND(JMAX-1).EQ.1) GO TO 200	GRE1123
115.	11	IF (JMAX-50) 10,6,6	GRE1124
116.	200	IF ((CENTR(JMAX)-CENTR(JMAX-1)).GE.(2.0 * FWHT)) GOTO 300	GRE1125
117.		C1=CENTR(JMAX-1)	GRE1126
118.		C2=ISTAR	GRE1127
119.		JDOUR=1	GRE1128
120.		CALL SURF2(IRIG)	GRE1129
121.		GOTO 301	GRE1130
122.	300	IF (ITOUR.EQ.1) GOTO 800	GRE1131
123.	301	JMAX=JMAX-1	GRE1132
124.		CENTR(JMAX)=CENTR1	GRE1133
125.		FWHM(JMAX)=FWHM1	GRE1134
126.		SIGMA(JMAX)=SIGMA1	GRE1135
127.		IND(JMAX)=INDICE	GRE1136
128.		ENERG(JMAX)=ENERG1	GRE1137
129.		IF (ITOUR.EQ.1) GOTO 60	GRE1138
130.		GOTO 10	GRE1139
131.	C		GRE1140
132.	C	ETUDE DU PIC AVEC IND(JMAX)=2	GRE1141
133.	C	RETOUR EN ARRIERE DU CANAL DE DEPART DE L'INTERVALLE	GRE1142
134.	C		GRE1143
135.	400	IF (JMAX.GT.1) GOTO 500	GRE1144
136.		IF (ITOUR.GT.0) GOTO 303	GRE1145
137.		ISTAR=CENTR1-5*FWHT	GRE1146
138.		ITOUR=ITOUR+1	GRE1147
139.		GOTO 12	GRE1148
140.	C		GRE1149
141.	C	RECHERCHE DU PIC DOUBLET	GRE1150
142.	C	1) FORME UN PIC DOUBLET AVEC PIC NORMAL IND(JMAX)=0	GRE1151
143.	C		GRE1152
144.	500	IF (IND(JMAX-1).EQ.1) GO TO 700	GRE1153
145.		IF ((CENTR(JMAX)-CENTR(JMAX-1)).GE.(2.0 * FWHT)) GOTO 310	GRE1154
146.		C2=CENTR(JMAX)	GRE1155
147.		C1=IFND	GRE1156
148.		JDOUR=2	GRE1157
149.		CALL SURF2(IRIG)	GRE1158
150.	302	IF (ITOUR.NE.1) GOTO 303	GRE1159
151.		JMAX=JMAX-1	GRE1160
152.		ITOUR=ITOUR+1	GRE1161
153.		GOTO 10	GRE1162
154.	310	IF (ITOUR.NE.1) GOTO 303	GRE1163
155.		JMAX=JMAX-2	GRE1164
156.		ITOUR=ITOUR+1	GRE1165
157.		GOTO 10	GRE1166
158.	303	JMAX=JMAX-1	GRE1167
159.		GO TO 10	GRE1168
160.	C		GRE1169
161.	C	2) FORME UN PIC DOUBLET AVEC PIC IND(JMAX)=1	GRE1170
162.	C		GRE1171
163.	700	IF ((CENTR(JMAX)-CENTR(JMAX-1)).GE.(2.0 * FWHT)) GOTO 801	GRE1172
164.		C1=CENTR(JMAX-1)	GRE1173
165.		C2=CENTR(JMAX)	GRE1174
166.		JDOUR=3	GRE1175
167.		CALL SURF2(IRIG)	GRE1176
168.	801	IF (ITOUR.NE.1) GOTO 800	GRE1177
169.		JMAX=0	GRE1178
170.		ITOUR=ITOUR+1	GRE1179
171.		GOTO 10	GRE1180
172.	800	JMAX=JMAX-2	GRE1181
173.		IF (ITOUR.EQ.1) GOTO 60	GRE1182
174.		GO TO 10	GRE1183
175.	C		GRE1184
176.	C	ETUDE DU PIC AVEC IND(JMAX)=1	GRE1185
177.	C		GRE1186
178.	900	IF (JMAX.EQ.1) GO TO 10	GRE1187
179.		IF (IND(JMAX-1).EQ.1) GO TO 1000	GRE1188
180.		GO TO 11	GRE1189
181.	1000	JMAX=JMAX-1	GRE1190
182.		CENTR(JMAX)=CENTR1	GRE1191
183.		FWHM(JMAX)=FWHM1	GRE1192
184.		SIGMA(JMAX)=SIGMA1	GRE1193
185.		IND(JMAX)=INDICE	GRE1194
186.		ENERG(JMAX)=ENERG1	GRE1195
187.		GOTO 10	GRE1196
188.	6	IF (JMAX.EQ.0) GOTO 60	GRE1197
189.		IF (IND(JMAX)-1) 60,62,60	GRE1198

```
IF(ITOURE.GT.0) GOTO 61
IEND=CENTR1+5*FWHT
ITOURE=ITOURE+1
I=JP
GOTO 10
JMAX=JMAX-1
CONTINUE
RETURN
END
```

```

:13(2.)
SUBROUTINE SURF (UCH,JS)
REAL*8 EL1,ELEM
COMMON EL1(10),EN1(10),ICAN(10),INT(10)
COMMON ELEM(50),EN(50),FWT(50),XLAM(50),AA(50)
COMMON COUNT(4096),SPSM(4096),DER(4096)
COMMON ISTAR,ITND,INDEX,JMAX,FMEDIO
COMMON AREA(50),ERR(50),FWHM(50)
COMMON LS(50),CENTR(50),LD(50),PMOL(50)
COMMON DECAY(10),VALA(10),PESO(10)
COMMON TDEC1,TIRR,TCON,ALFA,BETA,FLUX,WEIGHT
DIMENSION BOUNT(4096)

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IF(JS.EQ.1) GO TO 20
DO 30 J=1,NCH
  POINT(J)=COUNT(J)
GO TO 40

DO 10 I=1,NCH
  POINT(I)=SPSM(I)

DO 1 J=1,JMAX
  IC=CENTR(J)+0.5
  IL=IC-1
  IR=IC+1
  W=SORT(POINT(IC))
  W=POINT(IC)-W
  IND=0
  AR=POINT(IC)
  IF(POINT(IC-1).LT.W) GO TO 2
  IND=1
  IF(POINT(IC+1).LT.W) GO TO 3
  IF(IND.EQ.1) GO TO 4
  AR=AR+POINT(IC+1)
  IR=IR+1
  IF(IND.EQ.0) GO TO 4
  AR=AR+POINT(IC-1)
  IL=IL-1
  CONTINUE
  A=POINT(IL)+POINT(IR)
  INTV=CENTR(J)+FWHM(J)*2.1231+0.5
  DO 5 I=IC,INTV
    AR=AR+A
    IR=IR+1
    IL=IL-1
    A1=POINT(IR)+POINT(IL)
    IF(A1.GT.A) GO TO 6
    A=A1
  CONTINUE
  HM=IR-IL-1
  A2=(POINT(IR-1)+POINT(IL+1))/2.
  FONDO=A2*HM
  LD(J)=IR-1
  LS(J)=IL+1
  AR=A(J)=AR-FONDO
  ERR(J)=SORT(AR+HM*FONDO/2.0)

CONTINUE
RETURN
END

```

```

22:13(4, )
SUBROUTINE SURF2(IRIG)
REAL*8 EL1,ELEM,EL2D
INTEGER H1,H2
COMMON EL1(10),EN1(10),ICAN(10),INT(10)

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22:13(4,)  

SUBROUTINE SURF2(IRIG)  

REAL*8 EL1,ELEM,EL2D  

INTEGER H1,H2  

COMMON EL1(10),EN1(10),ICAN(10),INT(10)

```

5.	COMMON EL(50),EN(50),FWT(50),XLAM(50),AA(50)	GRE1275
6.	COMMON COUNT(4096),SPSM(4096),DER(4096)	GRE1276
7.	COMMON ISTAR,IFND,INDEX,JMAX,FMEDIO	GRE1277
8.	COMMON AREA(50),FBR(50),FWHM(50)	GRE1278
9.	COMMON LS(50),CFNTR(50),LD(50),PMOL(50)	GRE1279
10.	COMMON DFCAY(10),VALA(10),RESO(10)	GRE1280
11.	COMMON TDEC1,TIRR,TCON,ALFA,BETA,FLUX,WEIGHT	GRE1281
12.	COMMON STAND(6,50),NELST,IREL	GRE1282
13.	COMMON EWHI	GRE1283
14.	COMMON/ET1/NTRP,IO(5),MO(5),ANOME(5),ASAMPL(3)	GRE1284
15.	COMMON/ET2/BPM,ATOMI,PERC,AKEV,CPM,AREA0,UGM,LARG,H1,H2,J1	GRE1285
16.	COMMON/ET3/ENERG(50)	GRE1286
17.	COMMON/ET7/C1,C2,JDOUB,KKK,IND2	GRE1287
18.	COMMON/ET8/EDEC,ESAT,DENS	GRE1288
19.	COMMON/ET9/EL2D(50),VCD(50),PID(50),CCD(50),VID(50),CL1D(50),IPICD	GRE1289
20.	COMMON/ET11/EL2PREL(50),EL3PREL(50),EL2DRE(50)	GRE1290
21.	DIMENSION SIGMA(50)	GRE1291
22.	DATA STAN/,STAN/,/	GRE1292
23.	DATA BLANK/,/,/	GRE1293
24.		GRE1294
25.	JJJ=J1	GRE1295
26.	IC1=C1	GRE1296
27.	IC2=C2	GRE1297
28.	JC1=IC1+1	GRE1298
29.	JC2=IC2+1	GRE1299
30.	ASPM1=AMAX1(SPSM(IC1),SPSM(JC1))	GRE1300
31.	ASPM2=AMAX1(SPSM(IC2),SPSM(JC2))	GRE1301
32.	IF(ASPM1.NE.SPSM(IC1)) IC1=JC1	GRE1302
33.	IF(ASPM2.NE.SPSM(IC2)) IC2=JC2	GRE1303
34.	IND=1	GRE1304
35.	GOTO(10,20,30),JDOUB	GRE1305
36.	30 JDOUB=2	GRE1306
37.	D01 K1=1,JDOUB	
38.	KKK=K1	GRE1309
39.	GOTO(11,21),KKK	
40.	10 D01 K2=1,JDOUB	
41.	KKK=K2	GRE1312
42.	11 AR=SPSM(IC1)	GRE1313
43.	W=AR	GRE1314
44.	IL=IC1-1	GRE1315
45.	IPICD=IPICD+1	GRE1316
46.	GOTO 4	
47.	20 D01 K3=2,JDOUB	
48.	KKK=K3	GRE1319
49.	21 AR=SPSM(IC2)	GRE1320
50.	W=AR	GRE1321
51.	IL=IC2+1	GRE1322
52.	IND=1	GRE1323
53.	IPICD=IPICD+1	GRE1324
54.	4 IF(W-SPSM(IL))2,3,3	GRE1325
55.	3 IND=IND+1	GRE1326
56.	AR=AR+SPSM(IL)	GRE1327
57.	W=SPSM(IL)	GRE1328
58.	GOTO(40,50),KKK	GRE1329
59.	40 IL=IL-1	GRE1330
60.	GOTO 4	GRE1331
61.	50 IL=IL+1	GRE1332
62.	GOTO 4	GRE1333
63.	2 AR=AR*2	GRE1334
64.	IF(KKK.NE.1) GOTO 60	GRE1335
65.	IL=IL+1	GRE1336
66.	AR=AR-SPSM(IC1)	GRE1337
67.		GRE1338
68.	IF(IREL.EQ.1) GO TO 202	GRE1339
69.	EL2D(IPICD)=ELEM(JJJ)	GRE1340
70.	GO TO 203	GRE1341
71.	202 EL2DRE(IPICD)=STAND(1,JJJ)	GRE1342
72.	203 VCD(IPICD)=COUNT(IC1)	GRE1343
73.	P1D(IPICD)=ENERG(JMAX-1)	GRE1344
74.	CCD(IPICD)=C1	GRE1345
75.	VID(IPICD)=IL	GRE1346
76.	CL1D(IPICD)=COUNT(IL)	GRE1347
77.		GRE1348
78.	GOTO 70	GRE1349
79.	60 AR=AR-SPSM(IC2)	GRE1350
80.	IL=IL-1	GRE1351
81.		GRE1352
82.	IF(IREL.EQ.1) GO TO 200	GRE1353
83.	EL2D(IPICD)=ELEM(JJJ)	GRE1354
84.	GO TO 201	GRE1355
85.	200 EL2DRE(IPICD)=STAND(1,JJJ)	GRE1356
86.	201 VCD(IPICD)=COUNT(IC2)	GRE1357

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87. P1D(IPICD)=ENERG(JMAX)
88. CCD(IPICD)=C2
89. V1D(IPICD)=IL
90. CL1D(IPICD)=COUNT(IL)
91. C
92. 70 FIND=IND*2
93. FIND=FIND-1
94. FOND=W*FIND
95. ARE=AR-FOND
96. CPM=ARE/TCON
97. ERS=SQRT(AR+FIND*FOND/2.0)
98. ERR1=ERS/TCON
99. PERC=ERR1*100./CPM
100. AREAQ=CPM*FOND
101. IF(ASAMPL(1).EQ.STAN) STAN(6,JJJ)=AREAQ/STAN(5,JJJ)
102. IF(IREL.EQ.1) GO TO 100
103. UGM=AA(JJJ)*AREAQ/(FLUX*FSAT)
104. PPM=UGM/WEIGHT
105. ATOMT=6.02E17*UGM*DENS/(PMOL(JJJ)*WEIGHT)
106. GO TO 101
107. 100 IF(STAN(6,JJJ).NE.BLANK) GO TO 99
108. PPMERLANK
109. GO TO 101
110. 99 PPM=(1./STAN(6,JJJ))*(AREAQ/WEIGHT)
111. C
112. 101 H1=TL
113. IF(KKK.EQ.1) LARG=IC1-TL
114. IF(KKK.EQ.2) LARG=IL-IC2
115. IF(KKK.EQ.1) IWR=6
116. IF(KKK.EQ.2) IWR=7
117. CALL PRINT(IWR,IP1G)
118. 1 CONTINUE
119. RETURN
120. END

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TN 363 IBANK 99 DBANK 13922 COMMON
 SPECTRA,SENSIT
 R1 *06/16/82-22:13(2,)
 1. SUBROUTINE SENSIT(RAD)
 2. INTEGER H1,H2
 3. REAL*8 EL1,ELEM
 4. COMMON EL1(10),EN1(10),ICAN(10),INT(10)
 5. COMMON ELEM(50),EN(50),FWT(50),XLAM(50),AA(50)
 6. COMMON COUNT(4096),SPSM(4096),DER(4096)
 7. COMMON ISTAR,IEND,INDEX,JMAX,FMEDIO
 8. COMMON AREA(50),ERR(50),FWHM(50)
 9. COMMON LS(50),CENTR(50),LD(50),PMOL(50)
 10. COMMON DECAY(10),VALA(10),PESQ(10)
 11. COMMON TDEC1,TIRR,TCON,ALFA,BETA,FLUX,WEIGHT
 12. COMMON ET2/PPM,ATOMI,PERC,AKEV,CPM,AREAQ,UGM,LARG,H1,H2,J1
 13. C
 14. SP=0.
 15. CAPP=4./1.66
 16. DO 1 J=H1,H2
 17. 1 SP=SP+COUNT(J)
 18. RAD=CAPP*(1.+SQRT(1.+2.*SP))
 19. RETURN
 20. END

TN 34 IBANK 16 DBANK 13041 COMMON
 SPECTRA,PRINT
 R1 *06/16/82-22:13(2,)
 1. SUBROUTINE PRINT(IWR,IP1G)
 2. C ***** IN QUESTA ROUTINE SI FANNO TUTTE LE STAMPE *****
 3. C
 4. REAL*8 EL1,ELEM,EL2D
 5. INTEGER H1,H2
 6. C
 7. COMMON EL1(10),EN1(10),ICAN(10),INT(10)
 8. COMMON ELEM(50),EN(50),FWT(50),XLAM(50),AA(50)
 9. COMMON COUNT(4096),SPSM(4096),DER(4096)
 10. COMMON ISTAR,IEND,INDEX,JMAX,FMEDIO
 11. COMMON AREA(50),ERR(50),FWHM(50)
 12. COMMON LS(50),CENTR(50),LD(50),PMOL(50)
 13. COMMON DECAY(10),VALA(10),PESQ(10)
 14. COMMON TDEC1,TIRR,TCON,ALFA,BETA,FLUX,WEIGHT
 15. COMMON STAN(6,50),NELST,IREL
 16. COMMON ET1/NIRR,IO(5),MO(5),ANOME(5),ASAMPL(3)
 17. COMMON ET2/PPM,ATOMI,PERC,AKEV,CPM,AREAQ,UGM,LARG,H1,H2,J1
 18. COMMON ET3/ENERG(50)
 19. COMMON ET7/C1,C2,JDUR,KKK,IND2
 20. COMMON ET9/EL2D(50),VCD(50),P1D(50),CCD(50),V1D(50),CL1D(50),IPICD

21.	C	DATA BLANK/, ,/	GRF1433
22.		JJJ=J1	GRF1434
23.		J=J1	GRF1435
24.		IF (IRIG-53) 20,21,21	GRF1436
25.	21	WRITE (6,906)	GRF1437
26.		IRIG=0	GRF1438
27.		WRITE (6,902)	GRF1439
28.		WRITE (6,900)	GRF1440
29.		WRITE (6,901)	GRF1441
30.		WRITE (6,900)	GRF1442
31.		WRITE (6,902)	GRF1443
32.		IRIG=IRIG+5	GRF1444
33.			GRF1445
34.	C	20 GO TO (1,2,3,4,5,6,7,8,9,10),IWP	GRF1446
35.			GRF1447
36.	C	1 CONTINUE	GRF1448
37.		WRITE (6,1008) NIRR,TIRR,(TO(K),K=1,5),FLUX	GRF1449
38.		WRITE (6,1009) TCON,(MO(K),K=1,5),TDEC1	GRF1450
39.		WRITE (6,1012) (ASAMPL(K),K=1,3),(ANOMF(K),K=1,5)	GRF1451
40.		WRITE (6,1011)	GRF1452
41.		GO TO 100	GRF1453
42.			GRF1454
43.	C	2 CONTINUE	GRF1455
44.		WRITE (6,1008) NIRR,TIRR,(TO(K),K=1,5),FLUX	GRF1456
45.		WRITE (6,1009) TCON,(MO(K),K=1,5),TDEC1	GRF1457
46.		WRITE (6,1010) (ASAMPL(K),K=1,3),WEIGHT,ALFA,BETA,(ANOME(K),K=1,5)	GRF1458
47.		WRITE (6,1011)	GRF1459
48.		GO TO 100	GRF1460
49.			GRF1461
50.	C	3 IRIG=0	GRF1462
51.		WRITE (6,902)	GRF1463
52.		WRITE (6,900)	GRF1464
53.		WRITE (6,901)	GRF1465
54.		WRITE (6,900)	GRF1466
55.		WRITE (6,902)	GRF1467
56.		IRIG=IRIG+5	GRF1468
57.		GO TO 100	GRF1469
58.			GRF1470
59.	C	4 CONTINUE	GRF1471
60.		WRITE (6,900)	GRF1472
61.		IF (IREL.EQ.1) GO TO 200	GRF1473
62.		WRITE (6,904) ELEM(J),EN(J)	GRF1474
63.		GO TO 201	GRF1475
64.	200	WRITE (6,1904) STAND(1,J),EN(J)	GRF1476
65.	201	IRIG=IRIG+2	GRF1477
66.		GO TO 100	GRF1478
67.			GRF1479
68.	C	5 CONTINUE	GRF1480
69.		WRITE (6,900)	GRF1481
70.		IF (IREL.EQ.1) GO TO 202	GRF1482
71.		WRITE (6,905) ELEM(J),PPM,AA(J),EN(J),ATOMI	GRF1483
72.		GO TO 203	GRF1484
73.	202	IF (PPM.NE.BLANK) GO TO 300	GRF1485
74.		WRITE (6,900)	GRF1486
75.		WRITE (6,3000) STAND(1,J),EN(J)	GRF1487
76.		GO TO 203	GRF1488
77.	300	WRITE (6,1905) STAND(1,J),PPM,EN(J)	GRF1489
78.	203	IRIG=IRIG+2	GRF1490
79.		GO TO 100	GRF1491
80.			GRF1492
81.	C	6 CONTINUE	GRF1493
82.		WRITE (6,900)	GRF1494
83.		IF (IREL.EQ.1) GO TO 204	GRF1495
84.		WRITE (6,909) ELEM(JJJ),PPM,PERC,AA(JJJ),EN(JJJ),ENERG(JMAX-1),H1,L	GRF1496
85.		LARG,CPM,ENERG(JMAX-1),ENERG(JMAX),ATOMI	GRF1497
86.		GO TO 205	GRF1498
87.	204	IF (PPM.NE.BLANK) GO TO 301	GRF1499
88.		WRITE (6,900)	GRF1500
89.		WRITE (6,3000) STAND(1,JJJ),EN(JJJ)	GRF1501
90.		GO TO 205	GRF1502
91.	301	WRITE (6,1909) STAND(1,JJJ),PPM,PERC,EN(JJJ),ENERG(JMAX-1),H1,LARG,	GRF1503
92.		1CPM,ENERG(JMAX-1),ENERG(JMAX)	GRF1504
93.	205	IRIG=IRIG+2	GRF1505
94.		GO TO 100	GRF1506
95.			GRF1507
96.	C	7 CONTINUE	GRF1508
97.		WRITE (6,900)	GRF1509
98.		IF (IREL.EQ.1) GO TO 206	GRF1510
99.		WRITE (6,910) ELEM(JJJ),PPM,PERC,AA(JJJ),EN(JJJ),ENERG(JMAX),H1,LAR	GRF1511
100.		1G,CPM,ENERG(JMAX-1),ENERG(JMAX),ATOMI	GRF1512
101.		GO TO 207	GRF1513
102.			GRF1514

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103. 206 IF (PPM.NE.BLANK) GO TO 302
104. WRITE (6,900)
105. WRITE (6,3000) STAND(1,JJJ),EN(JJJ)
106. GO TO 207
107. 302 WRITE (6,1910) STAND(1,JJJ),PPM,PERC,EN(JJJ),FNERG(JMAX),H1,LARG,CP
108. 1M,FNERG(JMAX-1),ENERG(JMAX)
109. 207 IRIG=IRIG+2
110. GO TO 100
111. C
112. 8 CONTINUE
113. WRITE (6,900)
114. IRIG=IRIG+2
115. GO TO 100
116. C
117. 9 CONTINUE
118. WRITE (6,900)
119. IF (IREL.EQ.1) GO TO 208
120. WRITE (6,903) ELEM(J),PPM,PERC,AA(J),EN(J),AKEV,H1,LARG,CPM,AREA0,
121. 1UGM,ATOMI
122. GO TO 209
123. 208 IF (PPM.NE.BLANK) GO TO 303
124. WRITE (6,900)
125. WRITE (6,3000) STAND(1,J),EN(J)
126. GO TO 209
127. 303 WRITE (6,1903) STAND(1,J),PPM,PERC,EN(J),AKEV,H1,LARG,CPM,AREA0
128. 209 IRIG=IRIG+2
129. GO TO 100
130. C
131. 10 CONTINUE
132. WRITE (6,900)
133. IF (IREL.EQ.1) GO TO 210
134. WRITE (6,907) ELEM(J),PPM,PERC,AA(J),EN(J),AKEV,H1,LARG,CPM,AREA0,
135. 1UGM,ATOMI
136. GO TO 211
137. 210 IF (PPM.NE.BLANK) GO TO 304
138. WRITE (6,900)
139. WRITE (6,3000) STAND(1,J),EN(J)
140. GO TO 211
141. 304 WRITE (6,1907) STAND(1,J),PPM,PERC,EN(J),AKEV,H1,LARG,CPM,AREA0
142. 211 IRIG=IRIG+2
143. C
144. 900 FORMAT (1H,,I,,10X,,I,,11X,,I,,9X,,I,,11X,,I,,9X,,I,,10X
145. 1,,I,,11X,,I,,11X,,I,,11X,,I,,11X,,I,,)
146. 901 FORMAT (1H,,I,,ELEM,,I,,P.P.M.,,I,,FRR.0/0,,I,,I,,VALOPE,,A,,I,,KEV,,A
147. 1(T),,I,,KEV (S),,I,,LS,,WT,,I,,C.P.M.,,I,,AREA ZERO,,I,,UGM,,I,,A
148. 2I/CC,,I,,)
149. 902 FORMAT(1H,,I-----I-----I-----I-----I-----I-----I-----
150. 1-----I-----I-----I-----I-----I-----I-----
151. 2-----I,,)
152. 903 FORMAT (1H,,I,,A8,,I,,1PE9.3,,I,,0PF6.2,,I,,I,,1PE9.3,,I
153. 1,,0PF7.2,,I,,F7.2,,I,,2I4,,I,,1PE9.3,,I,,E9.3,,I,,F9.3,,
154. 2I,,E9.3,,I,,)
155. 904 FORMAT (1H,,I,,A8,,I,,OUT SCALE,,I,,9X,,I,,I,,11X,,I,,F7.2,,I,,
156. 18X,,I,,10X,,I,,11X,,I,,11X,,I,,11X,,I,,)
157. 905 FORMAT (1H,,I,,A8,,I,,(,,1PE9.3,,)I,,9X,,I,,I,,E9.3,,I,,0PF7.2,,
158. 1,,I,,9X,,I,,10X,,I,,11X,,I,,11X,,I,,11X,,I,,1PE9.3,,I,,)
159. 906 FORMAT(1H1)
160. 907 FORMAT (1H,,I,,A8,,I,,1PE9.3,,I,,0PF6.2,,I,,I,,1PE9.3,,I
161. 1,,0PF7.2,,I,,F7.2,,I,,2I4,,I,,1PE9.3,,I,,E9.3,,I,,F9.3,,
162. 2I,,E9.3,,I,,)
163. 909 FORMAT (1H,,I,,A8,,I,,1PE9.3,,I,,0PF6.2,,I,,I,,1PE9.3,,I
164. 1,,0PF7.2,,I,,F7.2,,I,,2I4,,I,,1PE9.3,,I,,PRIMO (,,0PF7.2,,1
165. 2X,F7.2,,),,I,,1PE9.3,,I,,)
166. 910 FORMAT (1H,,I,,A8,,I,,1PE9.3,,I,,0PF6.2,,I,,I,,1PE9.3,,I
167. 1,,0PF7.2,,I,,F7.2,,I,,2I4,,I,,1PE9.3,,I,,SECONDO(,,0PF7.2,,1
168. 2X,F7.2,,),,I,,1PE9.3,,I,,)
169. 1008 FORMAT (1H,,--IRRAGGIAMENTO --,I7,2X,F12.3,,MINUTI DATA,,5I
170. 15,5X,FLUSSO,3X,1PE10.4/)
171. 1009 FORMAT (1H,,--MISURA --TEMPO CONTEGGIO,,F8.3,,MINUTI DAT
172. 1A,,5I5.5X,DECADIMENTO,,F12.3,,MINUTI,/)
173. 1010 FORMAT (1H,,--CAMPIONE --,2X,3A4,5X,PESO,,F12.8,,6 CALIB.
174. 1Y=,,F7.3,,Xx,,F8.3,8X,PROBLEMA,,5A4)
175. 1011 FORMAT (1H,,93X,-----/)
176. 1012 FORMAT (1H,,--CAMPIONE --,2X,3A4,67X,PROBLEMA,,5A4)
177. C
178. 1903 FORMAT(1H,,I,,A4,4X,,I,,1PE9.3,,I,,0PF6.2,,I,,I,,9X,,I,,9
179. 1F7.2,,I,,F7.2,,I,,2I4,,I,,1PE9.3,,I,,E9.3,,I,,9X,,I,,9
180. 2X,,I,,)
181. 1904 FORMAT(1H,,I,,A4,4X,,I,,OUT SCALE,,I,,9X,,I,,I,,11X,,I,,F7.2,,I
182. 1,,8X,,I,,10X,,I,,11X,,I,,11X,,I,,11X,,I,,)
183. 1905 FORMAT (1H,,I,,A4,4X,,I,,(,,1PE9.3,,)I,,9X,,I,,I,,9X,,I,,0PF7.2
184.

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APPENDIX J

LISTING OF THE INPUT DATA FOR
SAMPLE PROBLEM

APPENDIX K

LISTING OF THE OUTPUT OF SAMPLE
PROBLEM

0	1	2	3	4	5	6	7	8	9	10
0	800.	12.	13.	48.	172.	104.	146.	136.	72.	60.
1	35.	15.	1.	4.	2.	1.	4.	4.	4.	1.
2	30.	47.	70.	17.	153.	278.	360.	608.	13.	24.
3	28.	1964.	6401.	1894.	4068.	3777.	3404.	222.	1540.	5366.
4	22.	2234.	1217.	1943.	2254.	2247.	2043.	222.	2563.	2177.
5	22.	2234.	2233.	2236.	2249.	2247.	2477.	225.	2252.	2152.
6	22.	2234.	2233.	2236.	2249.	2247.	2477.	225.	2252.	2152.
7	22.	2234.	2233.	2236.	2249.	2247.	2477.	225.	2252.	2152.
8	22.	2234.	2233.	2236.	2249.	2247.	2477.	225.	2252.	2152.
9	22.	2234.	2233.	2236.	2249.	2247.	2477.	225.	2252.	2152.
10	22.	2234.	2233.	2236.	2249.	2247.	2477.	225.	2252.	2152.

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- IRRAGGIAMENTO - 91 2400.000 MINUTI DATA 1974 12 15 24 0 FLUSSO .0000
 - MISURA - TEMPO CONTEGGIO 13.333 MINUTI DATA 1975 1 7 10 30 DECARINMENTO 32310.000 MINUTTY
 - CAMPIONE - CALIB PROBLEMA ALFALEA

N.	ELEM	EN	PESO (G)	CENTR	C P M	AREA ZFRO	A SPER	A TFOR	FATT CORR
1	EU1	245.00	.000000000	345.51	3.122+002	3.133+002	.000	1.285+002	.000
2	EU2	244.00	.000000000	345.81	3.122+002	3.133+002	.000	1.285+002	.000
3	EU3	778.40	.000000000	780.01	1.726+002	1.732+002	.000	1.286+002	.000
4	EU4	963.50	.000000000	964.06	1.519+002	1.524+002	.000	1.286+002	.000
5	CO2	1332.40	.000000000	1333.83	1.341+003	1.351+003	.000	1.286+002	.000
6	EU5	1408.00	.000000000	1409.44	1.473+002	1.478+002	.000	1.286+002	.000

VALORE MEDIO DEL FATTORE CORRETTIVO SU 6 PICCHI .000

CALIBRAZIONE Y= 1.000*x+ -1.712

[illegible]

IRRAGGIAMENTO - 91											
MISURA - TEMPO CONTEGGIO											
CAMPIONE - STANG2											
PF50 .0018900 6 CALIP. Y= 1.000*X+ -1.712											
DATA 1974 12 15 24 0 FLUSSO .0000											
DATA 1975 1 9 10 0 DECONTAMENTO 35160.000 MINUTI											
VALORE A KEV (T) KEV (S) LS WT C D M APFA ZEP0 UCM AT/CC											
CE				145.40	1333.16	1330	10	7.392+000	7.457+000		
CO	4.900+000	4.22		1332.48							
CR	9.000+000	60.03	*	320.08	320.53	320	5	1.362+000	2.502+000		
CS				795.80							
EU	1.500+000	7.57		1408.11	1408.61	1404	13	4.275+000	4.291+000		
HF	7.500+000	5.99		482.60	482.71	470	11	2.431+001	3.540+001		
HG				270.20							
IN				1300.00							
NI				810.81							
RB	2.340+002	19.85		1076.60	1077.16	1074	11	4.284+000	1.377+001		
SB				1691.00							
SC	3.900+000	1.35		889.25	889.65	886	11	1.083+002	1.323+002		
TA				1221.60							
TB				870.33							
TH	2.520+001	2.65		311.90	312.75	308	13	8.126+001	1.519+002		
YB	1.000+000	14.88		197.80	193.26	192	7	9.813+000	1.663+001		
YB	1.000+000	51.39	*	197.80	198.91	198	6	2.424+000	4.108+000		
ZN	7.490+001	9.37		1115.51	1115.75	1109	8	1.247+001	99100 (1115.75 1120.98)		
ZN	7.400+001	1.29		1115.51	1120.98	1118	10	8.395+001	9.995+001		
ZR	3.160+002	37.47		765.80	766.41	764	8	2.478+000	3.215+000		
BA	1.950+003	26.80		216.00	216.91	214	10	7.092+000	3.051+001		
FC	1.888+004	2.73		1099.27	1099.71	1094	15	4.296+001	6.236+001		

LIBRERIA METODO RELATIVO

ELEM	EN(KEV)	RES(KEV)	DEC(MIN)	UGM	CPMO/UGM
CE	145.40	3.000	1.502-005	3.136-001	
CO	1332.48	3.200	2.496-007	9.256-003	8.056+002
CR	320.00	3.070	1.730-005	1.700-002	1.472+002
CS	795.80	3.120	5.980-007	2.834-003	
EU	1408.11	3.220	1.055-007	2.834-003	1.514+003
HF	482.60	3.100	1.069-005	1.417-002	2.498+003
HG	279.20	3.000	1.035-005	8.312-005	
IN	1300.00	3.200	9.820-006	7.367-005	
NI	810.81	3.100	6.750-006	1.209-002	
RB	1076.60	3.150	3.320-005	4.420-001	3.114+001
SB	1691.00	3.250	8.030-006	1.001-004	
SC	889.25	3.120	5.675-006	7.367-003	1.795+004
TA	1221.60	3.200	4.185-006	1.719-003	
TB	879.33	3.100	6.592-006	9.823-004	
TH	311.90	3.050	1.780-005	4.760-002	3.192+003
YB	197.80	3.000	1.500-005	1.889-003	2.174+003
ZN	1115.51	3.170	1.964-006	1.415-001	6.358+002
ZR	765.89	3.100	7.404-006	5.969-001	5.386+000
BA	216.00	3.000	4.150-005	3.684+000	8.283+000
FE	1099.27	3.170	1.060-005	3.566+001	1.749+000

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77 78 79 80 81 82 83 84 85 86 87 88 89 90 91 92 93 94 95 96 97 98 99 100 101 102 103 104 105 106 107 108 109 110 111 112 113 114 115 116 117 118 119 120 121 122 123 124 125 126 127 128 129 130 131 132 133 134 135 136 137 138 139 140 141 142 143 144 145 146 147 148 149 150 151 152 153 154 155 156 157 158 159 160 161 162 163 164 165 166 167 168 169 170 171 172 173 174 175 176 177 178 179 180 181 182 183 184 185 186 187 188 189 190 191 192 193 194 195 196 197 198 199 200 201 202 203 204 205 206 207 208 209 210 211 212 213 214 215 216 217 218 219 220 221 222 223 224 225 226 227 228 229 230 231 232 233 234 235 236 237 238 239 240 241 242 243 244 245 246 247 248 249 250 251 252 253 254 255 256 257 258 259 260 261 262 263 264 265 266 267 268 269 270 271 272 273 274 275 276 277 278 279 280 281 282 283 284 285 286 287 288 289 290 291 292 293 294 295 296 297 298 299 300 301 302 303 304 305 306 307 308 309 310 311 312 313 314 315 316 317 318 319 320 321 322 323 324 325 326 327 328 329 330 331 332 333 334 335 336 337 338 339 340 341 342 343 344 345 346 347 348 349 350 351 352 353 354 355 356 357 358 359 360 361 362 363 364 365 366 367 368 369 370 371 372 373 374 375 376 377 378 379 380 381 382 383 384 385 386 387 388 389 390 391 392 393 394 395 396 397 398 399 400 401 402 403 404 405 406 407 408 409 410 411 412 413 414 415 416 417 418 419 420 421 422 423 424 425 426 427 428 429 430 431 432 433 434 435 436 437 438 439 440 441 442 443 444 445 446 447 448 449 450 451 452 453 454 455 456 457 458 459 460 461 462 463 464 465 466 467 468 469 470 471 472 473 474 475 476 477 478 479 480 481 482 483 484 485 486 487 488 489 490 491 492 493 494 495 496 497 498 499 500 501 502 503 504 505 506 507 508 509 510 511 512 513 514 515 516 517 518 519 520 521 522 523 524 525 526 527 528 529 530 531 532 533 534 535 536 537 538 539 540 541 542 543 544 545 546 547 548 549 550 551 552 553 554 555 556 557 558 559 560 561 562 563 564 565 566 567 568 569 570 571 572 573 574 575 576 577 578 579 580 581 582 583 584 585 586 587 588 589 590 591 592 593 594 595 596 597 598 599 600 601 602 603 604 605 606 607 608 609 610 611 612 613 614 615 616 617 618 619 620 621 622 623 624 625 626 627 628 629 630 631 632 633 634 635 636 637 638 639 640 641 642 643 644 645 646 647 648 649 650 651 652 653 654 655 656 657 658 659 660 661 662 663 664 665 666 667 668 669 670 671 672 673 674 675 676 677 678 679 680 681 682 683 684 685 686 687 688 689 690 691 692 693 694 695 696 697 698 699 700 701 702 703 704 705 706 707 708 709 710 711 712 713 714 715 716 717 718 719 720 721 722 723 724 725 726 727 728 729 730 731 732 733 734 735 736 737 738 739 740 741 742 743 744 745 746 747 748 749 750 751 752 753 754 755 756 757 758 759 760 761 762 763 764 765 766 767 768 769 770 771 772 773 774 775 776 777 778 779 780 781 782 783 784 785 786 787 788 789 790 791 792 793 794 795 796 797 798 799 800 801 802 803 804 805 806 807 808 809 810 811 812 813 814 815 816 817 818 819 820 821 822 823 824 825 826 827 828 829 830 831 832 833 834 835 836 837 838 839 840 841 842 843 844 845 846 847 848 849 850 851 852 853 854 855 856 857 858 859 860 861 862 863 864 865 866 867 868 869 870 871 872 873 874 875 876 877 878 879 880 881 882 883 884 885 886 887 888 889 890 891 892 893 894 895 896 897 898 899 900 901 902 903 904 905 906 907 908 909 910 911 912 913 914 915 916 917 918 919 920 921 922 923 924 925 926 927 928 929 930 931 932 933 934 935 936 937 938 939 940 941 942 943 944 945 946 947 948 949 950 951 952 953 954 955 956 957 958 959 960 961 962 963 964 965 966 967 968 969 970 971 972 973 974 975 976 977 978 979 980 981 982 983 984 985 986 987 988 989 990 991 992 993 994 995 996 997 998 999 1000

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- IRRAGGIAMENTO - 91									
- MISURA - TEMPO CONTEGGIO									
- CAMPIONE - ALFA 374									
ELEM	P.P.M.	ERR.0/0	VALORE A	KFV (T)	KFV (S)	LS	WT	C P M	AT/CC
CE				145.40					
CO	(4.115-003)			1332.48					
CR	(2.480-001)			320.08					
CS				795.80					
EU	(1.345-003)			1408.11					
HF	(9.197-003)			482.60					
HG				279.20					
IN				1300.00					
NI				810.81					
PR	5.240+000	12.69		1076.60	1079.00	1078	6	1.423+000	9.601+000
SB				1691.00					
SC	(8.720-004)			889.25					
TA				1221.60					
TR				870.33					
TH	1.108+000	.92		311.90	313.32	310	11	7.102+001	1.863+002
YR	(1.862-002)			197.80					
ZN	3.966+000	.41		1115.51	1113.55	1111	9	1.195+002	1.329+002
ZR	(2.935+000)			765.80					
BA	(1.995+001)			216.00					
FE	(8.062+000)			1099.27					

- IRRAGGIAMENTO - 91
 - MISURA - TEMPO CONTEGGIO
 - CAMPIONE - ALFA 374
 2400.000 MINUTI
 1851.950 MINUTI
 PESO .05270000 G
 DATA 1974 12 15 24
 DATA 1975 1 22 15
 CALIP. Y= 1.000*X+
 PROBLEMA ALFALEA
 FLUSSO 0
 DECAPIMENTO 54180.000 MINUTI
 AT/CC

[illegible]

IRRAGGIAMENTO ~		2400.000 MINUTI		DATA 1974		12		15		24		0		FLUSSO		.0000	
MISURA - TEMPO CONTEGGIO ~		30.183 MINUTI		DATA 1975		1		23		8		0		DECADIMENTO		55200.000 MINUTI	
CAMPIONE - ALFA 656		PES0		.05630000 G		CALIB.		Y= 1.000*X4		-1.712		C P M		AREA ZERO		ALFALEA	
FLEM	P.P.M.	EPR.0/0	VALORE A	KEV (T)	KEV (S)	LS	WT	C P M	AREA ZERO	UGM	AT/CP						
CE				145.40													
CE				145.40													
Co	(1.245-001)			1332.48													
CR	(7.242+000)			320.08													
CS				795.80													
CS				795.80													
EU	(3.069-002)			1408.11													
HF	(2.737-001)			482.60													
HG				279.20													
IN				1300.00													
NI				810.81													
RB	(4.703+001)			1076.60													
SB				1691.00													
SC	(2.650-002)			889.25													
TA				1221.60													
TR				879.33													
TH	1.583+001	1.87		311.90	313.02	310	11	1.065+003	2.844+003								
YR	1.684+000	17.61		197.80	190.95	196	3	9.006+001	185.64	190.95							
YB	3.029+000	12.67		197.80	185.64	183	9	1.620+002	3.709+002								
ZN	5.953+001	.92		1115.51	1113.19	1111	9	1.912+003	2.131+003								
ZR	(8.876+001)			765.80													
BA	(5.950+002)			216.00													
FE	(2.499+002)			1099.27													

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- IRRAGGIAMENTO - 91											
- MISURA - TEMPO CONTEGGIO 151.017 MINUTI											
- CAMPIONE - ALFA 154											
PF50 .06150000 G CALIP. Y= 1.000*X+ -1.712											
PROBLEMA ALFALEA											
ELEM	P.P.M.	ERR.0/0	VALORE A	KFV (T)	KEV. (S)	LS	WT	C P M	AREA ZERO	UAM	AT/CC
CE											
CO	4.250+000	1.16		145.40	1726.61	1319	19	2.077+002	2.106+002		
CR	(3.126+000)			1332.48							
CS				320.08							
EU	(1.587-002)			795.80							
HF	(1.045-001)			1408.11							
HG				482.60							
IN				270.20							
NI				1300.00							
NT				810.81							
RR	2.598+002	5.01		1076.60	1070.61	1067	11	7.857+001	4.975+002		
RU	6.095+001	16.07		1076.60	1079.79	1078	8	1.844+001	1.167+002		
SA				1691.00							
SC	(1.147-002)			880.25							
TA				1221.60							
TR				870.33							
TH	1.570+001	.76		311.90	314.26	311	11	1.146+003	3.082+003		
YB	(2.123-001)			197.90							
ZN	1.921+001	1.25		1115.51	1109.64	1099	12	6.735+002	PRIMO (1109.64	1114.40)	
ZN	4.143+001	.43		1115.51	1114.40	1112	9	1.453+003	1.620+003		
ZR	(3.258+001)			765.90							
BA	(2.358+002)			216.00							
FE	(9.838+001)			1090.27							

- IRRAGGIAMENTO - 91 2400.000 MINUTI 133.833 MINUTI 04790000 G CALIF. Y= 1.000*X+ -1.712 PROBLEMA ALFA ALFA
- MISURA - TEMPO CONTEGGIO 133.833 MINUTI 04790000 G DATA 1974 12 15 24 0 FLUSSO .0000
- CAMPIONE - ALFA 19 PFSO .04790000 G DATA 1975 1 24 7 30 DECADIMENTO 56610.000 WNUITY

ELEM	P.P.M.	PR. 3/0	VALORF A	KFV (T)	KFV (S)	LS	WT	C P M	AREA ZERO	USM	AT/CC
CF				145.40							
CO	(7.026-002)			1332.48							
CR	(3.844+000)			320.08							
CS				795.80							
EU	(1.930-002)			1408.11							
HF	(1.262-001)			432.60							
HG				279.20							
IN				1300.00							
NT				810.81							
RR	2.572+002	6.25		1076.60	1070.52	1067	11	5.858+001	3.837+002		
RA	6.991+001	19.13		1076.60	1079.88	1077	9	1.592+001	1.043+002		
SA				1691.00							
SB				1691.00							
SC	(1.466-002)			889.25							
TA				1221.60							
TR				879.33							
TH	1.560+001	.93		311.90	314.15	311	11	8.709+002	2.306+003		
YR	(2.615-001)			197.80							
ZN	7.102+001	.36		1115.51	1114.22	1106	21	1.935+003	2.163+003		
ZR	(3.890+001)			765.80							
RA	(2.981+002)			216.00							
FE	(1.212+002)			1090.27							

77 78 79 80 81 82 83 84 85 86 87 88 89 90 91 92 93 94 95 96 97 98 99 100 101 102 103 104 105 106 107 108 109 110 111 112 113 114 115 116 117 118 119 120 121 122 123 124 125 126 127 128 129 130 131 132 133 134 135 136 137 138 139 140 141 142 143 144 145 146 147 148 149 150 151 152 153 154 155 156 157 158 159 160 161 162 163 164 165 166 167 168 169 170 171 172 173 174 175 176 177 178 179 180 181 182 183 184 185 186 187 188 189 190 191 192 193 194 195 196 197 198 199 200 201 202 203 204 205 206 207 208 209 210 211 212 213 214 215 216 217 218 219 220 221 222 223 224 225 226 227 228 229 230 231 232 233 234 235 236 237 238 239 240 241 242 243 244 245 246 247 248 249 250 251 252 253 254 255 256 257 258 259 260 261 262 263 264 265 266 267 268 269 270 271 272 273 274 275 276 277 278 279 280 281 282 283 284 285 286 287 288 289 290 291 292 293 294 295 296 297 298 299 300 301 302 303 304 305 306 307 308 309 310 311 312 313 314 315 316 317 318 319 320 321 322 323 324 325 326 327 328 329 330 331 332 333 334 335 336 337 338 339 340 341 342 343 344 345 346 347 348 349 350 351 352 353 354 355 356 357 358 359 360 361 362 363 364 365 366 367 368 369 370 371 372 373 374 375 376 377 378 379 380 381 382 383 384 385 386 387 388 389 390 391 392 393 394 395 396 397 398 399 400 401 402 403 404 405 406 407 408 409 410 411 412 413 414 415 416 417 418 419 420 421 422 423 424 425 426 427 428 429 430 431 432 433 434 435 436 437 438 439 440 441 442 443 444 445 446 447 448 449 450 451 452 453 454 455 456 457 458 459 460 461 462 463 464 465 466 467 468 469 470 471 472 473 474 475 476 477 478 479 480 481 482 483 484 485 486 487 488 489 490 491 492 493 494 495 496 497 498 499 500 501 502 503 504 505 506 507 508 509 510 511 512 513 514 515 516 517 518 519 520 521 522 523 524 525 526 527 528 529 530 531 532 533 534 535 536 537 538 539 540 541 542 543 544 545 546 547 548 549 550 551 552 553 554 555 556 557 558 559 560 561 562 563 564 565 566 567 568 569 570 571 572 573 574 575 576 577 578 579 580 581 582 583 584 585 586 587 588 589 590 591 592 593 594 595 596 597 598 599 600 601 602 603 604 605 606 607 608 609 610 611 612 613 614 615 616 617 618 619 620 621 622 623 624 625 626 627 628 629 630 631 632 633 634 635 636 637 638 639 640 641 642 643 644 645 646 647 648 649 650 651 652 653 654 655 656 657 658 659 660 661 662 663 664 665 666 667 668 669 670 671 672 673 674 675 676 677 678 679 680 681 682 683 684 685 686 687 688 689 690 691 692 693 694 695 696 697 698 699 700 701 702 703 704 705 706 707 708 709 710 711 712 713 714 715 716 717 718 719 720 721 722 723 724 725 726 727 728 729 730 731 732 733 734 735 736 737 738 739 740 741 742 743 744 745 746 747 748 749 750 751 752 753 754 755 756 757 758 759 760 761 762 763 764 765 766 767 768 769 770 771 772 773 774 775 776 777 778 779 780 781 782 783 784 785 786 787 788 789 790 791 792 793 794 795 796 797 798 799 800 801 802 803 804 805 806 807 808 809 810 811 812 813 814 815 816 817 818 819 820 821 822 823 824 825 826 827 828 829 830 831 832 833 834 835 836 837 838 839 840 841 842 843 844 845 846 847 848 849 850 851 852 853 854 855 856 857 858 859 860 861 862 863 864 865 866 867 868 869 870 871 872 873 874 875 876 877 878 879 880 881 882 883 884 885 886 887 888 889 890 891 892 893 894 895 896 897 898 899 900 901 902 903 904 905 906 907 908 909 910 911 912 913 914 915 916 917 918 919 920 921 922 923 924 925 926 927 928 929 930 931 932 933 934 935 936 937 938 939 940 941 942 943 944 945 946 947 948 949 950 951 952 953 954 955 956 957 958 959 960 961 962 963 964 965 966 967 968 969 970 971 972 973 974 975 976 977 978 979 980 981 982 983 984 985 986 987 988 989 990 991 992 993 994 995 996 997 998 999 1000 1001 1002 1003 1004 1005 1006 1007 1008 1009 1010 1011 1012 1013 1014 1015 1016 1017 1018 1019 1020 1021 1022 1023 1024 1025 1026 1027 1028 1029 1030 1031 1032 1033 1034 1035 1036 1037 1038 1039 1040 1041 1042 1043 1044 1045 1046 1047 1048 1049 1050 1051 1052 1053 1054 1055 1056 1057 1058 1059 1060 1061 1062 1063 1064 1065 1066 1067 1068 1069 1070 1071 1072 1073 1074 1075 1076 1077 1078 1079 1080 1081 1082 1083 1084

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URRAGGIAMENTO - 01											
MISURA - TEMPO CONTEGGIO 2409.000 MINUTI 2409.000 MINUTI 2409.000 MINUTI											
CAMPIONE - ALFA 756 PESO 36.100 VALORE A											
FLEM	P.P.M.	ERR.0/0	VALORE A	KEY (T)	KEY (S)	LS	WT	C P W	ARFA ZERO	UGM	AT/CC
CE				145.40		1324	13				
CO	7.922+000	.85		1332.48	1328.86	313	11	3.586+002	3.639+002		
CR	3.754+002	1.10		320.08	316.56			1.162+003	3.149+003		
CS				795.80							
EU	5.778+001	3.59		1409.11	1404.41	1401	11	4.957+001	4.987+001		
HF	3.655+000	3.18		482.60	478.35	476	9	2.811+002	5.205+002		
HG				279.20							
IN				1300.00							
NT				810.81		1068	13	1.164+002	7.884+002		
RR	4.441+002	6.06		1076.60	1072.94						
SB				1691.00	885.28	879	17	3.915+003	5.430+003		
SC	5.306+000	.38		889.25							
TA				1221.60							
TR				879.33							
TH	3.596+000	4.60		311.90	308.42	306	9	2.345+002	6.542+002		
YB	2.394+000	12.22		197.80	194.54	201	5	1.250+002	SECOROT 199.06	104.50	
YH	2.956+000	8.54		197.80	189.06	187	9	1.543+002	3.664+002		
ZH	1.214+002	.29		1115.51	1116.70	1108	21	3.930+003	4.401+003		
ZR	2.975+002	16.55		765.80	761.75	758	12	5.961+001	9.134+001		
BA	(4.201+002)			216.00							
FE	1.466+004	.97		1099.27	1095.39	1091	13	7.934+002	1.462+003		

	1	2	3	4	5	6	7	8	9	10
0	6539	10563	10554	8652	6906	5370	3700	2391	1216	1015
1	5537	487	372	282	207	189	151	140	153	126
2	1073	114	116	111	121	167	152	126	280	485
3	7548	1851	7358	36688	14248	65484	154381	198507	106306	18427
4	138265	173581	136712	166689	162248	124130	121664	150618	145850	142670
5	82365	111389	101308	137978	103065	103526	103161	101614	117120	114670
6	82365	108579	100908	101086	90452	88052	916175	90086	89511	88874
7	82365	83944	82844	971964	80535	80281	91177	80559	79515	78773
8	82365	72844	72844	771704	771704	771704	771704	72207	71616	71250
9	82365	66670	66670	66670	66670	66670	66670	65336	64647	64537
10	82365	66154	66154	66154	66154	66154	66154	65558	64870	64537
11	82365	64746	64746	64746	64746	64746	64746	64573	64164	63898
12	82365	43314	43314	46744	46744	46744	46744	45173	44164	44307
13	82365	41844	41844	41844	41844	41844	41844	40846	40036	39280
14	82365	38771	38771	37719	37719	37719	37719	37381	36516	35820
15	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
16	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
17	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
18	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
19	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
20	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
21	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
22	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
23	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
24	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
25	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
26	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
27	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
28	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
29	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
30	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
31	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
32	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
33	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
34	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
35	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
36	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
37	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
38	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
39	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
40	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
41	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
42	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
43	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
44	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
45	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
46	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
47	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
48	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
49	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
50	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
51	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
52	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
53	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
54	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
55	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
56	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
57	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
58	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
59	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
60	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
61	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
62	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
63	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
64	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
65	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
66	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
67	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
68	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
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70	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
71	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
72	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
73	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
74	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
75	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020
76	82365	35528	35528	35528	35528	35528	35528	35413	34618	34020

-	IRRAGGIAMENTO -	91	2400.000	MINUTI	DATA	1974	12	15	24	0	FLUSSO	.00000
-	MISURA -	TEMPO CONTEGGIO	108.983	MINUTI	DATA	1975	1	25	17	30	DECADIMENTO	58650.000
-	CAMPIONE -	ALFA 515	PESO	.06230000	G	CALIP.	Y=	1.000	*X+	-1.712	PROBLEMA	ALFALEA

ELEM	P.P.M.	EPR.0/0	VALORE A	KEV (T)	KEV (S)	LS	WT	C P M	ADFA ZERO	UEM	AT/CC
CE				145.40							
CO	7.156+000	.83		1332.48	1328.74	1324	13	3.540+002	3.502+002		
CR	3.788+002	.99		320.08	316.44	313	11	1.259+003	3.474+003		
CS				795.80							
CU	7.119+001	3.13		1408.11	1404.31	1400	13	6.675+001	6.716+001		
FE	3.051+000	3.82		482.60	478.23	474	13	3.202+002	5.004+002		
HG				279.20							
IN				1300.00							
NI				810.81							
RR	4.027+002	5.46		1076.60	1072.75	1069	11	1.115+002	7.914+002		
SB				1691.00							
SC	1.791+002	37.39		889.25	893.30	892	3	1.436+001	PRIMO (893.30	898.12)	
SC	5.057+001	5.41		889.25	898.12	919	20	4.606+002	SECUNDO (893.30	898.12)	
SC	5.666+000	.30		880.25	885.15	880	15	4.543+003	6.337+003		
TA				1221.60							
TH				879.33							
TH	4.186+000	3.60		311.90	308.29	306	9	2.930+002	8.324+002		
YB	2.084+000	15.08		197.80	194.42	205	9	1.621+002	SECUNDO (189.00	194.42)	
YB	3.409+000	5.29		197.80	189.02	188	7	1.916+002	4.618+002		
ZN	2.758+001	1.26		1115.51	1111.92	1102	11	9.737+002	PRIMO (1111.92	1116.57)	
Zn	8.355+001	.33		1115.51	1116.57	1114	9	2.940+003	3.307+003		
ZP	2.627+002	14.61		765.80	761.69	759	10	5.700+001	9.914+001		
BA	(3.993+002)			216.00							
FE	1.532+004	.96		1099.27	1085.24	1099	15	8.962+002	1.660+003		

[illegible]

10290. 10661. 10672. 10683. 10694. 10705. 10716. 10727. 10738. 10749. 10760. 10771. 10782. 10793. 10804. 10815. 10826. 10837. 10848. 10859. 10870. 10881. 10892. 10903. 10914. 10925. 10936. 10947. 10958. 10969. 10980. 10991. 11002. 11013. 11024. 11035. 11046. 11057. 11068. 11079. 11090. 11101. 11112. 11123. 11134. 11145. 11156. 11167. 11178. 11189. 11200. 11211. 11222. 11233. 11244. 11255. 11266. 11277. 11288. 11299. 11310. 11321. 11332. 11343. 11354. 11365. 11376. 11387. 11398. 11409. 11420. 11431. 11442. 11453. 11464. 11475. 11486. 11497. 11508. 11519. 11530. 11541. 11552. 11563. 11574. 11585. 11596. 11607. 11618. 11629. 11640. 11651. 11662. 11673. 11684. 11695. 11706. 11717. 11728. 11739. 11750. 11761. 11772. 11783. 11794. 11805. 11816. 11827. 11838. 11849. 11860. 11871. 11882. 11893. 11904. 11915. 11926. 11937. 11948. 11959. 11970. 11981. 11992. 12003. 12014. 12025. 12036. 12047. 12058. 12069. 12080. 12091. 12102. 12113. 12124. 12135. 12146. 12157. 12168. 12179. 12190. 12201. 12212. 12223. 12234. 12245. 12256. 12267. 12278. 12289. 12300. 12311. 12322. 12333. 12344. 12355. 12366. 12377. 12388. 12399. 12410. 12421. 12432. 12443. 12454. 12465. 12476. 12487. 12498. 12509. 12520. 12531. 12542. 12553. 12564. 12575. 12586. 12597. 12608. 12619. 12630. 12641. 12652. 12663. 12674. 12685. 12696. 12707. 12718. 12729. 12740. 12751. 12762. 12773. 12784. 12795. 12806. 12817. 12828. 12839. 12850. 12861. 12872. 12883. 12894. 12905. 12916. 12927. 12938. 12949. 12960. 12971. 12982. 12993. 13004. 13015. 13026. 13037. 13048. 13059. 13070. 13081. 13092. 13103. 13114. 13125. 13136. 13147. 13158. 13169. 13180. 13191. 13202. 13213. 13224. 13235. 13246. 13257. 13268. 13279. 13290. 13301. 13312. 13323. 13334. 13345. 13356. 13367. 13378. 13389. 13400. 13411. 13422. 13433. 13444. 13455. 13466. 13477. 13488. 13499. 13510. 13521. 13532. 13543. 13554. 13565. 13576. 13587. 13598. 13609. 13620. 13631. 13642. 13653. 13664. 13675. 13686. 13697. 13708. 13719. 13730. 13741. 13752. 13763. 13774. 13785. 13796. 13807. 13818. 13829. 13840. 13851. 13862. 13873. 13884. 13895. 13906. 13917. 13928. 13939. 13950. 13961. 13972. 13983. 13994. 14005. 14016. 14027. 14038. 14049. 14060. 14071. 14082. 14093. 14104. 14115. 14126. 14137. 14148. 14159. 14170. 14181. 14192. 14203. 14214. 14225. 14236. 14247. 14258. 14269. 14280. 14291. 14302. 14313. 14324. 14335. 14346. 14357. 14368. 14379. 14390. 14401. 14412. 14423. 14434. 14445. 14456. 14467. 14478. 14489. 14500. 14511. 14522. 14533. 14544. 14555. 14566. 14577. 14588. 14599. 14610. 14621. 14632. 14643. 14654. 14665. 14676. 14687. 14698. 14709. 14720. 14731. 14742. 14753. 14764. 14775. 14786. 14797. 14808. 14819. 14830. 14841. 14852. 14863. 14874. 14885. 14896. 14907. 14918. 14929. 14940. 14951. 14962. 14973. 14984. 14995. 15006. 15017. 15028. 15039. 15050. 15061. 15072. 15083. 15094. 15105. 15116. 15127. 15138. 15149. 15160. 15171. 15182. 15193. 15204. 15215. 15226. 15237. 15248. 15259. 15270. 15281. 15292. 15303. 15314. 15325. 15336. 15347. 15358. 15369. 15380. 15391. 15402. 15413. 15424. 15435. 15446. 15457. 15468. 15479. 15490. 15501. 15512. 15523. 15534. 15545. 15556. 15567. 15578. 15589. 15600. 15611. 15622. 15633. 15644. 15655. 15666. 15677. 15688. 15699. 15710. 15721. 15732. 15743. 15754. 15765. 15776. 15787. 15798. 15809. 15820. 15831. 15842. 15853. 15864. 15875. 15886. 15897. 15908. 15919. 15930. 15941. 15952. 15963. 15974. 15985. 15996. 16007. 16018. 16029. 16040. 16051. 16062. 16073. 16084. 16095. 16106. 16117. 16128. 16139. 16150. 16161. 16172. 16183. 16194. 16205. 16216. 16227. 16238. 16249. 16260. 16271. 16282. 16293. 16304. 16315. 16326. 16337. 16348. 16359. 16370. 16381. 16392. 16403. 16414. 16425. 16436. 16447. 16458. 16469. 16480. 16491. 16502. 16513. 16524. 16535. 16546. 16557. 16568. 16579. 16590. 16601. 16612. 16623. 16634. 16645. 16656. 16667. 16678. 16689. 16700. 16711. 16722. 16733. 16744. 16755. 16766. 16777. 16788. 16799. 16810. 16821. 16832. 16843. 16854. 16865. 16876. 16887. 16898. 16909. 16920. 16931. 16942. 16953. 16964. 16975. 16986. 16997. 17008. 17019. 17030. 17041. 17052. 17063. 17074. 17085. 17096. 17107. 17118. 17129. 17140. 17151. 17162. 17173. 17184. 17195. 17206. 17217. 17228. 17239. 17250. 17261. 17272. 17283. 17294. 17305. 17316. 17327. 17338. 17349. 17360. 17371. 17382. 17393. 17404. 17415. 17426. 17437. 17448. 17459. 17470. 17481. 17492. 17503. 17514. 17525. 17536. 17547. 17558. 17569. 17580. 17591. 17602. 17613. 17624. 17635. 17646. 17657. 17668. 17679. 17690. 17701. 17712. 17723. 17734. 17745. 17756. 17767. 17778. 17789. 17800. 17811. 17822. 17833. 17844. 17855. 17866. 17877. 17888. 17899. 17910. 17921. 17932. 17943. 17954. 17965. 17976. 17987. 17998. 18009. 18020. 18031. 18042. 18053. 18064. 18075. 18086. 18097. 18108. 18119. 18130. 18141. 18152. 18163. 18174. 18185. 18196. 18207. 18218. 18229. 18240. 18251. 18262. 18273. 18284. 18295. 18306. 18317. 18328. 18339. 18350. 18361. 18372. 18383. 18394. 18405. 18416. 18427. 18438. 18449. 18460. 18471. 18482. 18493. 18504. 18515. 18526. 18537. 18548. 18559. 18570. 18581. 18592. 18603. 18614. 18625. 18636. 18647. 18658. 18669. 18680. 18691. 18702. 18713. 18724. 18735. 18746. 18757. 18768. 18779. 18790. 18801. 18812. 18823. 18834. 18845. 18856. 18867. 18878. 18889. 18900. 18911. 18922. 18933. 18944. 18955. 18966. 18977. 18988. 18999. 19010. 19021. 19032. 19043. 19054. 19065. 19076. 19087. 19098. 19109. 19120. 19131. 19142. 19153. 19164. 19175. 19186. 19197. 19208. 19219. 19230. 19241. 19252. 19263. 19274. 19285. 19296. 19307. 19318. 19329. 19340. 19351. 19362. 19373. 19384. 19395. 19406. 19417. 19428. 19439. 19450. 19461. 19472. 19483. 19494. 19505. 19516. 19527. 19538. 19549. 19560. 19571. 19582. 19593. 19604. 19615. 19626. 19637. 19648. 19659. 19670. 19681. 19692. 19703. 19714. 19725. 19736. 19747. 19758. 19769. 19780. 19791. 19802. 19813. 19824. 19835. 19846. 19857. 19868. 19879. 19890. 19901. 19912. 19923. 19934. 19945. 19956. 19967. 19978. 19989. 20000.

| IRRAGGIAMENTO - GI | | 2400.000 | | MINUTI | | DATA | | 1974 | | 1975 | | 1976 | | 1977 | | 1978 | | 1979 | | 1980 | | 1981 | | 1982 | | 1983 | | 1984 | | 1985 | | 1986 | | 1987 | | 1988 | | 1989 | | 1990 | | 1991 | | 1992 | | 1993 | | 1994 | | 1995 | | 1996 | | 1997 | | 1998 | | 1999 | | 2000 | | 2001 | | 2002 | | 2003 | | 2004 | | 2005 | | 2006 | | 2007 | | 2008 | | 2009 | | 2010 | | 2011 | | 2012 | | 2013 | | 2014 | | 2015 | | 2016 | | 2017 | | 2018 | | 2019 | | 2020 | | 2021 | | 2022 | | 2023 | | 2024 | | 2025 | | 2026 | | 2027 | | 2028 | | 2029 | | 2030 | | 2031 | | 2032 | | 2033 | | 2034 | | 2035 | | 2036 | | 2037 | | 2038 | | 2039 | | 2040 | | 2041 | | 2042 | | 2043 | | 2044 | | 2045 | | 2046 | | 2047 | | 2048 | | 2049 | | 2050 | | 2051 | | 2052 | | 2053 | | 2054 | | 2055 | | 2056 | | 2057 | | 2058 | | 2059 | | 2060 | | 2061 | | 2062 | | 2063 | | 2064 | | 2065 | | 2066 | | 2067 | | 2068 | | 2069 | | 2070 | | 2071 | | 2072 | | 2073 | | 2074 | | 2075 | | 2076 | | 2077 | | 2078 | | 2079 | | 2080 | | 2081 | | 2082 | | 2083 | | 2084 | | 2085 | | 2086 | | 2087 | | 2088 | | 2089 | | 2090 | | 2091 | | 2092 | | 2093 | | 2094 | | 2095 | | 2096 | | 2097 | | 2098 | | 2099 | | 2100 | | 2101 | | 2102 | | 2103 | | 2104 | | 2105 | | 2106 | | 2107 | | 2108 | | 2109 | | 2110 | | 2111 | | 2112 | | 2113 | | 2114 | | 2115 | | 2116 | | 2117 | | 2118 | | 2119 | | 2120 | | 2121 | | 2122 | | 2123 | | 2124 | | 2125 | | 2126 | | 2127 | | 2128 | | 2129 | | 2130 | | 2131 | | 2132 | | 2133 | | 2134 | | 2135 | | 2136 | | 2137 | | 2138 | | 2139 | | 2140 | | 2141 | | 2142 | | 2143 | | 2144 | | 2145 | | 2146 | | 2147 | | 2148 | | 2149 | | 2150 | | 2151 | | 2152 | | 2153 | | 2154 | | 2155 | | 2156 | | 2157 | | 2158 | | 2159 | | 2160 | | 2161 | | 2162 | | 2163 | | 2164 | | 2165 | | 2166 | | 2167 | | 2168 | | 2169 | | 2170 | | 2171 | | 2172 | | 2173 | | 2174 | | 2175 | | 2176 | | 2177 | | 2178 | | 2179 | | 2180 | | 2181 | | 2182 | | 2183 | | 2184 | | 2185 | | 2186 | | 2187 | | 2188 | | 2189 | | 2190 | | 2191 | | 2192 | | 2193 | | 2194 | | 2195 | | 2196 | | 2197 | | 2198 | | 2199 | | 2200 | | 2201 | | 2202 | | 2203 | | 2204 | | 2205 | | 2206 | | 2207 | | 2208 | | 2209 | | 2210 | | 2211 | | 2212 | | 2213 | | 2214 | | 2215 | | 2216 | | 2217 | | 2218 | | 2219 | | 2220 | | 2221 | | 2222 | | 2223 | | 2224 | | 2225 | | 2226 | | 2227 | | 2228 | | 2229 | | 2230 | | 2231 | | 2232 | | 2233 | | 2234 | | 2235 | | 2236 | | 2237 | | 2238 | | 2239 | | 2240 | | 2241 | | 2242 | | 2243 | | 2244 | | 2245 | | 2246 | | 2247 | | 2248 | | 2249 | | 2250 | | 2251 | | 2252 | | 2253 | | 2254 | | 2255 | | 2256 | | 2257 | | 2258 | | 2259 | | 2260 | | 2261 | | 2262 | | 2263 | | 2264 | | 2265 | | 2266 | | 2267 | | 2268 | | 2269 | | 2270 | | 2271 | | 2272 | | 2273 | | 2274 | | 2275 | | 2276 | | 2277 | | 2278 | | 2279 | | 2280 | | 2281 | | 2282 | | 2283 | |
|--------------------|--|----------|--|--------|--|
|--------------------|--|----------|--|--------|--|

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- IRRAGGIAMENTO - 91 2400.000 MINUTI DATA 1974 12 15 24 0 FLUSSO .0000
 - MISURA - TEMPO CONTEGGIO 1'41.783 MINUTI DATA 1975 1 27 4 30 DECADECIMENTO 60750.000 MINUTI
 - CAMPIONE - ALFA 274 PESO .055000000 G CALIP. Y= 1.000*Y+ -1.712 PROPLEVA ALFALEA

| ELEM | P.P.V. | ERR.0/0 | VALORE A | KEV (T) | KEV (S) | LS | WT | C P M | AREA 7ERO | UCM | AT/CC |
|------|-------------|---------|----------|---------|---------|------|----|-----------|-----------|-----|-------|
| CE | | | | 145.40 | | | | | | | |
| CO | (6.233-002) | | | 1332.48 | | | | | | | |
| CP | (3.609+000) | | | 320.08 | | | | | | | |
| CS | | | | 795.80 | | | | | | | |
| EU | (2.006-002) | | | 1408.11 | | | | | | | |
| HF | (1.202-001) | | | 482.60 | | | | | | | |
| HG | | | | 279.20 | | | | | | | |
| IN | | | | 1300.00 | | | | | | | |
| NI | | | | 810.81 | | 1076 | 11 | 3.969+001 | 2.982+002 | | |
| RR | 1.741+002 | 10.04 | | 1076.60 | 1079.50 | | | | | | |
| SB | | | | 1691.00 | | | | | | | |
| SC | (1.339-002) | | | 880.25 | | | | | | | |
| TA | | | | 1221.60 | | | | | | | |
| TB | | | | 879.33 | | | | | | | |
| TH | 1.309+001 | 1.10 | | 311.90 | 313.84 | 310 | 11 | 7.793+002 | 2.202+003 | | |
| YB | (2.513-001) | | | 197.80 | | | | | | | |
| ZN | 7.916+001 | .33 | | 1115.51 | 1113.89 | 1103 | 25 | 2.457+003 | 2.769+003 | | |
| ZR | (3.772+001) | | | 765.80 | | | | | | | |
| BA | (3.211+002) | | | 216.00 | | | | | | | |
| FE | (1.164+002) | | | 1099.27 | | | | | | | |

[illegible]

| | | | | | | | | | | |
|---|-------------|---------|----------|---------|---------|------|----|-----------|-----------|-----|
| - IRRAGGIAMENTO - 91 | | | | | | | | | | |
| - MISURA - TEMPO CONTEGGIO ~ 139.983 MINUTI | | | | | | | | | | |
| - CAMPIONE - ALFA 467 | | | | | | | | | | |
| PESO .06220000 G | | | | | | | | | | |
| CALIB. YE 1.000X+ | | | | | | | | | | |
| -1.712 | | | | | | | | | | |
| PROBLEMA ALFALEA | | | | | | | | | | |
| 63210.000 MINUTI | | | | | | | | | | |
| FLUSSO 0 | | | | | | | | | | |
| AT/CC | | | | | | | | | | |
| ELEM | P.P.M. | FRR.0/0 | VALORE A | KEV (T) | KEV (S) | LS | WT | C P M | AREA ZERO | UCM |
| CE | | | | 145.40 | | | | | | |
| CO | (5.114-002) | | | 1332.48 | | | | | | |
| CR | (3.482+000) | | | 320.08 | | | | | | |
| CS | | | | 795.80 | | | | | | |
| EU | (1.706-002) | | | 1408.11 | | | | | | |
| HF | (1.251-001) | | | 482.60 | | | | | | |
| HG | | | | 279.20 | | | | | | |
| IN | | | | 1300.00 | | | | | | |
| NI | | | | 810.81 | | | | | | |
| RI | 1.395+002 | 14.84 | | 1076.60 | 1078.99 | 1075 | 12 | 3.313+001 | 2.702+002 | |
| SH | | | | 1691.00 | | | | | | |
| SC | (1.169-002) | | | 889.25 | | | | | | |
| TA | | | | 1221.60 | | | | | | |
| TE | | | | 879.33 | | | | | | |
| TH | 1.209+001 | 1.18 | | 311.90 | 313.43 | 310 | 11 | 7.789+002 | 2.300+003 | |
| YB | (2.515-001) | | | 197.80 | | | | | | |
| ZN | 6.017+001 | .35 | | 1115.51 | 1113.64 | 1111 | 9 | 2.102+003 | 2.370+003 | |
| ZR | (3.889+001) | | | 765.80 | | | | | | |
| BA | (3.449+002) | | | 216.00 | | | | | | |
| FE | (1.109+002) | | | 1099.27 | | | | | | |

[illegible]

[illegible]

- IRRAGGIAMENTO - 91
 - MISURA - TEMPO CONTEGGIO 169.317 MINUTI
 - CAMPIONE - ALFA 910 PESO .06320000 G CALIP. Y= 1.000*X+ -1.712

| ELEM | P.P.M. | FRR.0/0 | VALORE A | KFV (T) | KFV (S) | LS | WT | C P M | AREA ZERO | UCM | ATACC |
|------|-------------|---------|----------|---------|---------|------|----|-----------|-----------|--------|-------|
| CE | | | | 145.40 | | | | | | | |
| Co | (4.702-002) | | | 1332.48 | | | | | | | |
| Cr | (2.939+000) | | | 320.08 | | | | | | | |
| CS | | | | 795.80 | | | | | | | |
| CU | (1.385-002) | | | 1408.11 | | | | | | | |
| HF | (9.976-002) | | | 482.60 | | | | | | | |
| HG | | | | 279.20 | | | | | | | |
| IN | | | | 1300.00 | | | | | | | |
| IN | | | | 1300.00 | | | | | | | |
| IN | | | | 1300.00 | | | | | | | |
| NT | | | | 810.81 | | | | | | | |
| RB | (2.088+001) | | | 1076.60 | | | | | | | |
| SB | | | | 1691.00 | | | | | | | |
| SC | (1.067-002) | | | 889.25 | | | | | | | |
| TA | | | | 1221.60 | | | | | | | |
| TH | | | | 879.33 | | | | | | | |
| TH | 9.351+000 | 1.21 | | 311.90 | 314.18 | 311 | 11 | 5.006+002 | 1.886+003 | | |
| YB | 1.607+000 | 13.78 | | 197.80 | 192.42 | 201 | 7 | 8.302+001 | 156.00 | 192.42 | |
| YB | 1.891+000 | 6.85 | | 197.80 | 186.82 | 185 | 8 | 9.978+001 | 2.599+002 | | |
| ZH | 4.197+001 | .38 | | 1115.51 | 1114.11 | 1112 | 9 | 1.406+003 | 1.687+003 | | |
| ZR | (3.060+001) | | | 765.80 | | | | | | | |
| BA | 7.073+002 | 33.50 | | 216.00 | 209.43 | 216 | 5 | 2.548+001 | 156.00 | 209.43 | |
| BA | 1.089+002 | 69.76 | | 216.00 | 208.47 | 204 | 6 | 7.167+000 | 1.041+002 | | |
| AFIN | (9.362+001) | | | 1090.27 | | | | | | | |